

# **BATS-R-US: a Multi-Physics and Multi-Application Code**

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## **M** Space Weather Modeling Framework

## **M** BATS-R-US

## **M** MHD equations with anisotropic proton pressure

- 🌐 Magnetosphere application

## **M** Electron physics

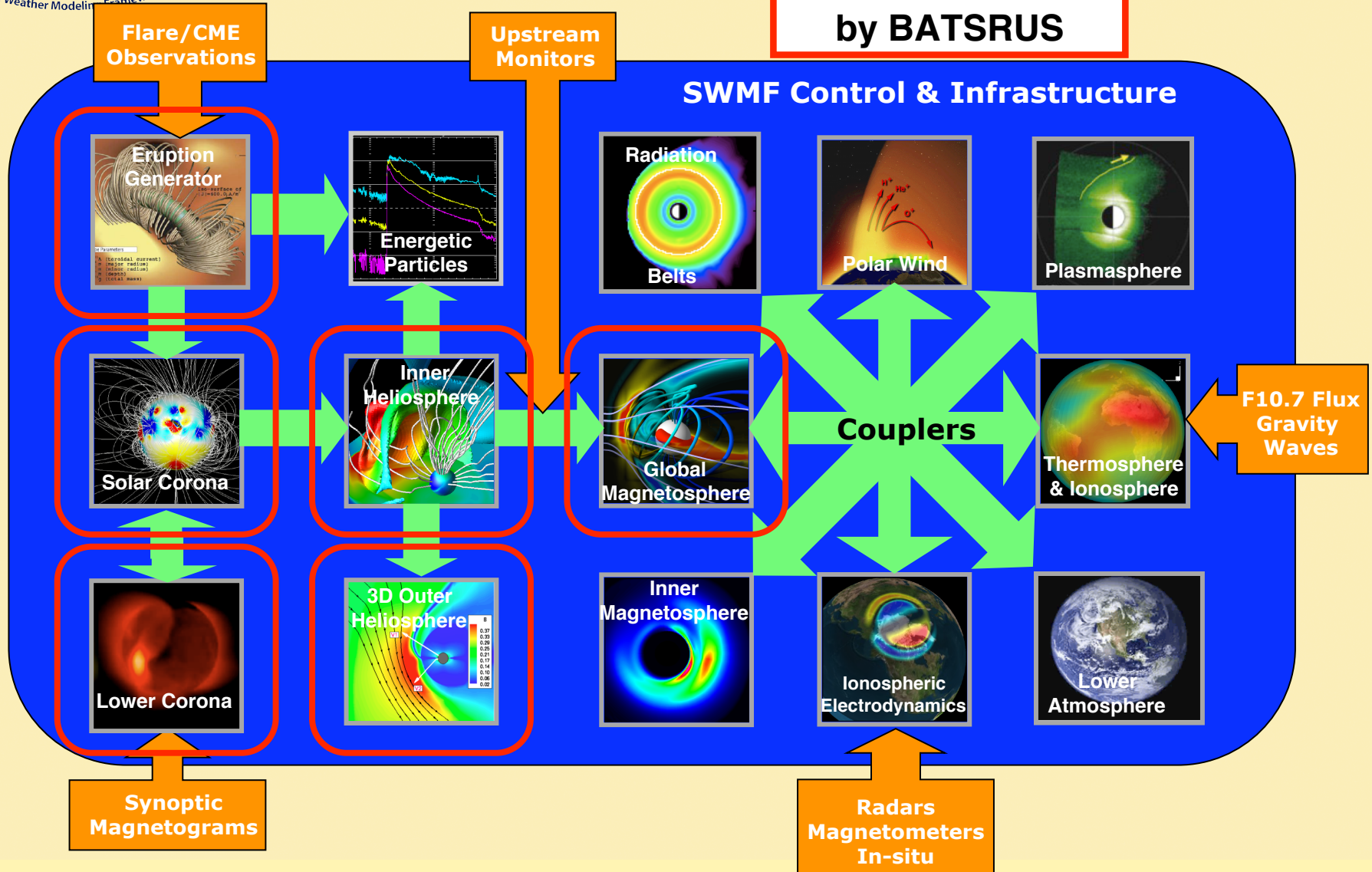
- 🌐 Solar wind application

## **M** Summary

# Space Weather Modeling Framework



domains modeled  
 by **BATSRUS**



SWMF is freely available at <http://csem.engin.umich.edu> and via CCMC

# BATS-R-US

## Block Adaptive Tree Solar-wind Roe Upwind Scheme



### Physics

- Classical, semi-relativistic and Hall MHD
- Multi-species, multi-fluid, **anisotropic ion pressure**
- **(Anisotropic) heat conduction, Alfvén wave turbulence**
- Radiation hydrodynamics with grey/multigroup diffusion
- Multi-material, non-ideal equation of state

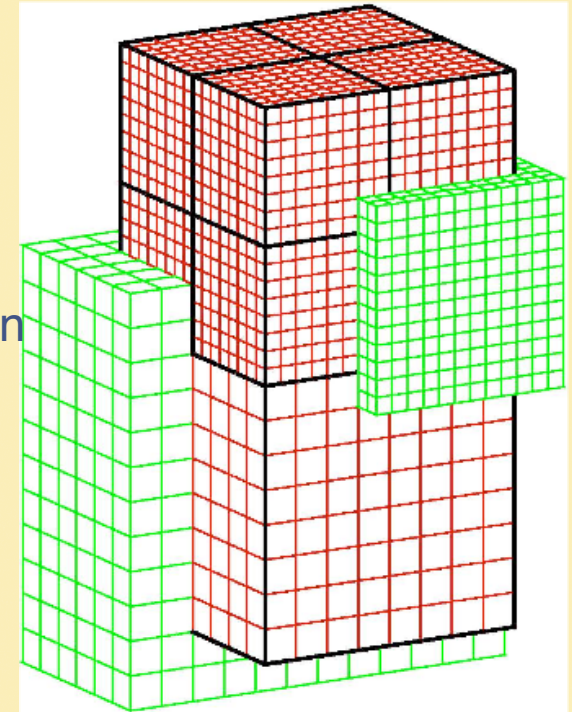
### Numerics

- Conservative finite-volume discretization
- Parallel block-adaptive grid
- Cartesian and generalized coordinates
- Splitting the magnetic field into  $B_0 + B_1$
- Divergence B control: 8-wave, CT, projection, parabolic/hyperbolic cleaning
- Shock-capturing TVD schemes: Rusanov, HLLE, AW, Roe, HLLD
- Explicit, point-implicit, semi-implicit, fully implicit time stepping

### Applications

- Sun, heliosphere, magnetospheres, unmagnetized planets, moons, comets...

**100,000+ lines of Fortran 90 code with MPI parallelization**



## Anisotropic MHD



**M Different pressures parallel and perpendicular to the magnetic field**

**M Space physics applications**

- Reconnection
- Magnetosphere
- Coupling with inner magnetosphere models
- Solar wind heating

**M Difficulties**

- What is a proper conservative form?
- Physical instabilities: fire-hose, mirror, proton cyclotron

**M Combinations with more physics**

- Separate electron pressure
- Hall MHD, semi-relativistic, multi-ion

# Resistive MHD with electrons and anisotropic ion pressure



Mass conservation:  $\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0$

Momentum:  $\frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot (\rho \mathbf{u} \mathbf{u} + P) = \mathbf{J} \times \mathbf{B}$

$$P = (p_{\perp} + p_e)I + (p_{\parallel} - p_{\perp})\mathbf{b}\mathbf{b} \quad p = \frac{2p_{\perp} + p_{\parallel}}{3}$$

Induction:  $\frac{\partial \mathbf{B}}{\partial t} + \nabla \times \mathbf{E} = 0$

