

Acceleration and Transport of High Energy Particles in Galaxy Clusters

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Focus of this talk

$$r_g(p) = \frac{pc}{eB} \simeq 3 \times 10^{12} \left(\frac{pc}{\text{GeV}} \right) \left(\frac{B}{\mu\text{G}} \right) \text{cm}$$
$$= 10^{-6} \text{ pc}$$

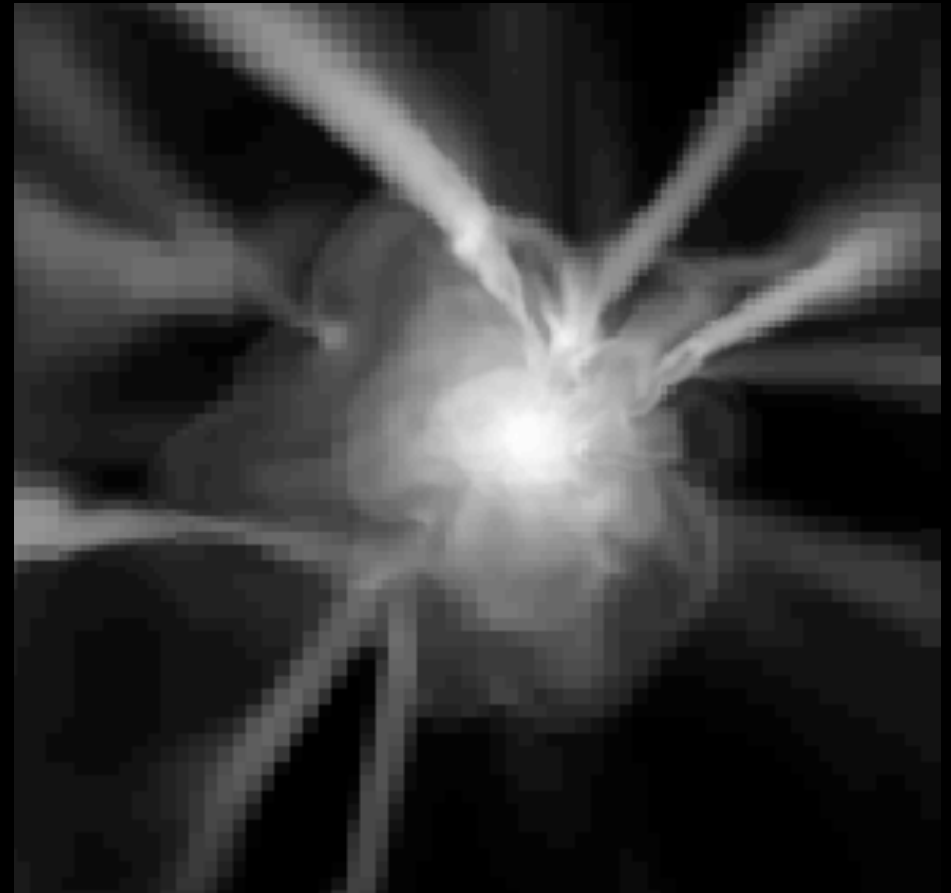
Microscale-Astrophysics of Galaxy Clusters

Cluster Shocks

Structure formation shocks have high Mach numbers (Miniati+ 2000, Ryu+ 2003, Pfrommer+ 2005, Skillman+ 2008)

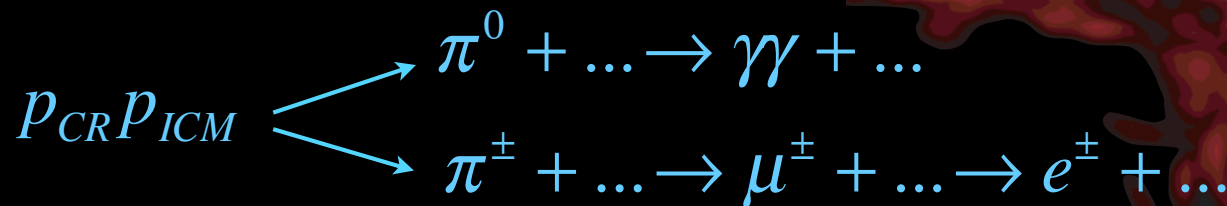
Shocks classes :
Internal $M \sim 4-5$
External $M \geq 10$

