SUSY, Technicolor, Little Higgs and Inverse problem

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OUTLINE

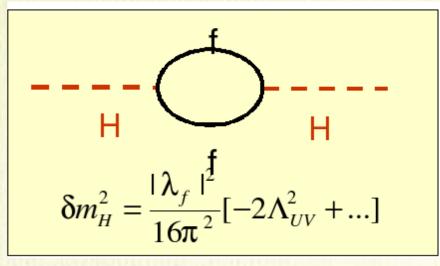
4What is the Inverse problem? Discriminating SM, SUSY and Technicolor through the Higgs boson signatures **+Little Higgs model with T-parity (LHT)** Phenomenology and Signatures +The problem of Discriminating from SUSY +Conclusions

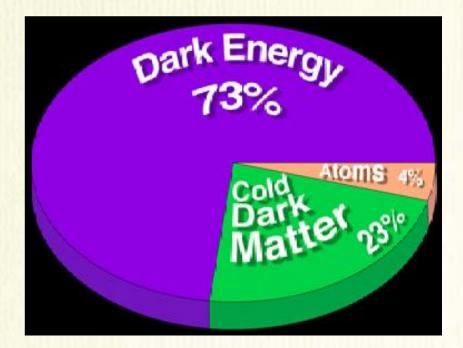
Inverse Problem

- Actually, without specifying the model, there is no inverse problem
- Once your best model (e.g. SUSY) or set of best alternatives (including Little Higgs models, Technicolor, models with extra dimensions, or general model based on the effective Lagrangian approach) is specified, inverse problem comes into play
 Before solving the inverse problem – understanding the underlying theory, we should understand all possible patterns, i.e. do study inverse to inverse

Problems of the Standard Model

- Problem of large quantum corrections to the Higgs mass in SM (Hierarchy problem)
- Coupling Unification
- The origin of Electroweak Symmetry breaking. Higgs boson is not found yet
- The origin and Nature of Dark Matter and Dark Energy No explanation within SM ...
 Baryogenesis problem



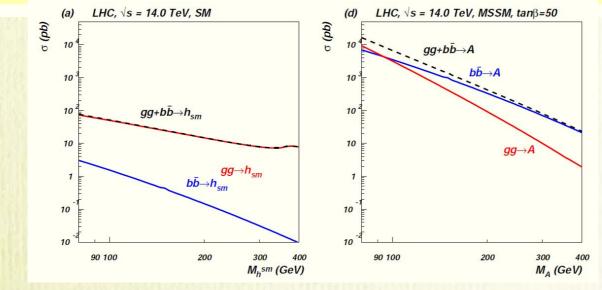


Discriminating the nature of EWSB: SUSY and Technicolor

Related talks given by J. Wells and K. Tobe SUSY pattern

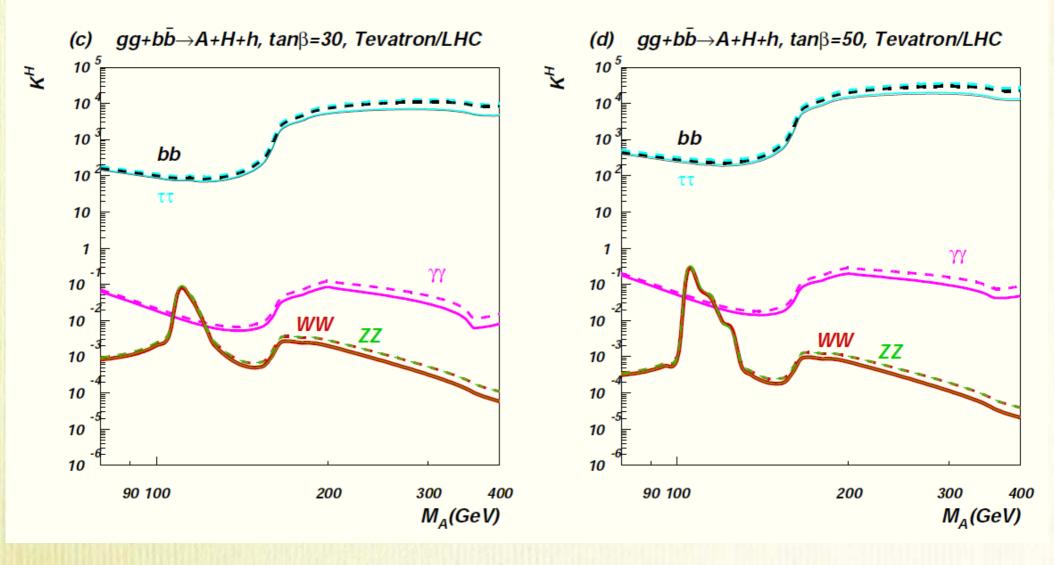
 $\begin{array}{ll} \tan\beta = v_u/v_d \text{ and } M_A & \text{define the Higgs sector at tree level} \\ \text{Large } M_A & \Rightarrow Y_{Hb\bar{b}}/Y_{hb\bar{b}}^{SM} = Y_{H\tau\bar{\tau}}/Y_{h\tau\bar{\tau}}^{SM} \simeq \tan\beta, \\ \text{Small } M_A \simeq M_h & \Rightarrow Y_{hb\bar{b}}/Y_{hb\bar{b}}^{SM} = Y_{h\tau\bar{\tau}}/Y_{h\tau\bar{\tau}}^{SM} \simeq \tan\beta \end{array}$