Pressure:  $\frac{\partial p_{\perp}}{\partial t} + \nabla \cdot (p_{\perp} \mathbf{u}) = \frac{1}{3\tau} (p_{\parallel} - p_{\perp}) + \frac{2}{\tau} (p_e - p) - p_{\perp} \nabla \cdot \mathbf{u} + p_{\perp} \mathbf{b} \cdot (\nabla \mathbf{u}) \cdot \mathbf{b}$

$$\frac{\partial p_{\parallel}}{\partial t} + \nabla \cdot (p_{\parallel} \mathbf{u}) = \frac{2}{3\tau} (p_{\perp} - p_{\parallel}) + \frac{2}{\tau} (p_e - p) - 2p_{\parallel} \mathbf{b} \cdot (\nabla \mathbf{u}) \cdot \mathbf{b}$$

Electron pressure:

$$\tau_{ie} = \frac{2}{3} \frac{M_i}{\eta e^2 n_e}$$

$$\frac{\partial p_e}{\partial t} + \nabla \cdot (p_e \mathbf{u}) = (\gamma - 1) [-p_e \nabla \cdot \mathbf{u} + \eta \mathbf{J}^2 + \nabla \cdot (\kappa \mathbf{b}\mathbf{b} \cdot \nabla T_e)] + \frac{2}{\tau_{ie}} (p - p_e)$$

Electric field:  $\mathbf{E} = -\mathbf{u} \times \mathbf{B} + \eta \mathbf{J}$

$$\mathbf{b} = \mathbf{B}/B$$

Current:  $\mathbf{J} = \nabla \times \mathbf{B}$

## Conservative Form?



- ⌘ Shock capturing schemes require conservation laws to get proper jump conditions
  - 🌐 Double adiabatic invariants are not the right conservative variables
  - 🌐 Energy conservation only replaces one of the two pressure equations
  - 🌐 The anisotropy behind a shock is determined by instabilities
- ⌘ We use the energy equation and instability criteria to get proper jump conditions.

## M Instabilities

🌐 Fire-hose:  $\frac{p_{\parallel}}{p_{\perp}} > 1 + \frac{B^2}{p_{\perp}}$  (destabilized Alfvén wave)

🌐 Mirror:  $\frac{p_{\perp}}{p_{\parallel}} > 1 + \frac{B^2}{2p_{\perp}}$

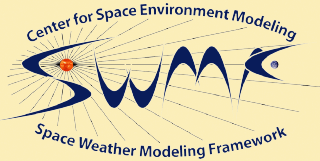
🌐 Proton cyclotron:  $\frac{p_{\perp}}{p_{\parallel}} > 1 + 0.847 \left( \frac{B^2}{2p_{\parallel}} \right)^{0.48}$

