

Probing dark matter couplings to top and bottom at the LHC

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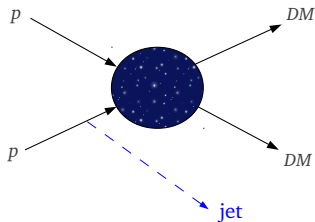
with Liantao Wang and Rocky Kolb, arXiv:1303.6638

April 17, 2013

Dark matter at colliders

“Mono” searches

- **Monojets:** DM production (E_T^{miss}) plus one or two jets
- Can arise from initial state radiation of quarks, gluons

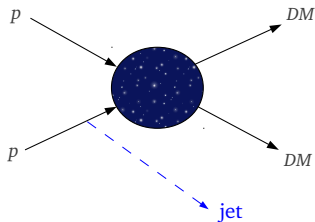


- Useful collider probe for light dark matter, and operators with SD scattering
- Related: monophoton, mono- Z , mono-leptons, **mono- b**

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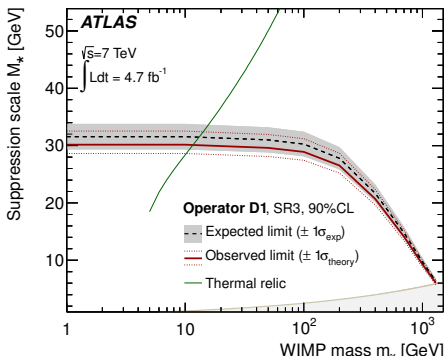
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Constraints on operators

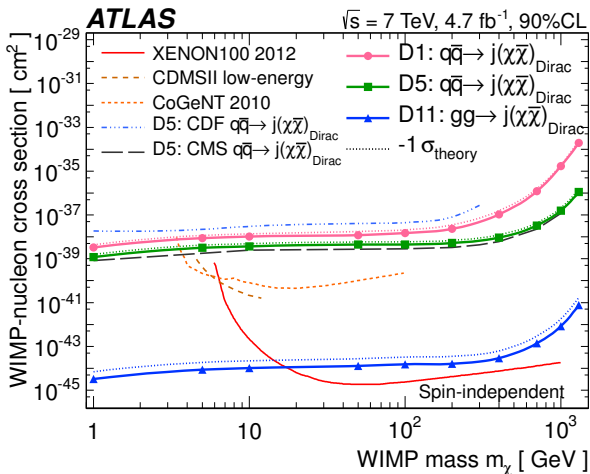
Model-independent effective field theory (EFT) approach: simple connection to direct detection and relic abundance (annihilation cross section)

Single scale M_* parameterizes size of interactions:

$$\text{D1/Scalar} : \mathcal{O} = \frac{m_q}{M_*^3} \bar{q}q \bar{X}X$$

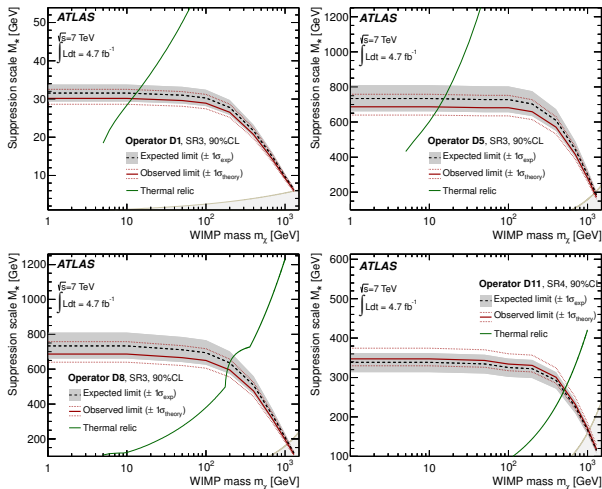


Map to direct detection



Constraints on operators

Constraints on M_* much worse for scalar operator:



Scalar operator

A closer look at scalar operators:

$$\mathcal{O}_S \equiv \frac{m_q}{M_*^3} \bar{q}q \bar{X}X \sim \frac{y_q g_X}{m_\phi^2} \bar{q}q \bar{X}X$$

