Skirmishes on the LDM frontier

territorial disputes over the low-recoil region in direct detection

> Josef Pradler Johns Hopkins University

"Light Dark Matter" University of Michigan April 15, 2013

Outline

- I. New feeble signals at the direct DM detection threshold "Claiming territory over the low recoil region"
 - Dark Photons => Haipeng An's talk this morning [304.3461 (PLB)]

1302.3884 with Haipeng An and Maxim Pospelov

- New neutrino signals in direct detection

I203.0545 (PRD) with Maxim Pospelov

2. A critical look at DAMA's Dark Matter claim

I210.5501 (PLB) with Balraj Singh and Itay YavinI210.7548 with Itay Yavin

A vision of a true neutrino observatory

PHYSICAL REVIEW D

VOLUME 30, NUMBER 11

1 DECEMBER 1984

Principles and applications of a neutral-current detector for neutrino physics and astronomy

A. Drukier and L. Stodolsky Max-Planck-Institut für Physik und Astrophysik, Werner-Heisenberg-Institut für Physik, Munich, Federal Republic of Germany (Received 21 November 1983)

- superconducting grains in filler material in magnetic field
- at low temperatures specific heat ~ T^3

=> single scatter of neutrino can make grain conducting

=> magnetic field collapses, induces electric signal in detector

coherent neutrinonucleus scattering

$$\frac{d\sigma}{d\cos\theta} = \frac{1}{8\pi} G_F^2 E_\nu^2 \left[Z(4\sin^2\theta_W - 1) + N \right]^2 \left(1 + \cos\theta \right)$$

• coherent enhancement N^2 for MeV-scale neutrinos from

=> spallation sources, supernovae, reactors, sun, earth

- cross section grows quadratically with neutrino energy
- helicity conservation forbids back-scattering

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(this process has not yet been observed)

=> direct DM detection

PHYSICAL REVIEW D

VOLUME 31, NUMBER 12

15 JUNE 1985

Detectability of certain dark-matter candidates

Mark W. Goodman and Edward Witten Joseph Henry Laboratories, Princeton University, Princeton, New Jersey 08544 (Received 7 January 1985)

We consider the possibility that the neutral-current neutrino detector recently proposed by **Drukier and Stodolsky** could be used to detect some possible candidates for the dark matter in galactic halos. This may be feasible if the galactic halos are made of particles with coherent weak interactions and masses $1-10^6$ GeV; particles with spin-dependent interactions of typical weak strength and masses $1-10^2$ GeV; or strongly interacting particles of masses $1-10^{13}$ GeV.





Witten 1985

=> direct DM detection

• nuclear recoil can be picked up in various channels:





scintillation



ionization

WIMPs vs. solar neutrinos

flux $\Phi_{pp} = 6 \times 10^{10} \,\mathrm{cm}^{-2} \mathrm{s}^{-1}$ $\Phi_{DM} = \frac{\rho_0 v}{m_{DM}} \sim 10^5 \,\mathrm{cm}^{-2} \mathrm{s}^{-1} \left(\frac{100 \,\mathrm{GeV}}{m_{DM}}\right)$ $\Phi_{^{8}\mathrm{B}} = 6 \times 10^{6} \, \mathrm{cm}^{-2} \mathrm{s}^{-1}$ cross section $\sigma \simeq 10^{-44} \,\mathrm{cm}^2 \times N^2 \left(\frac{E_\nu}{1 \,\mathrm{MeV}}\right)^2$ $\sigma = 10^{-44} \,\mathrm{cm}^2 \times \sigma_{44} A^2 \left(\frac{\mu_N}{\mu_m}\right)^2$ recoil $E_R^{\max} = \frac{(2\mu_N v)^2}{2m_N} \sim \begin{cases} 20 \,\text{keV}\left(\frac{A}{20}\right) & E_R^{\max} = \frac{(2E_\nu)^2}{2m_N} \\ \left(m_N \ll m_{DM}\right) & e_R^{\max} = \frac{(2E_\nu)^2}{2m_N} \\ 4 \,\text{keV}\left(\frac{m_{DM}}{20 \,\text{GeV}}\right)^2 \left(\frac{100}{A}\right) & \sim 0.1 \,\text{keV}\left(\frac{20}{A}\right) \left(\frac{E_\nu}{1 \,\text{MeV}}\right)^2 \\ \left(m_{DM} \ll m_N\right) & e_R^{\max} = \frac{(2E_\nu)^2}{2m_N} \end{cases}$