🌐 In unstable regions we reduce anisotropy so it becomes stable

## M Ion-ion, ion-electron and/or wave-ion interactions:

🌐 Push ion pressure towards isotropic distribution with time rate  $\tau$



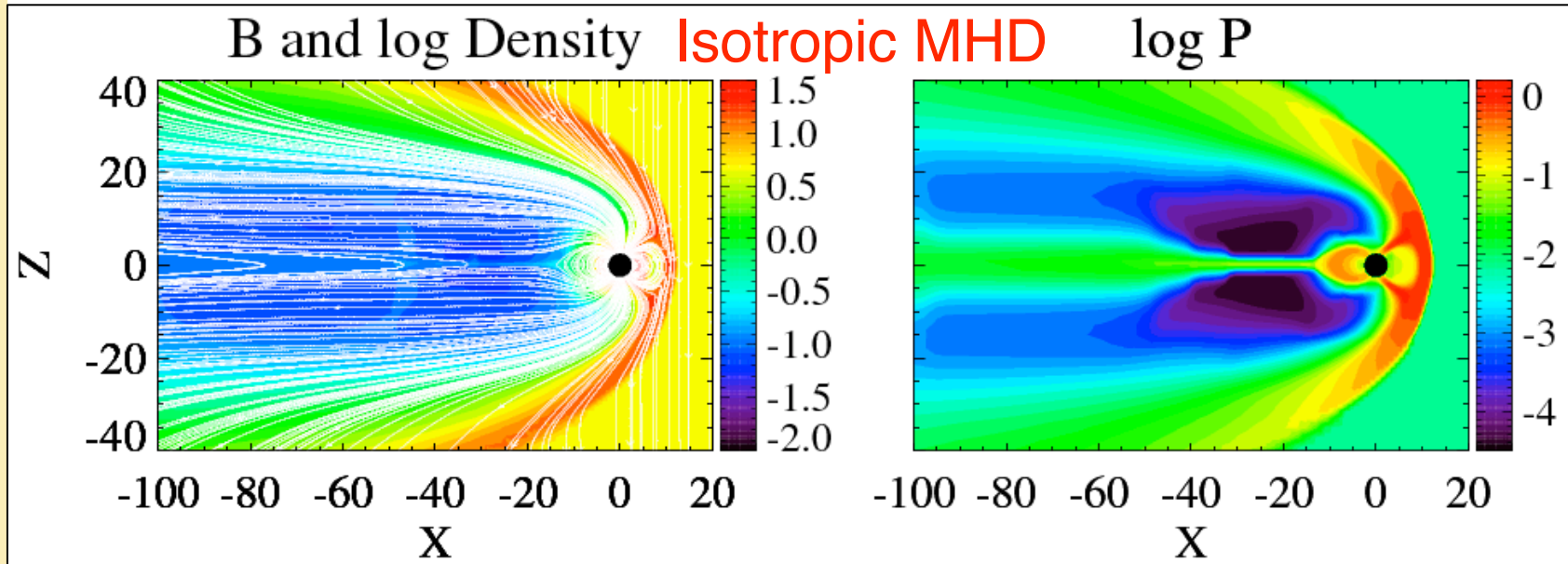
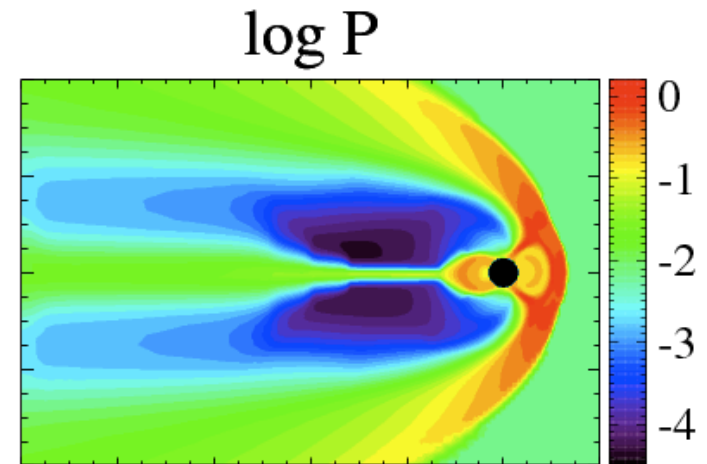
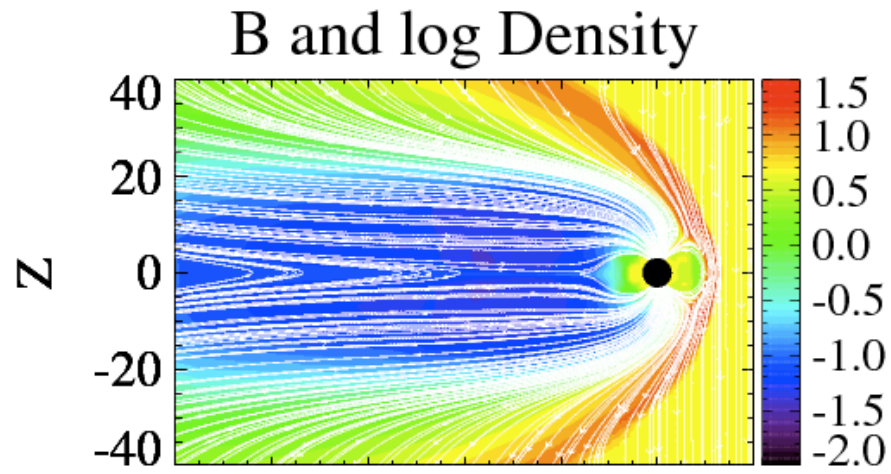


## Idealized Steady Magnetosphere Run (Xing Meng)



- dipole axis aligned with Z
- Steady solar wind:  $n = 5 \text{ /cc}$ ,  $v = 400 \text{ km/s}$ ,  $B_z = -5 \text{ nT}$
- Solve for energy and parallel pressure near bow shock
- Enforce stability conditions
- Relaxation rate towards isotropy:  $\tau = 20 \text{ s}$