 $(Y_{hb\bar{b}}, Y_{Ab\bar{b}})$ or $(Y_{hb\bar{b}}, Y_{Ab\bar{b}})$ are enhanced at large $an \beta!$



Total enhancement of $xx ightarrow \mathcal{H} ightarrow yy$ channel

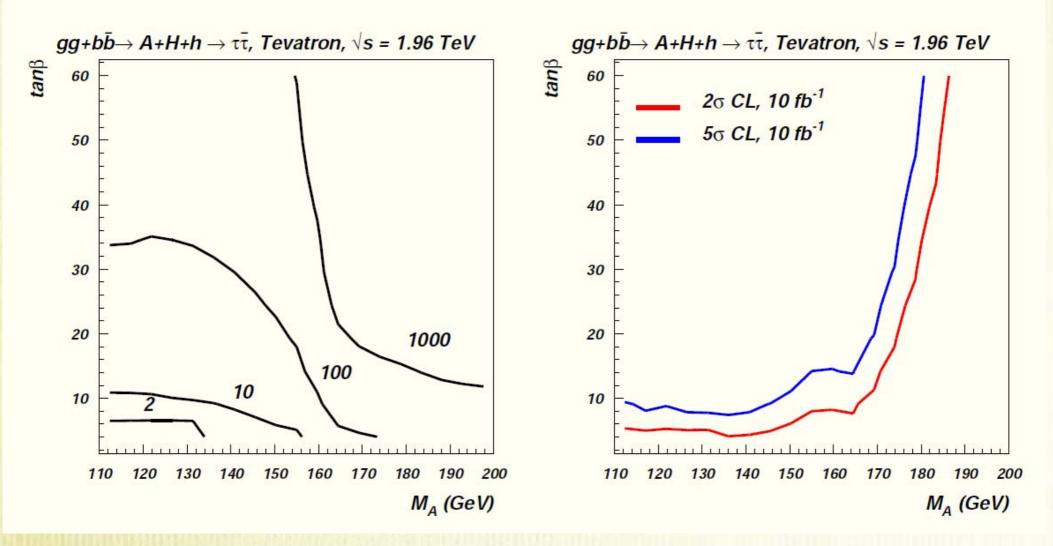
$$\kappa^{\mathcal{H}}_{total/xx} = [\kappa^{\mathcal{H}}_{gg/xx} + \kappa^{\mathcal{H}}_{bb/xx}R_{bb:gg}$$