Collisionless shocks dissipate into CRs
(Krymsky 1977, Skadron+ 1977,
Bell 1978, Balndford & Ostriker 1978)



Non-thermal Emission

Straightforward Predictions:

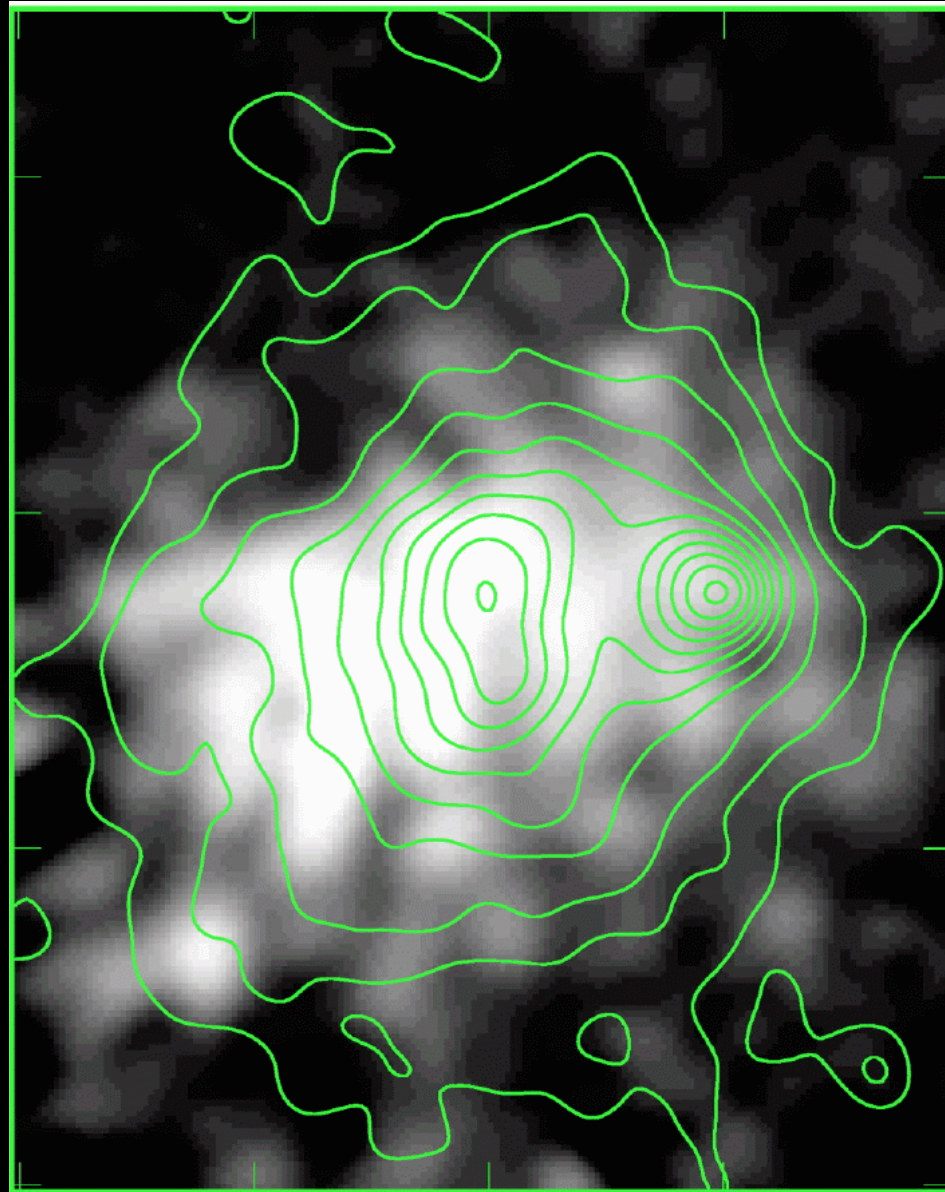


- i) gamma-ray emission