# Idealized Steady Magnetosphere Run (Xing Meng)

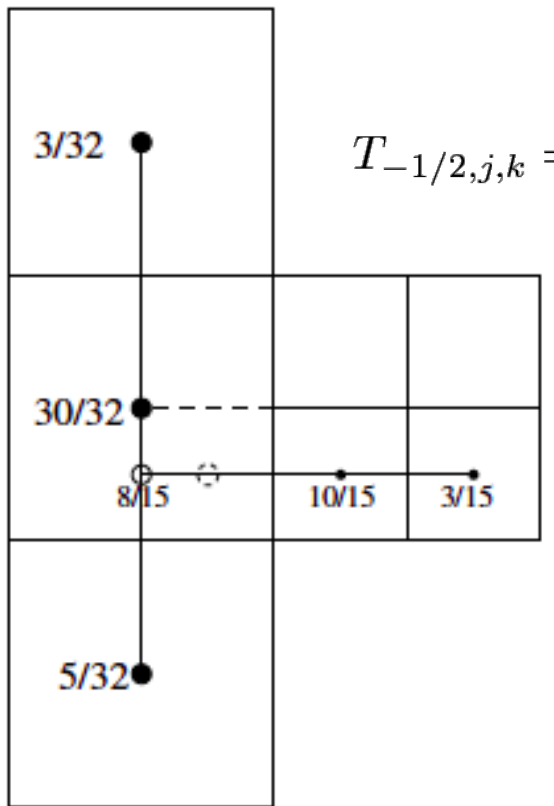


## Electron thermal heat conduction



$$C_V \frac{\partial T_e}{\partial t} = -\nabla \cdot \left[ \kappa_e T_e^{5/2} \frac{\mathbf{B}\mathbf{B}}{B^2} \cdot \nabla T_e \right]$$

- At resolution changes: interpolation of  $T_e$  at fine AMR block has to be third order to make the scheme second order. Use finite difference approach.



$$T_{-1/2,j,k} = \frac{5T_{-1/2,j-3/2,k-3/2} + 30T_{-1/2,j+1/2,k+1/2} - 3T_{-1/2,j+5/2,k+5/2}}{32}$$

$$T_{0,j,k} = \frac{8T_{-1/2,j,k} + 10T_{1,j,k} - 3T_{2,j,k}}{15}$$

- At block faces:

- Similar interpolation schemes are used for the block edges

- M Collisional electron heat conduction from inner coronal boundary to approximately  $5R_{\text{sun}}$ , smoothly diminishes between  $5R_{\text{sun}}$  and  $10R_{\text{sun}}$**
- M Heating of protons by Alfvén wave dissipation**
- M Heating of electrons by collisional coupling with protons**

$$\frac{\partial p_i}{\partial t} + \nabla \cdot (p_i \mathbf{u}) + (\gamma - 1)p_i \nabla \cdot \mathbf{u} = (\gamma - 1) [ Q_i + \lambda_{ei}(T_e - T_i) ],$$

$$\frac{\partial p_e}{\partial t} + \nabla \cdot (p_e \mathbf{u}) + (\gamma - 1)p_e \nabla \cdot \mathbf{u} = (\gamma - 1) [ -\nabla \cdot \mathbf{q}_e + \lambda_{ei}(T_i - T_e) ],$$

$$\mathbf{q}_e = -\kappa_e T_e^{5/2} \frac{\mathbf{B}\mathbf{B}}{B^2} \cdot \nabla T_e$$