Visibility of MSSM Higgs bosons: $\tau \tau$ channel

Predicted Tevatron reach, based on the $h_{SM}
ightarrow au^+ au^-$ studies

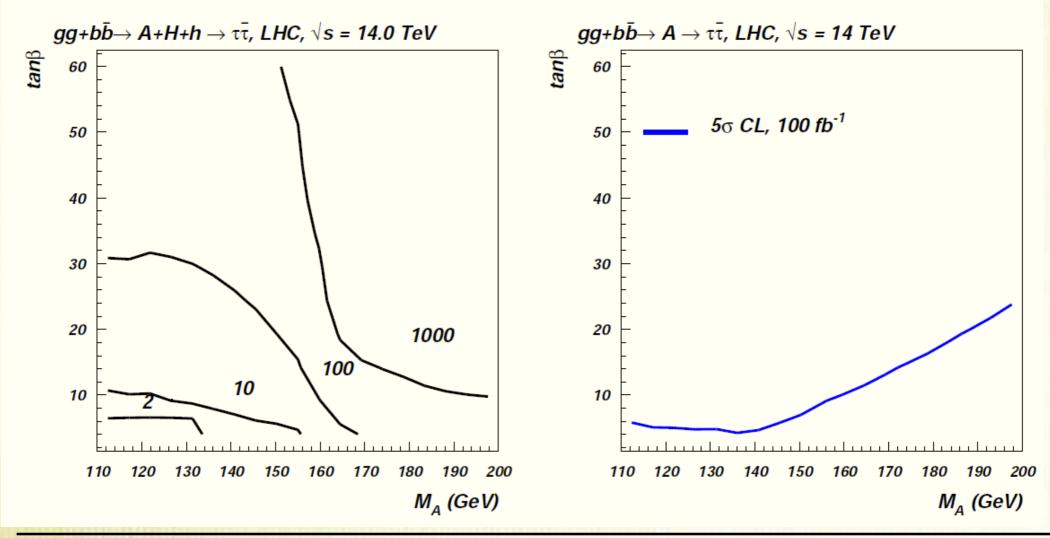
by A.B., T.Han, R.Rosenfeld, hep-ph/0204210



Visibility of MSSM Higgs bosons: $\tau \tau$ channel

Predicted LHC reach, based on the $h_{SM} ightarrow au^+ au^-$ studies

by D.Cavalli et al, hep-ph/0203056



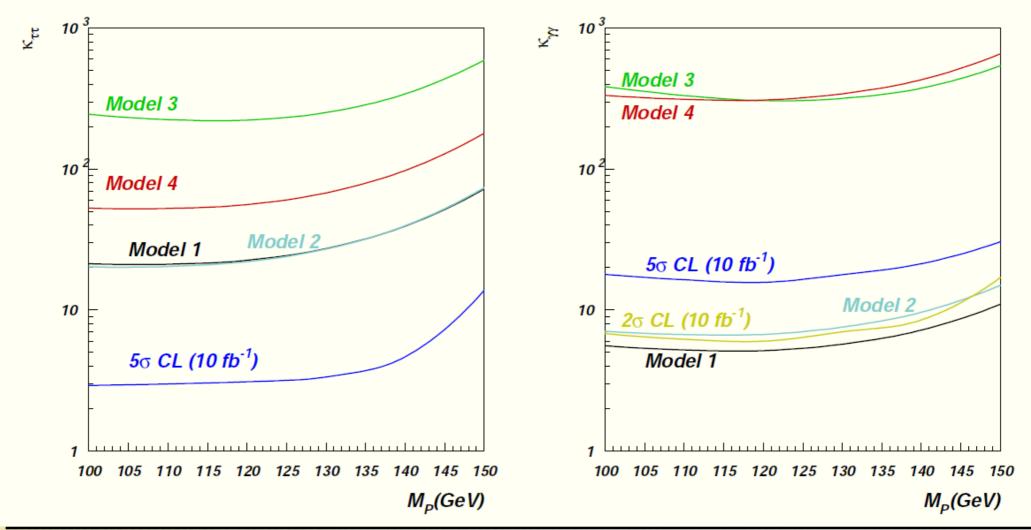
What happens in alternative models of EWSB? Technicolor