- ii) radio emission

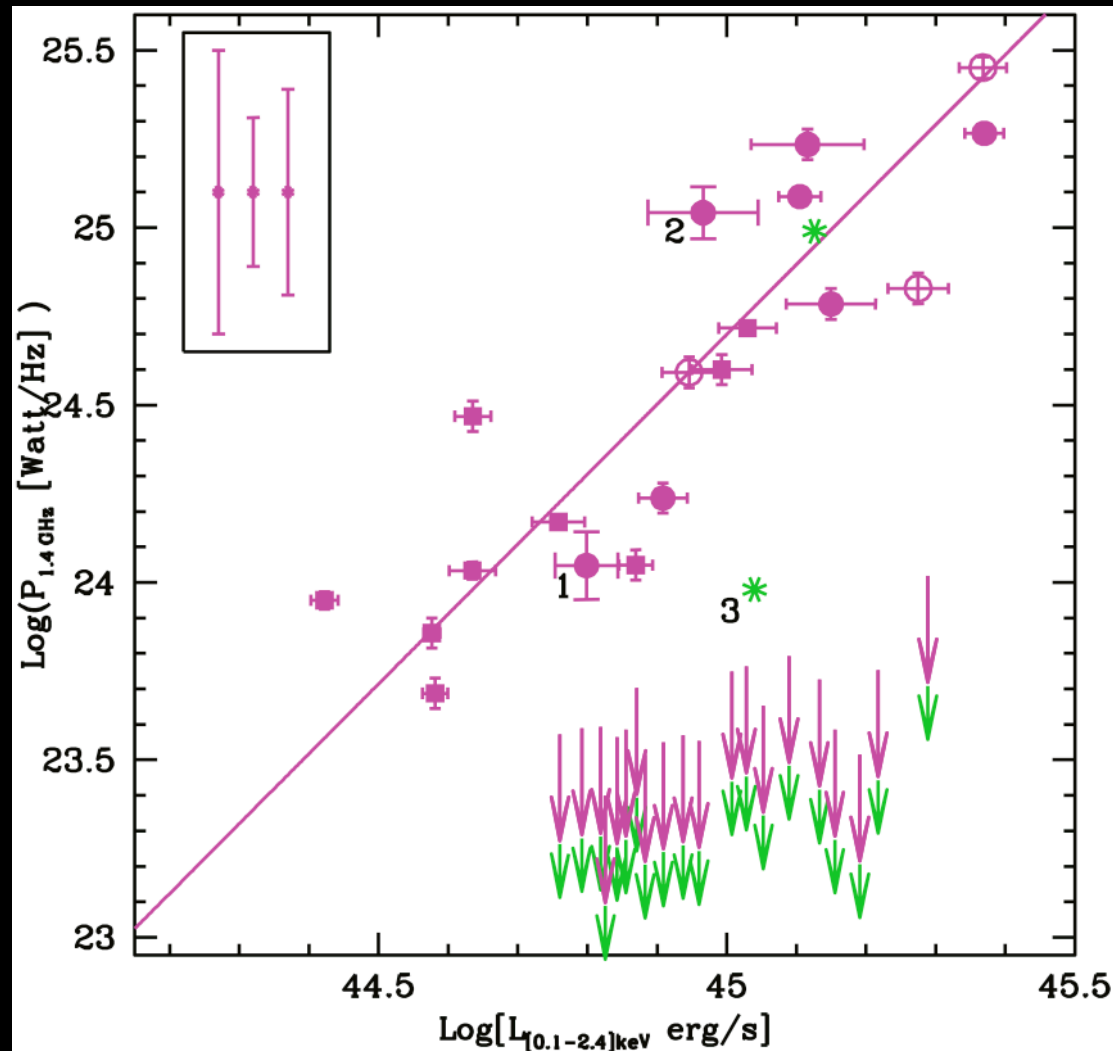
gamma-ray upper limits from
Fermi, H.E.S.S, MAGIC

Miniati 2003

Bullet Cluster at Radio WL



Bimodality in Cluster Diffuse Radio Emission



Brunetti et al. (2007)

What's wrong ?