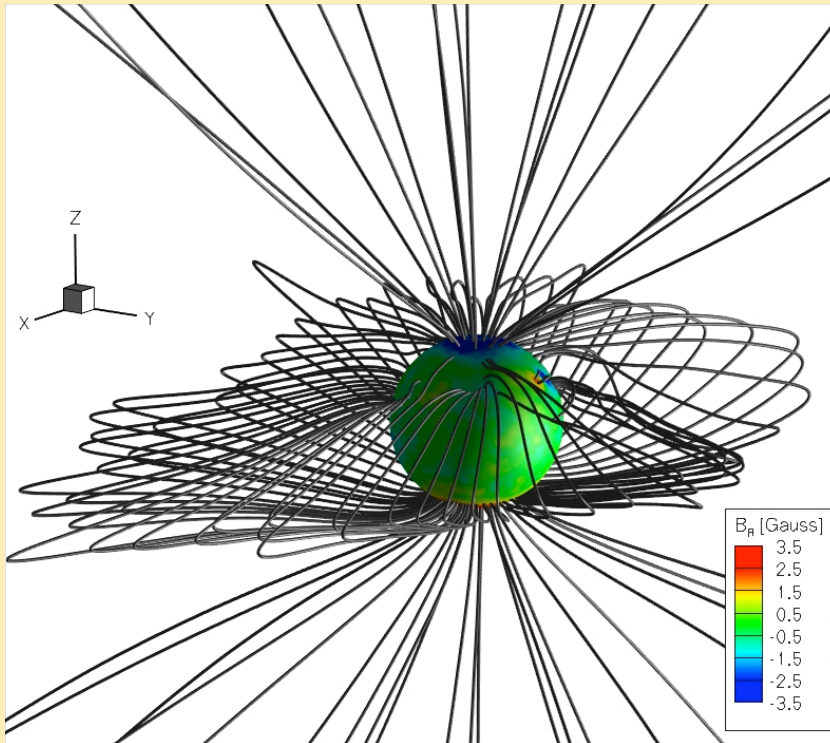
- M Wind acceleration: work done by wave pressure force**
- M Coronal heating: formulation of the Kolmogorov dissipation by Hollweg (1986)**

$$\frac{\partial E_w^+}{\partial t} + \nabla \cdot [E_w^+ (\mathbf{u} + \mathbf{u}_A)] + p_w^+ \nabla \cdot \mathbf{u} = -Q^+,$$

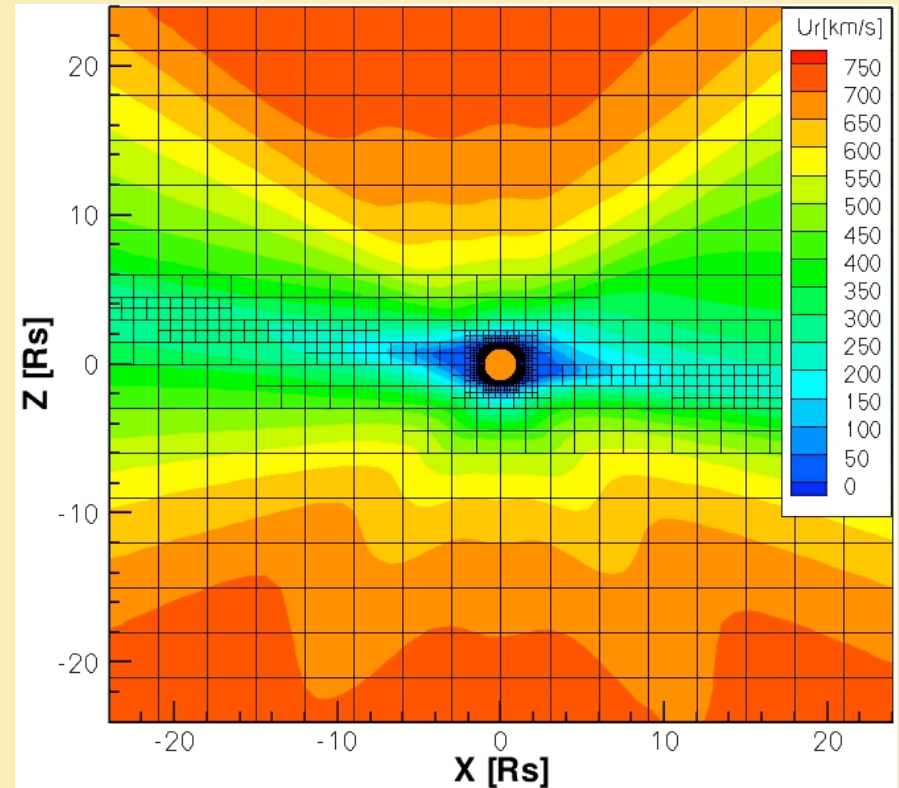
$$\frac{\partial E_w^-}{\partial t} + \nabla \cdot [E_w^- (\mathbf{u} - \mathbf{u}_A)] + p_w^- \nabla \cdot \mathbf{u} = -Q^-,$$