- Scalar states involved in EWSB are manifestly composite at scales not much above the electroweak scale $v \sim 250$ GeV
- A new asymptotically free strong gauge interaction, Technicolor, (Susskind, Weinberg) breaks the chiral symmetries of massless fermions
 additional light neutral pseudo Nambu-Goldstone bosons: "technipions" in Technicolor models

$$\begin{split} &\Gamma(P \to gg) = \frac{m_P^3}{8\pi} \left(\frac{\alpha_s N_{TC} \mathcal{A}_{gg}}{2\pi F_P} \right)^2, m_P = 130 \text{ GeV case} \\ \hline & 1) \text{ one-family } 2) \text{ variant one-family } 3) \text{ multiscale } 4) \text{ low-scale } \\ \hline & \mathcal{A}_{gg} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{6}} & \sqrt{2} & \frac{1}{\sqrt{3}} \\ \hline & \mathcal{A}_{\gamma\gamma} & -\frac{4}{3\sqrt{3}} & \frac{16}{3\sqrt{6}} & \frac{4\sqrt{2}}{3} & \frac{34}{9} \\ \hline & 1) \text{ one family } 2) \text{ variant one-family } 3) \text{ multiscale } 4) \text{ low scale } \\ \hline & \kappa_{gg \ prod} & 48 & 6 & 1200 & 120 \\ \hline & \kappa_{prod}^P & 48 & 6 & 1200 & 120 \\ \hline & \kappa_{prod}^P & 47 & 5.9 & 1100 & 120 \\ \hline & F_P^{(1)} = \frac{v}{2}, \ F_P^{(2)} = v, \ F_P^{(4)} = \frac{v}{\sqrt{10}}, \ F_P^{(3)} = \frac{v}{4} \end{split}$$

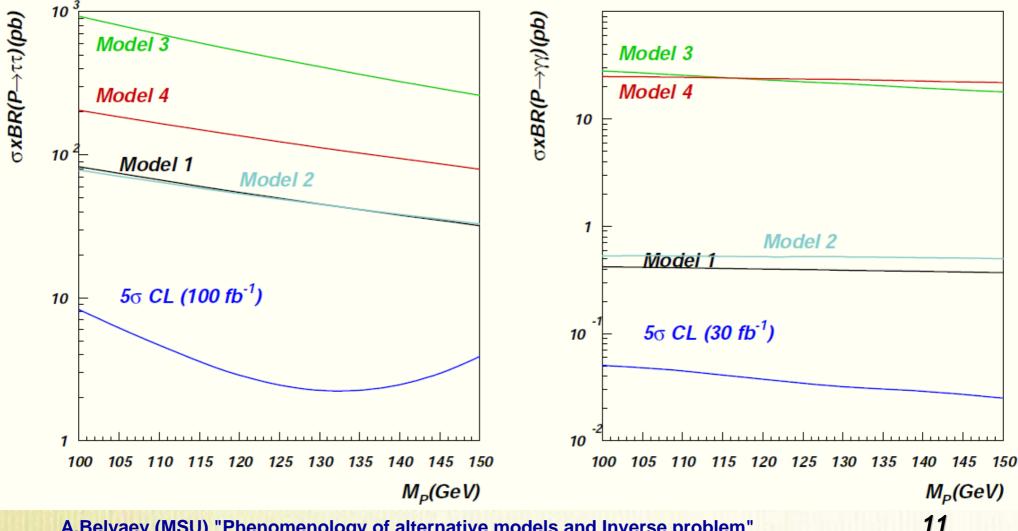
Visibility of Technipions: $\tau \tau$ and $\gamma \gamma$ channels

Predicted Tevatron reach, based on the $h_{SM} \rightarrow \tau^+ \tau^-$ studies by A.B., T.Han, R.Rosenfeld, hep-ph/0204210 and on the $h_{SM} \rightarrow \gamma \gamma$ studies by S. Mrenna and J. D. Wells, hep-ph/0001226



Visibility of Technipions: $\tau \tau$ and $\gamma \gamma$ channels

Predicted LHC reach, based on the $h_{SM} \rightarrow \tau^+ \tau^-$ studies by D.Cavalli et al, hep-ph/0203056 and on the $h_{SM}
ightarrow \gamma\gamma$ studies by R. Kinnunen, S. Lehti, A. Nikitenko and P. Salmi,hep-ph/0503067



Distinguishing SUSY from Technicolor models

- Tevatron and LHC have the potential to observe the light (pseudo) scalar states characteristic of both supersymmetry and models of dynamical symmetry breaking τ⁺τ⁻ channel!
- SUSY case: $\tau^+\tau^-$ channel is enhanced while the $\gamma\gamma$ channel is suppressed, and this suppression is strong enough that even the LHC would not observe the $\gamma\gamma$ signature.
- In contrast, for the dynamical symmetry breaking models studied we expect simultaneous enhancement of both the $\tau^+\tau^-$ and $\gamma\gamma$ channels.