"baryonic" neutrinos ν_b

M. Pospelov PRD 2011

- introduce new left-handed neutrino species ν_b together with gauged $U(1)_b$
- ν_b couples to quarks, but not to leptons
- breaking of $U(1)_b$ gives new gauge field V_μ mass

$$\mathcal{L}_B = \overline{\nu}_b \gamma^{\mu} (i\partial_{\mu} - g_l q_b V_{\mu}) \nu_b - \frac{1}{3} g_b \sum_q \bar{q} \gamma^{\mu} q V_{\mu} - \frac{1}{4} V_{\mu\nu} V^{\mu\nu} + \frac{1}{2} m_V^2 V_{\mu} V^{\mu} + \mathcal{L}_m.$$

$$\nu_b \checkmark \text{ sterile under SM-gauge group} \\ \text{active under } U(1)_b$$

"baryonic" neutrinos ν_b

M. Pospelov PRD 2011

• for $Q^2 \ll m_V^2$ effective Lagrangian reads

$$\mathcal{L}_{\text{eff}} = -G_B j_{NCB}^{\mu} \sum_{N=n,p} \overline{N} \gamma_{\mu} N, \qquad G_B = q_b \frac{g_b g_l}{m_V^2}$$

$$j^{\mu}_{NCB} = \overline{\nu}_b \gamma^{\mu} \nu_b$$

• measure interaction strength in units of G_F :

$$\mathcal{N} = \frac{|G_B|}{G_F} \simeq 100 \times \left(\frac{3\,\text{GeV}}{m_V}\right)^2 \left(\frac{g_l g_b}{10^{-2}}\right)$$

"baryonic" neutrinos ν_b

M. Pospelov PRD 2011

• crucial insight:

$$\frac{\sigma_{\nu_b N}(\text{elastic})}{\sigma_{\nu_b N}(\text{inelastic})} \sim \frac{A^2}{E_{\nu}^4 R_N^4} \sim \mathcal{O}(10^8)$$

this ratio makes direct detection experiments competitive with large scale neutrino experiments

For solar flux, deuteron breakup in SNO does not constrain scenario

direct detection of solar ν_b

like SM-neutrinos with $G_F^2(N/2)^2 \rightarrow G_B^2 A^2$ $\frac{dR(t)}{dE_R} = N_T \left[\frac{L_0}{L(t)}\right]^2 \sum_i \Phi_i \int_{E_{\text{trin}}} dE_\nu \, \frac{df_i}{dE_\nu} \frac{d\sigma}{dE_R} P_b(t, E_\nu)$ 1 1 appearance probability overall flux average over modulation neutrino spectrum i $L_0 = 1 \,\mathrm{AU}$ $L(t) = L_0 \left\{ 1 - \epsilon \cos \left| \frac{2\pi(t - t_0)}{1 \operatorname{vr}} \right| \right\}$ $t_0 \simeq 3 \text{ Jan} \text{ (perihelion)}$ $\epsilon = 0.0167$ (eccentricity)

direct detection of ν_b

$$\frac{dR(t)}{dE_R} = N_T \left[\frac{L_0}{L(t)}\right]^2 \sum_i \Phi_i \int_{E_\nu^{\min}} dE_\nu \frac{df_i}{dE_\nu} \frac{d\sigma}{dE_R} P_b(t, E_\nu)$$

$$\uparrow$$
more modulation here

$$\frac{L_{\rm osc}}{L_0} \simeq 0.5 \times \left(\frac{10^{-10} \,\text{eV}}{\Delta m^2}\right) \left(\frac{E_{\nu}}{10 \,\text{MeV}}\right)$$

oscillation-length on the order sun-earth distance

=> flip phase for high energy part of the neutrino spectrum? explain DAMA?

Appearance probability

• considering small values in Δm_b^2 standard solar story unfolds

$$P_b(\text{earth}) \simeq \sin^2(2\theta_b) \sin^2\left[\frac{\Delta m_b^2 L(t)}{4E}\right]$$
$$\bigvee$$
$$\mathcal{N}_{\text{eff}}^2 \equiv \frac{\mathcal{N}^2}{2} \times \sin^2 2\theta_b$$

(from a tribimaximal ansatz assuming mixing to ν_2) [see also arXiv:1103.3261]

=> for fast oscillations $P_b G_B^2 \to \mathcal{N}_{eff}^2 G_F^2$



direct detection of ν_b



14



modulation amplitude



CRESST-II signal

Angloher et al EPJC 2012

- 8 CaWO₄ crystals, measure scintillation light and phonons from nuclear recoil
- in a nutshell: 67 events in acceptance region half of which are attributed to backgrounds



CRESST-II

fits

- we follow CRESST in their modeling of backgrounds
 - => e/gamma events known
 => other bkg. essentially flat