$$\text{Ion heating } Q_i = Q^+ + Q^- = \frac{E_w^{+3/2}}{L\sqrt{\rho}} + \frac{E_w^{-3/2}}{L\sqrt{\rho}}, \quad L = C/\sqrt{B},$$

- M Free parameter C in heating scale height L**

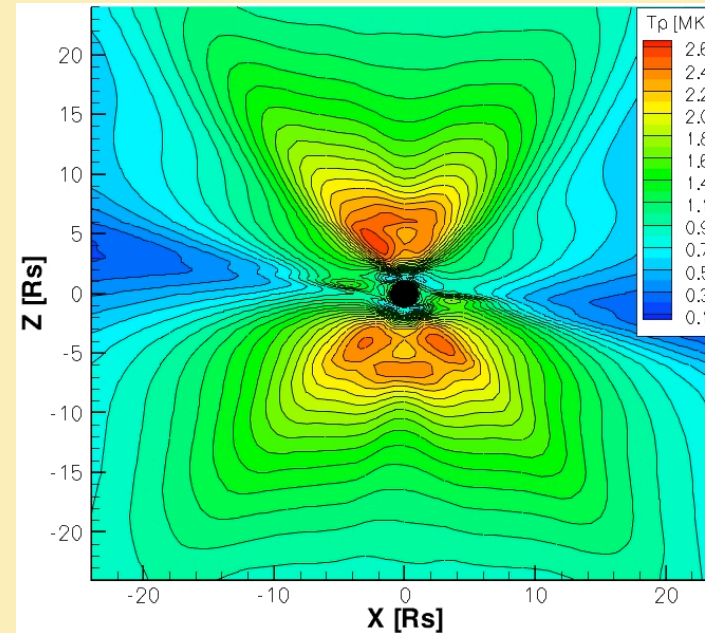
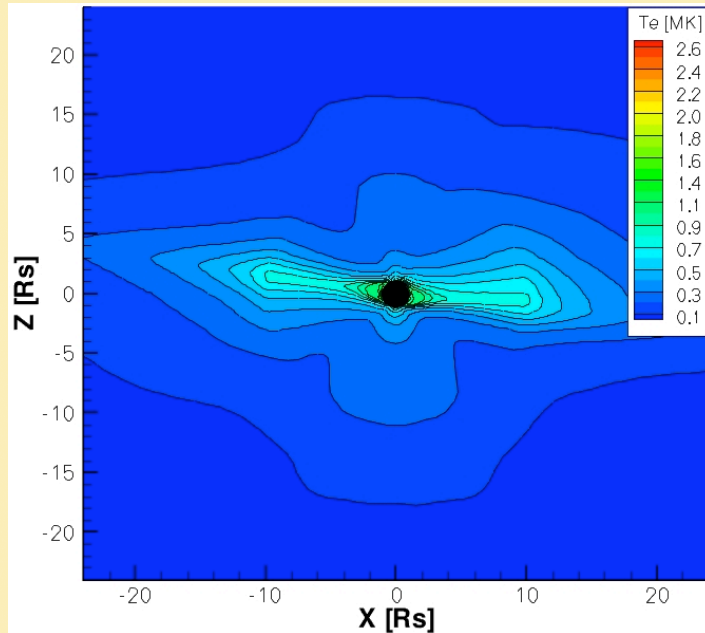