Little Higgs Model as alternative to SUSY

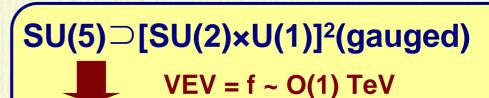
+Little Higgs scenario

New-Physics < 10TeV + EW ~ 0.1 TeV

Special arrangement for the cancellation of Λ^2 at 1-loop

+Littlest Higgs model

Arkani-Hamed, Cohen, Katz and Nelson (2002)



SO(5)⊃SU(2)×U(1) (diagonal)

New heavy particles

bosons,top-quarks,

scalars

λŧt

 T_R

 $-\lambda_t/(2f)$

h

 T_L

cancel the one-loop

quadratic divergences

h

 W', Z', γ'

Little Higgs Model with T -parity 4 Original LHM is suffered from severe constraints from EW observables. Tree-level corrections are due to the exchange of additional heavy gauge bosons and non-vanishing VEV of triplet higgs. f must be larger than 5 TeV, fine-tuning again ! **4**T-parity: Cheng, Low 2003, Z₂ symmetry $SU(2)_1 \times U(1)_1 \leftrightarrow SU(2)_2 \times U(1)_2$ **SM** particles \rightarrow + SM particles $(W_H, Z_H, A_H, \Phi, Q) \rightarrow$ $-(W_H, Z_H, A_H, \Phi, Q)$ No tree-level to EW observables The lightest T-odd particle is a good DM candidate New scale f can be lower then 1 TeV interesting phenomenology!

LHT Model

Hsin-Chia Cheng, Ian Low, Jay Hubisz, Patrick Meade, Andrew Noble, Maxim Perelstein, Claudio O. Dib, Rogerio Rosenfeld, Alfonso Zerwekh, Seung J. Lee, Gil Paz, Chuan-Ren Chen, Kazuhiro Tobe, C.-P. Yuan, Andreas Birkedal, ...