- Why m_q ? Flavor constraints \rightarrow Minimal Flavor Violation
- Suppressed couplings to light quarks means suppressed monojet signal from gluons and light quarks
- ATLAS published result (2012) - only consider up to m_c

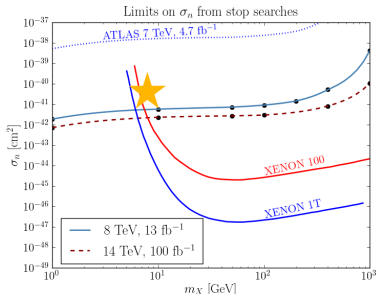
This operator gives rise to spin-independent scattering:

$$\sigma_n = \frac{(0.38 m_n)^2 \mu_X^2}{\pi M_*^6} \approx 2 \times 10^{-38} \text{cm}^2 \left(\frac{30 \text{ GeV}}{M_*} \right)^6$$

(Other operators like $\bar{q}\gamma^5 q \bar{X}\gamma^5 X$ give similar collider constraints but velocity-suppressed direct detection cross sections.)

Searches with bottoms and tops

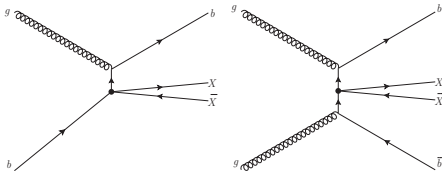
- Include heavy quarks in monojet signal
- Significant improvement in constraints on scalar interactions with b -tagging and $t\bar{t}$ final states
 - heavy flavor production: y_q^2 wins over pdf suppression
 - b -enriched final states from bottoms and tops
- Direct search for couplings to third generation



DM production

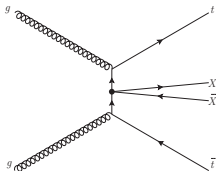
Direct b production:

- $bg \rightarrow bX\bar{X}$, $gg \rightarrow b\bar{b}X\bar{X}$



DM production in association with tops:

- $bg \rightarrow bX\bar{X}$, $gg \rightarrow b\bar{b}X\bar{X}$



Factor of ~ 10 larger cross section than direct b production.

b -tagging on monojets

Monojet:

- $E_T^{miss} > 350\text{GeV}$, $p_T^1 > 100\text{ GeV}$, No lepton
- No more than 2 jets with $p_T > 50\text{ GeV}$

Mono- b :

- **b -tag on leading jet**

b -tagging efficiency (calibrated on tops) $\sim 60\%$, mistag $\sim 0.2\%$ for light quarks at ATLAS

Dominant SM backgrounds for monojet searches: Z +jets (W + jets) cut from 500 fb to around 15 fb with b -tag.

Mono- b

	Process	Monojet	b -tag	b -tag on j_1
Signal	$\bar{X}X + \text{jets}$	11 fb	0.9 fb	0.7 fb
	$\bar{X}X + b + \text{jets}$	65 fb	40 fb	33 fb
	$\bar{X}X + t\bar{t}$	120 fb	63 fb	41 fb

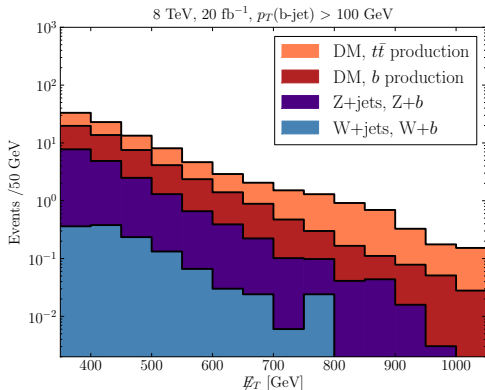
Table: **Inclusive mono- b search at 8 TeV:** Cross sections after cuts applied. For the signal we take $M^* = 50$ GeV.