Simulation results based on following simple assumptions:

- A. Accretion shocks accelerate CR protons
- B. Magnetic fields are seeded at high- z more or less uniformly
- C. CRs diffusion negligible (CRs tied up to ICM by B)

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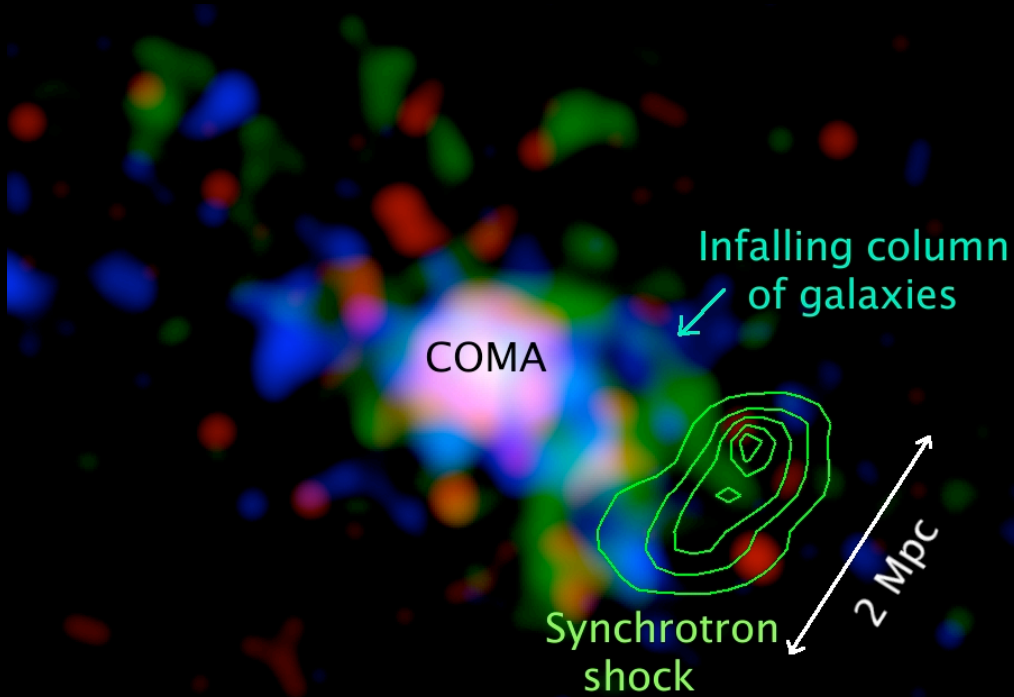
(A): Accretion shocks accelerate CR protons(?)

- SNRs
- radio relics show that e^- are accelerated at cosmic shocks
- Particle-in-Cell simulations indicate that at least relativistic collision-less shocks produce supra-thermal particles even if initially unmagnetized (e.g. Spitovsky 2008)

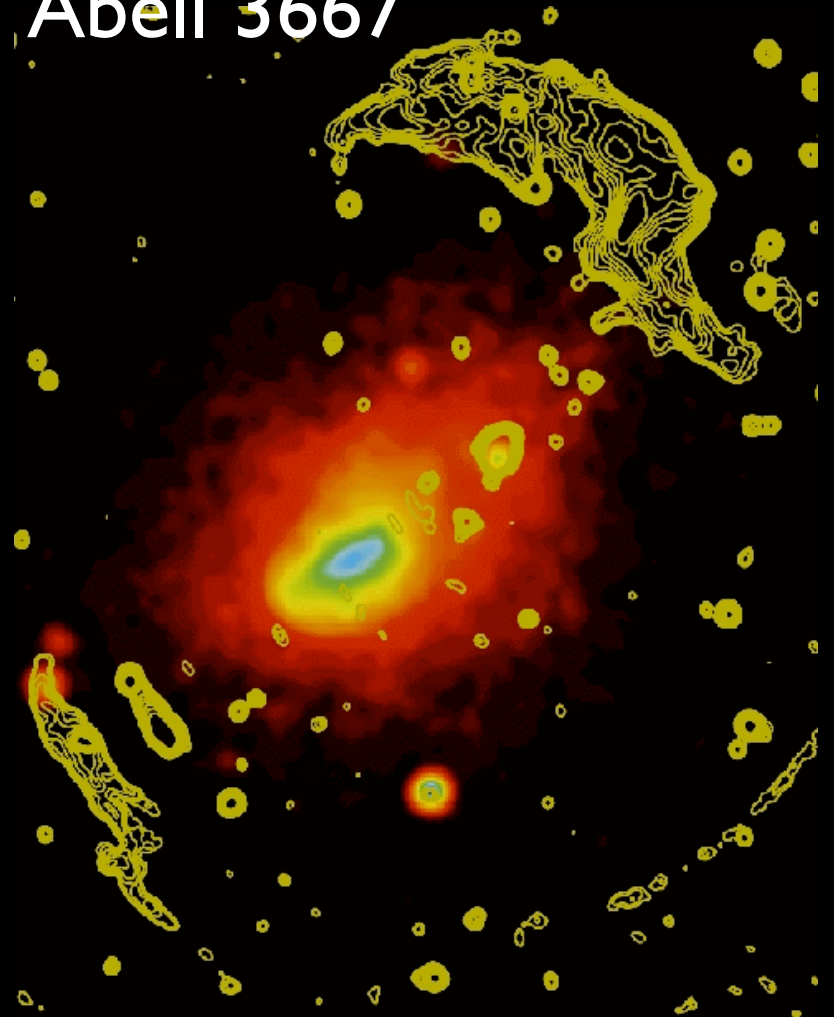
CR are accelerated at filament termination shocks