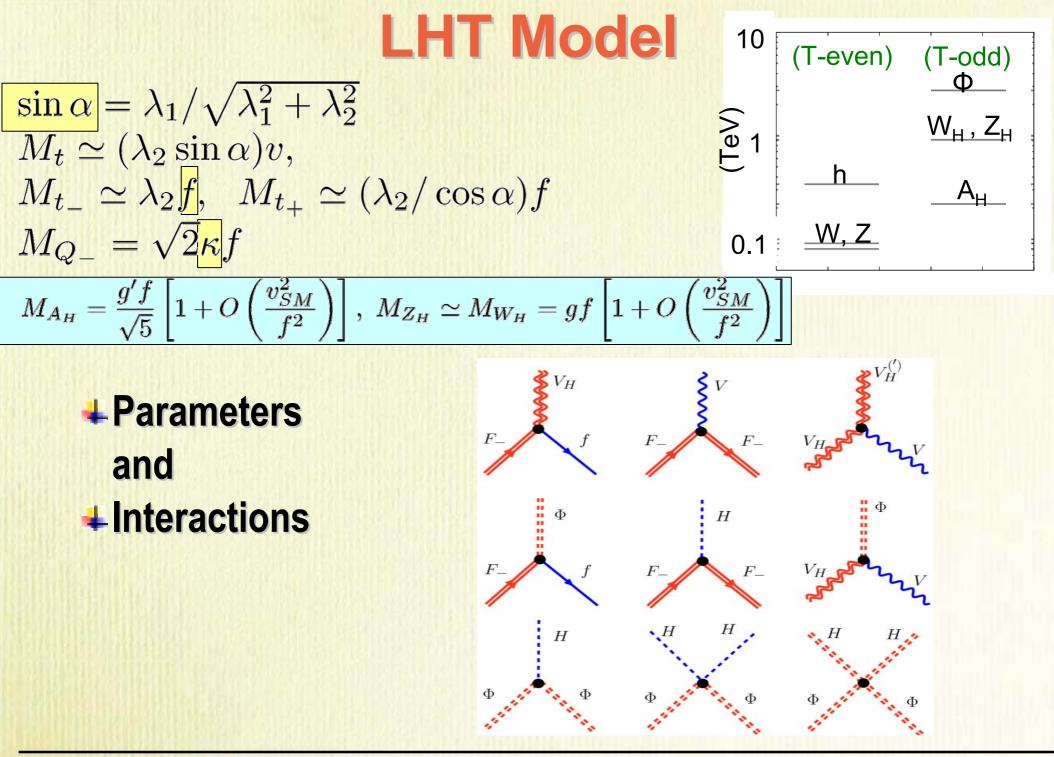
LHT Model is a good alternative, but

- Can LHT model phenomenology be distinguished from SUSY?
- The model can easily mimic SUSY signatures: cascade decays down to WIMP dark matter

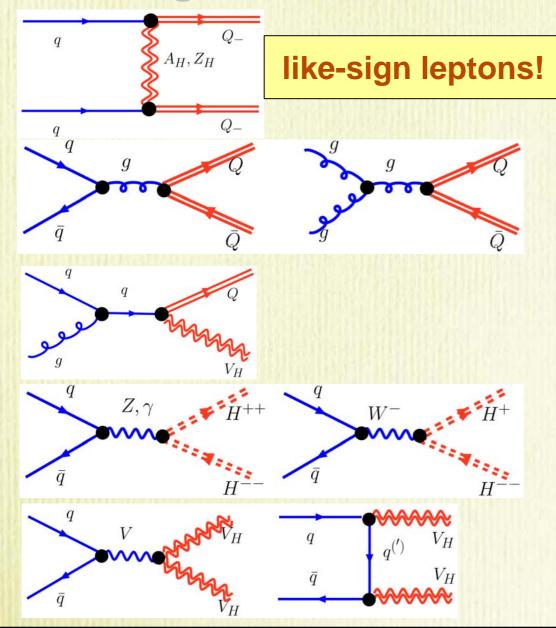
The pattern of LHT Model

- Hiss boson is light pseudo-Goldstone boson of SB global symmetry
- SM 1-loop quadratic divergences are canceled by new particles
- **4**Z₂ symmetry relaxes EW constraints and provides DM

Vector bosons: W_H^{\pm} , Z_H , A_H Fermions: $SU(2)_L$ singlets $-T_+, T_ SU(2)_L$ doublets $-Q_L^-, L_L^-$ Scalars: scalar $SU(2)_L$ triplet $-\phi^{\pm\pm}, \phi^{\pm}, \phi^0, \phi^P$

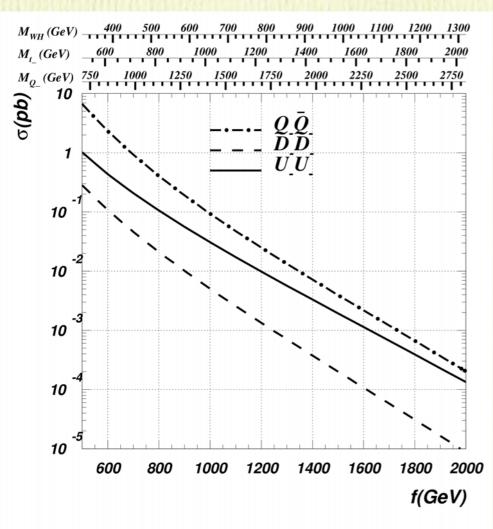


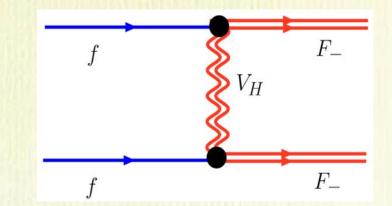
Phenomenology of LHT model Promising Processes classification