Improvement in bounds from: (1) overall monojet production rate is bigger (2) monojets are b -enriched

Simulation of Events:

- MadGraph5 + PYTHIA6 + Delphes
- NLO cross sections from MCFM-Dark (Fox and Williams)

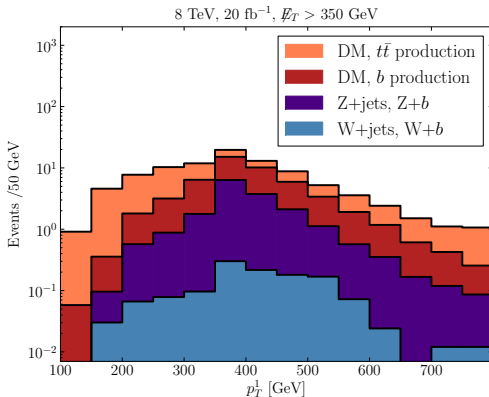
Mono- b MET distribution



$$(M_* = 50 \text{ GeV})$$

- Direct b -production isn't significantly harder
- Z +jets often initiated by at least one valence quark
- $t\bar{t}$ has different kinematics and a harder spectrum; it isn't coming from ISR

Mono- b p_T distribution

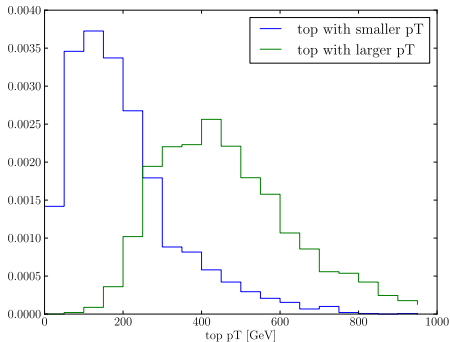


($M_* = 50$ GeV)

- For b production, typically just one hard b -jet
- For top production, second jet may be harder, b -tagged 30% of the time

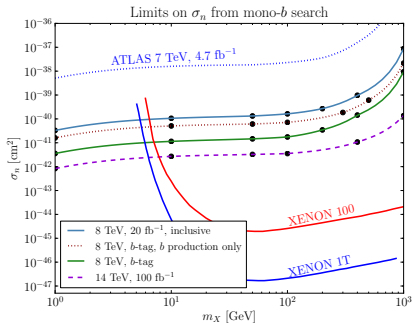
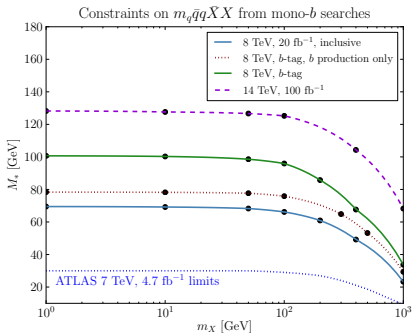
Top kinematics: events passing mono- b cuts

Associated production of $t\bar{t}$ can look like “mono- b ” cuts if one top is boosted and the other is not too hard.



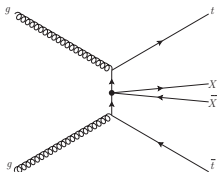
Efficiency WRT veto on 3rd jet: 0.17

Mono-b limits



$t\bar{t}$ final state

A direct search for $t\bar{t} + \text{MET}$ would do better in this case:

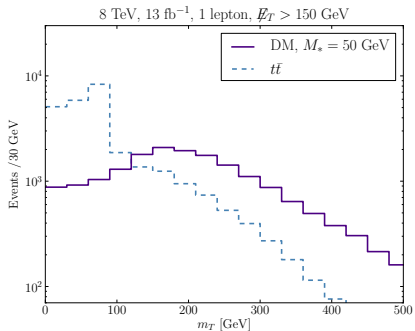
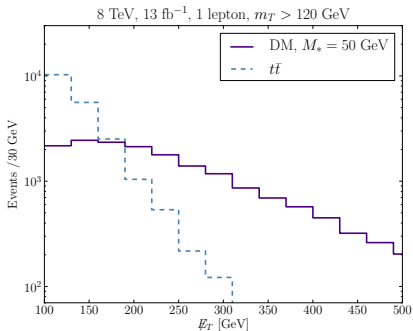