Relic Radio Emission

Abell 3667

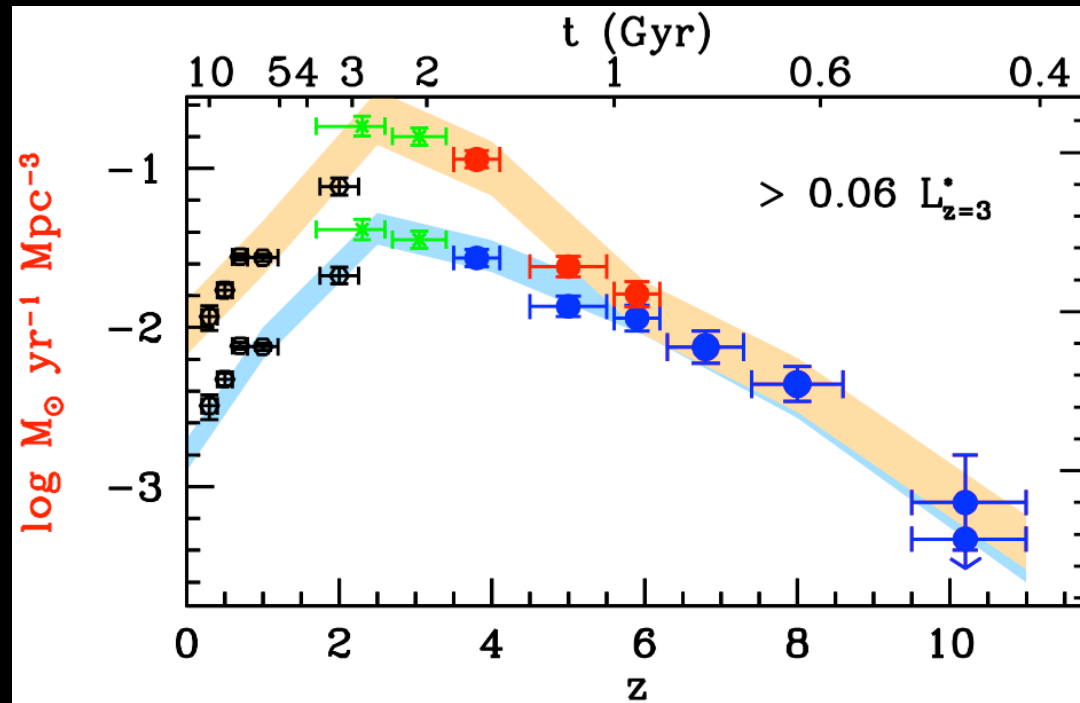


Bown & Rudnick 2010



credits to M. Johnston-Hollitt

CRs from Galaxies



Bouwens et al 2010

For example, Voelk et al (1996) find $P_{\text{CR}}/P_{\text{th}} \sim 5\text{-}10\%$