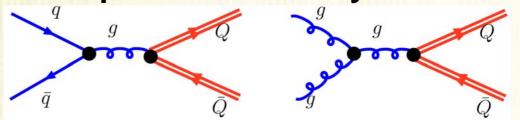
Heavy quarks production rates and signatures

$$\sin \alpha = 1/\sqrt{2} \ (\lambda_1 = \lambda_2), \kappa = 1$$

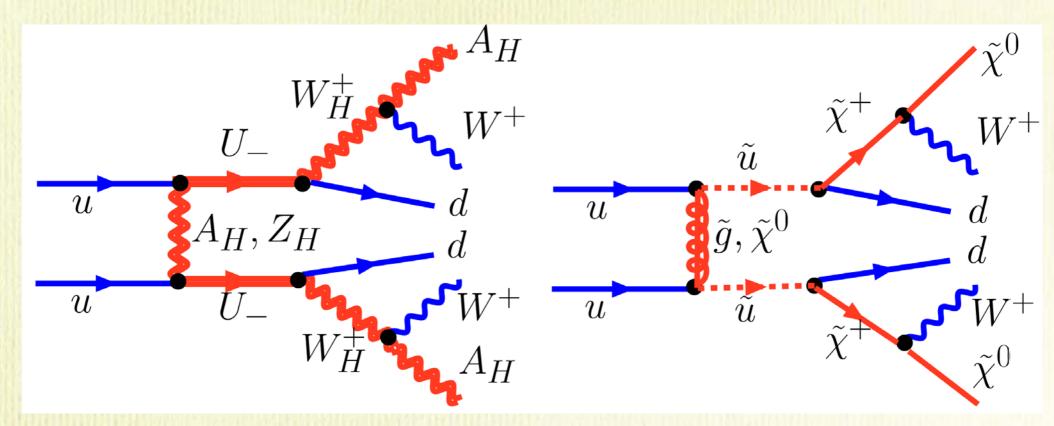




EW production due to the initial double valence quarks leads to like sign lepton signature (LSL), it is comparable to strong production and becomes even more important for heavier masses due to parton luminosity behavior!



LHT and SUSY cascade decays

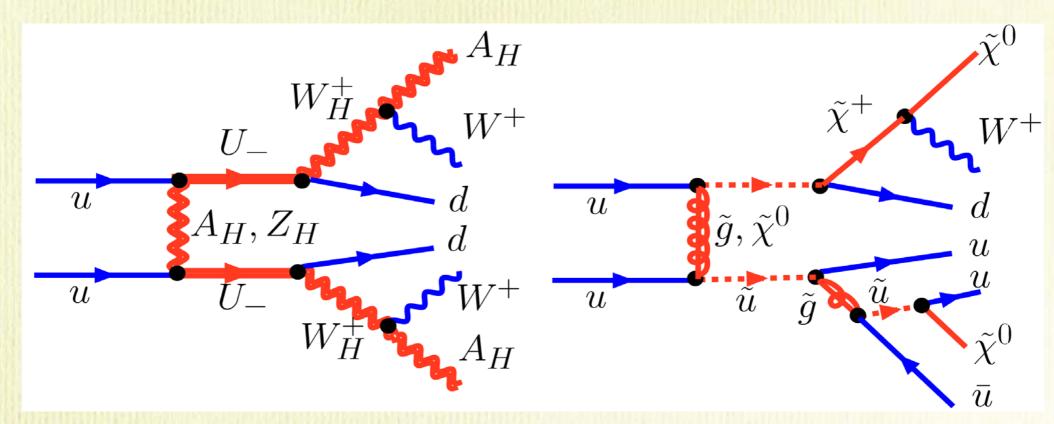


Both, SUSY and LHT give a very similar signature pattern

 $egin{aligned} \lambda_1 &= \lambda_2 = 1 \ f &= 1 \; TeV, \;\; \kappa = 1 \ Br(Q &
ightarrow W_H q') = 0.62 \ Br(W_H &
ightarrow WA_H) = 1 \end{aligned}$