→ SUSY top-partner searches

We apply the ATLAS 8 TeV, 13 fb^{-1} study with 1 lepton in the final state:

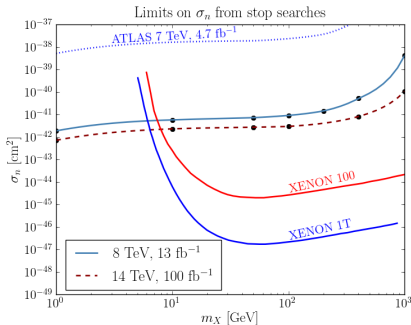
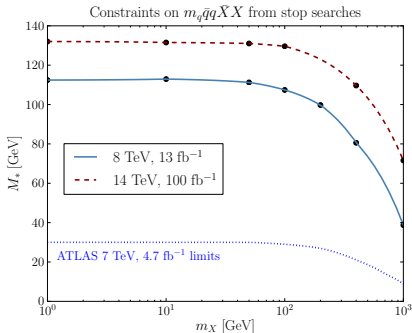
- 4 jets with $p_T > (80, 60, 40, 25)$ GeV with 1 b -tag
- We use signal region with $E_T^{\text{miss}} > 225$ GeV, $m_T > 130$ GeV

$t\bar{t}$ final state



Dominant SM background is $t\bar{t}$ dileptonic decay, where one lepton is either missed or a τ .

Limits from recasting stop search

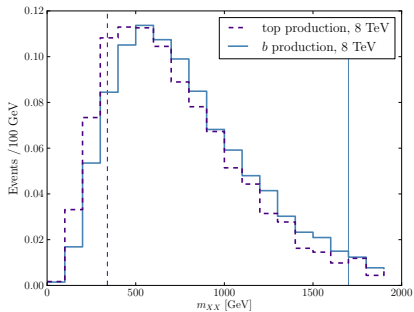
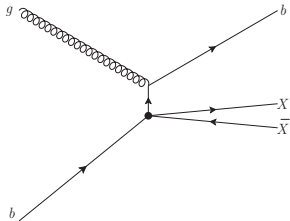


~ 20% improvement in M_* constraint compared to mono- b .
Another 10% possible with all-hadronic final state.

Unitarity

Bounds on M_* still low - what do these constraints mean?
(Shoemaker and Vecchi; Fox, Harnik, Primulando, Yu):

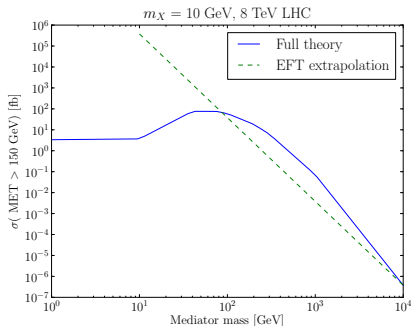
$$\frac{\sqrt{3}}{8\pi} \frac{m_q}{M_*^3} m_{\chi\bar{\chi}}^2 \lesssim \frac{1}{2}$$



b -production seems safe, for top production it depends...

Example: neutral scalar mediator

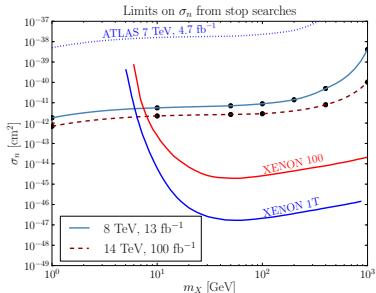
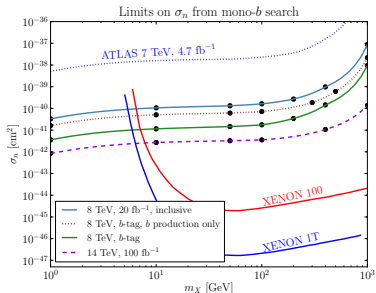
- Interactions through a mediator of mass M_ϕ can go through a resonance:



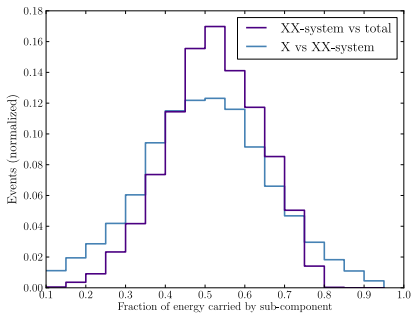
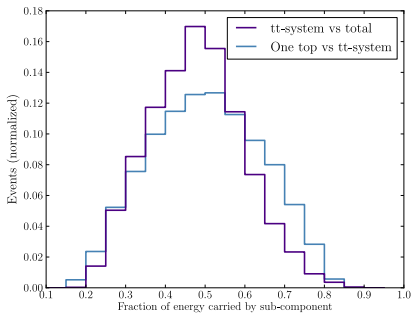
- Other constraints can also apply depending on the model.
- If the scalar is the Higgs, the constraints on Higgs invisible decay - $Br(H \rightarrow \chi\chi)$ - are already stronger.

Conclusions

- Monojets searches useful to study DM-SM interactions
- For scalar operators, final states are b enriched. “Mono- b ” gives significant improvement in bounds.
- Top production even more important: direct $t\bar{t} + E_T^{miss}$ search
- Direct probe of couplings to third generation

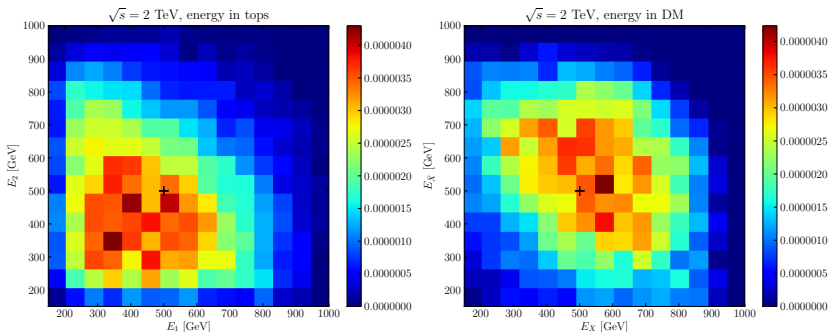


Top kinematics



Top kinematics

Compared to $t\bar{t}$ or stop pair production, in this case each particle carries off about $\sim 1/4$ of the energy in roughly uncorrelated directions \rightarrow large E_T^{miss} .



2D parton level distributions for energies of tops (left) and DM (right)

Loops

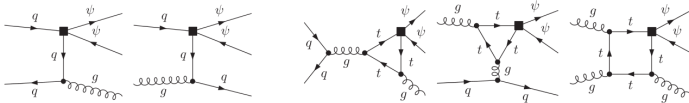
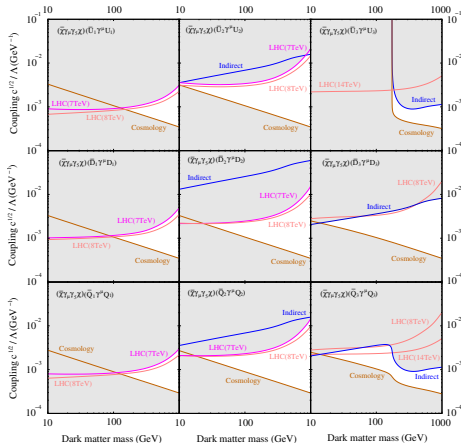


FIG. 1. Typical tree-level (left) and loop-level (right) diagrams leading to monojet events. The black squares denote insertions of four-fermion operators.

Haisch, Kahlhoefer, Unwin: M_* limit raised to 150 GeV.

Scalar operator constraints

- Recent paper by Bhattacharjee et al. [1212.5013], “Model Independent Analysis of Interactions between Dark Matter and Various Quarks”:



- Kamenik and Zupan 2011 - “monotop” from MFV and dark matter