CRESST-II

fits

- we follow CRESST in their modeling of backgrounds
 - => e/gamma events known
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- use Poisson log-Likelihood to fit ν_b

$$\chi_P^2 = 2\sum_i \left[y_i - n_i + n_i \ln\left(\frac{n_i}{y_i}\right) \right]$$

best fit yields

 $\chi_P^2/d.o.f. = 27.8/28$ (recoil spectrum only)



direct detection experiments as neutrino, observatories



reg CDMS-II Si as of today







time series



phase off by one month! :(

Outlook

direct detection





direct detection

• even lighter targets (He) are even better! Neutrinos give more recoil than WIMPs (χ)

$$\frac{E_R^{\max}(\nu)}{E_R^{\max}(\chi)} = \frac{E_\nu^2}{v^2 \mu_\chi^2} \simeq 25 \times \left(\frac{E_\nu}{10 \,\mathrm{MeV}}\right)^2 \left(\frac{4 \,\mathrm{GeV}}{\mu_\chi}\right)^2$$

• directional detection

$$\frac{dR}{d\cos\theta_N}$$
 has strongest "A_{FB}" in direction $\mathbf{v}_{avg}(\chi)$ (WIMPs)
in direction of sun (neutrinos)

Outlook

solar neutrino experiments

• elastic scattering off scintillating mineral oil with ultra-pure setups



Outlook astrophysical signatures

• supernovae production of $\nu_b =>$ signal in direct detection possible?





• may affect dynamics of explosions

insulating scattering sphere

- sensitivity to truly tiny mass splittings over cosmological distances
- Neff = ?

with Liang Dai (JHU)

Part II

apropos DAMA

DAMA signal interpretation in the presence of backgrounds

observed modulation amplitude $S_m \simeq 0.02 \,\mathrm{cpd/kg/keV}$

$$s_m^{\text{obs}} = \frac{S_m}{R} = \frac{S_m}{B + S_0} \simeq 2\%$$
 $S = S_0 + S_m \cos \omega (t - t_0)$

the higher the background, the stronger the signal must be modulated

$$s_m^{\max} \ge s_m^{\text{obs}} \left(1 + \frac{B_0}{S_0} \right) \approx 2\% \times \left(1 + \frac{B_0}{S_0} \right) \qquad s_m^{\max} = S_m / S_0$$

=> take a closer look at the DAMA backgrounds to see what is needed

DAMA signal interpretation in the presence of backgrounds





needs MC to find rate at 3 keV
=> we employ the results from
Kudryavtsev et al 2010



(nuclear recoil $\ll 1 \, \mathrm{keV}$)









angular momentum change by 4 units, "3rd forbidden unique weak decay" => the ONLY such EC realized in nature



(nuclear recoil $\ll 1\,{\rm keV}$)



Simulated DAMA spectrum using reported contaminations



Simulated DAMA spectrum using reported contaminations



 strong indication of a flat background component

 $B_{\rm flat} \simeq 0.85 \ {\rm cpd/kg/keV}$

• β^- and Compton background at low energies are **flat!**

=> work out implication for modulation fraction

required modulation fraction if a flat background is present



challenges standard
 WIMP scenario with
 Maxwellian halo:

 $s_m \lesssim 10\%$

 for I3 ppb potassium contamination

$$s_m \gtrsim 20\%$$

required!



• critique 1: potassium is measured at ${}^{nat}K = 13 \, ppb$

critique 2: EC to g.s. is only 10%
 => our discussion is "captious"

 critique 3: upper limit on signal claimed S₀ ≤ 0.25 cpd/kg/keV
 => allows for 6-10% modulation!



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> => let's check. Requires "slide-forensic". Number not in print.

referenced by DAMA's reply to our paper:



talk by Nozzoli for DAMA, TAUP 2009

Bernabei et al. 1210.6199 1211.6346



none of our questions/ concerns have been addressed.

instead our assumption of a flat background was criticized as being ad hoc

=> their own model is not supported by data

interpretation of the signal in terms of DM is seen to be very sensitive to assumptions on the background....

This is how it's done

Background model for a NaI (Tl) detector devoted to dark matter searches

S. Cebrián^{a,b}, C. Cuesta^{a,b}, J. Amaré^{a,b}, S. Borjabad^b, D. Fortuño^a, E. García^{a,b}, C. Ginestra^{a,b}, H. Gómez^{a,b}, M. Martínez^{a,b}, M.A. Oliván^{a,b}, Y. Ortigoza^{a,b}, A. Ortiz de Solórzano^{a,b}, C. Pobes^{a,b}, J. Puimedón^{a,b}, M.L. Sarsa^{a,b,*}, J.A. Villar^{a,b}

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DAMA should show us the K40 shoulder

A count rate in DAMA much greater than 0.04 cpd/kg/keV will severely undermine a DM interpretation of the signal.

ANAIS collaboration

Conclusions

 new neutrinos with enhanced baryonic currents can be tested in direct detection experiments

=> "DM-like" signals from new neutrino physics can explain DM anomalies CoGeNT and CRESST-II, unchallenged by other searches

=> upcoming experimental results will conclusively probe the most interesting parameter space

 DAMA data speaks against a "vanilla" Dark Matter interpretation
 => a minimum of 20% modulation in any putative signal may at least be required

=> after a decade of modulation maybe it is time to take a more global look at the data set