**Selected field lines showing streamer belt**



**Meridional slice showing bimodal wind due to Alfvén waves**

## December 2008 Solar Wind

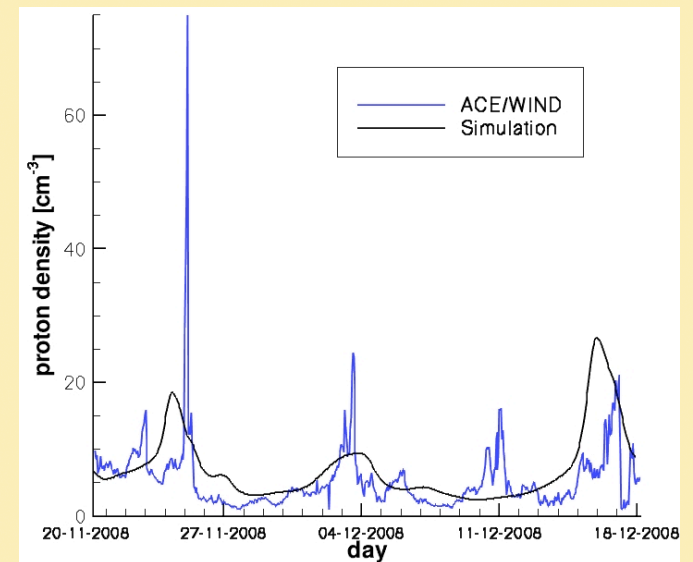
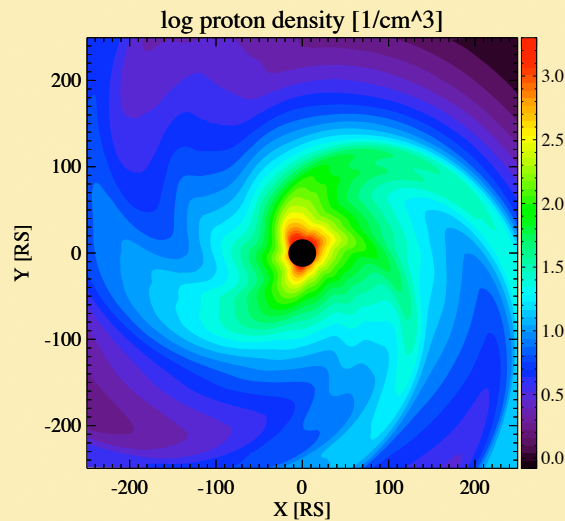
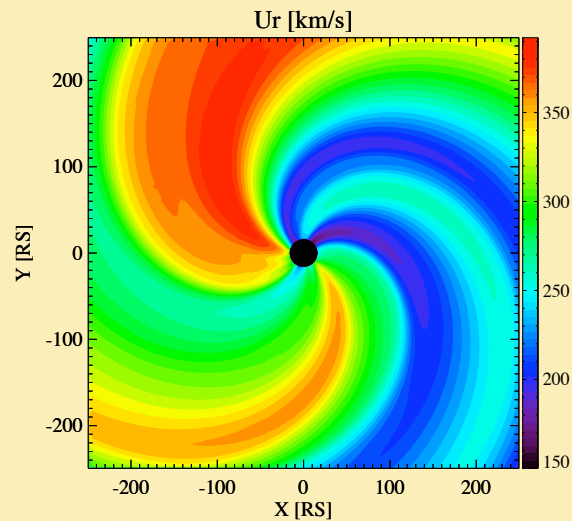


- M** High electron temperature above streamer due to heat conduction, cool electrons in fast wind due to adiabatic expansion
- M** Protons mostly heated in coronal hole due to Alfvén waves

# Corotating Interaction Regions in Inner Heliosphere



- Non-alignment of magnetic- and rotation-axis gives appearance of fast and slow streams
- Result: compression of plasma seen as spiral arms
- Comparison with ACE satellite at L1 point





## **M Anisotropic proton pressures in BATSRUS**

- 🌐 Enforcing stability limits
- 🌐 Optional (ad hoc) relaxation term
- 🌐 Comparison of magnetosphere simulations with observations in progress

## **M Electron Physics in BATSRUS**

- 🌐 Separate electron pressure and electron thermal heat conduction
- 🌐 Comparison of solar wind at 1AU with observations is rather good

## **M Plans**

- 🌐 Combined electron pressure and anisotropic proton pressure in both the solar wind and magnetosphere
- 🌐 Include counter-propagating Alfvén waves (partial reflection due to inhomogeneities in background)