One should look closely: various decay channels, spin correlations, couplings

LHT and SUSY cascade decays

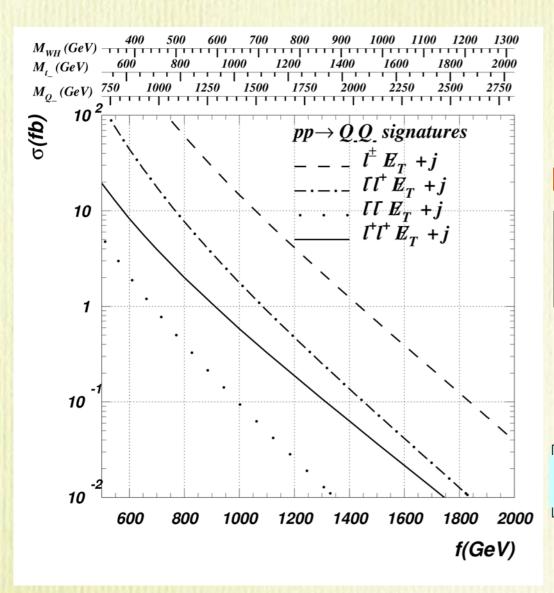


Gluon has no partner in LHT model!

 $egin{aligned} \lambda_1 &= \lambda_2 = 1 \ f &= 1 \; TeV, \;\; \kappa = 1 \ Br(Q &\to W_H q') &= 0.62 \ Br(W_H &\to W A_H) &= 1 \end{aligned}$

Study of spins and couplings is quite a challenge at the LHC

Heavy quarks production rates and signatures



$$egin{aligned} \lambda_1 &= \lambda_2 = 1 \ f &= 1 \; TeV, \;\; \kappa = 1 \ Br(Q &
ightarrow W_H q') = 0.62 \ Br(W_H &
ightarrow WA_H) = 1 \end{aligned}$$

Like-sign lepton signature (LSL)

$$\begin{array}{l} qq \to QQ \\ (Q \to W_H^+ q') \to W_H^+ W_H^+ q'q' \end{array}$$

Opposite sign lepton signature and 1-lepton signature (1L)

$$q\bar{q}(gg) \to Q\bar{Q} \to W_H^+ W_H^- q'\bar{q}'$$



- Understanding the Underlying Model can be a serious problem since SUSY alternatives having partial SUSY properties lead to a very similar signatures
- To discriminate underlying theory (not just SUSY!) one should actually perform
 - studies opposite to inverse problem -
 - we should understand the signal pattern for specific models to identify them in the end!
- Study of fine effects such as spin and couplings of new particles is necessary.
 - This task might be problematic at the LHC. If this is the case, then ILC would be precious but important tool to accomplish it.
- **Lets keep up our efforts!**

Technicolor enhancement factor for production and decay

$\Gamma(P o gg)=rac{m_P^3}{8\pi}\left(rac{lpha_s N_{TC} {\cal A}_{gg}}{2\pi F_P} ight)^2$, $m_P=130$ GeV case					
	1) one-family	2) variant one-family	3) multiscale	4) low-scale	
\mathcal{A}_{gg}	$\frac{1}{\sqrt{3}}$	$\frac{1}{\sqrt{6}}$	$\sqrt{2}$	$\frac{1}{\sqrt{3}}$	
Arr	$-\frac{4}{3\sqrt{3}}$	$\frac{16}{3\sqrt{6}}$	$\frac{4\sqrt{2}}{3}$	$\frac{34}{9}$	
	1) one family	2) variant one-family	3) multiscale	4) low scale	
$\kappa^P_{gg\ prod}$	48	6	1200	120	
$\kappa^P_{bb\ prod}$	4	0.67	16	10	
κ^P_{prod}	47	5.9	1100	120	
Decay	1) one family	2) variant	3) multiscale	4) low scale	SM Higgs
Channel		one family			
$b\overline{b}$	0.60	0.53	0.23	0.60	0.53
$c\overline{c}$	0.05	0	0.03	0.05	0.02
$ au^+ au^-$	0.03	0.25	0.01	0.03	0.05
<i>99</i>	0.32	0.21	0.73	0.32	0.07
$\gamma\gamma$	$2.7 imes10^{-4}$	$2.9 imes10^{-3}$	$6.1 imes10^{-4}$	$6.4 imes10^{-3}$	$2.2 imes10^{-3}$
<i>w</i> + <i>w</i> -	0	0	0	0	0.29