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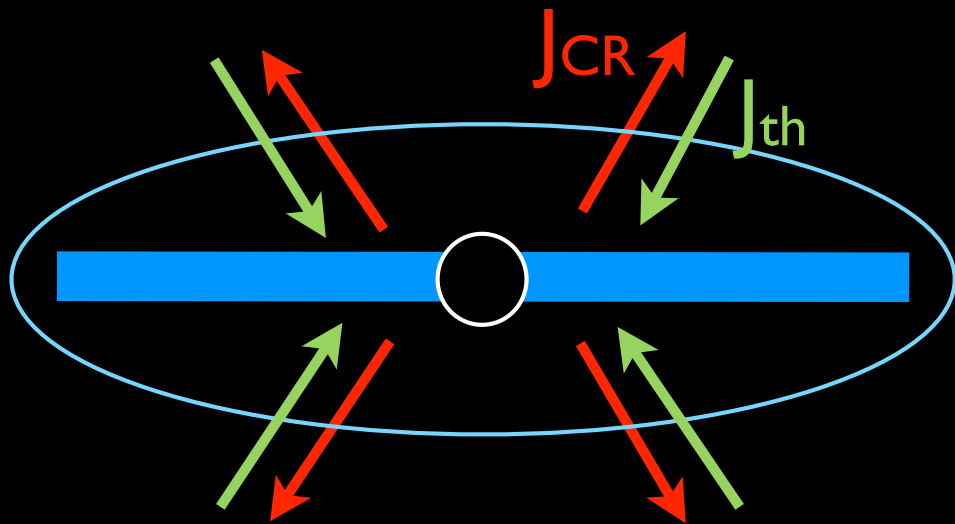
(B): Bimodality in B-field

$$j_{synch} \propto n_{CR}(E) \frac{B^2}{B^2 + B_{CMB}^2}$$

- non radio-halos GCs have μG magnetic fields (RM, from Clarke et al. 2001, 2004)
- magnetic seeds are probably generated uniformly in the proto-cluster regions (Miniati & Bell 2010, Kulsrud+ 1997, Gnedin + 2000, Vogt+2005, Donnert+2009)

Resistive Mechanism

Miniati & Bell 2010 (arXiv:1001.2011)



CR current, j_{cr} , drives a return current in the thermal plasma, j_{th} , that tends to cancel j_{cr} itself.

$$\vec{E}' = \frac{\vec{j}_{th}}{\sigma},$$

$$\sigma_{Spitzer} \approx 10^7 \left(\frac{T}{K} \right)^{3/2} s^{-1}$$

$$\frac{\partial \vec{B}}{\partial t} = -c \vec{\nabla} \times \vec{E} \approx 10^{-16} \frac{G}{Gyr}$$

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(C:) CR diffusion negligibly small ?

Voelk et al (1996):

- fluctuations $\frac{\delta B}{B} \sim 1$ driven by galaxy motions
- fluctuations extend down to the gyration radius of CR particles, i.e.:

$$r_g(p) = \frac{pc}{eB} \simeq 3 \times 10^{12} \left(\frac{pc}{\text{GeV}} \right) \left(\frac{B}{\mu\text{G}} \right) \text{cm}$$

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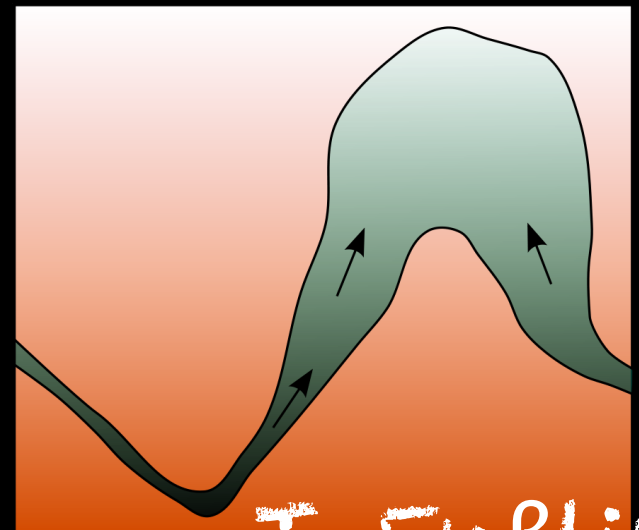
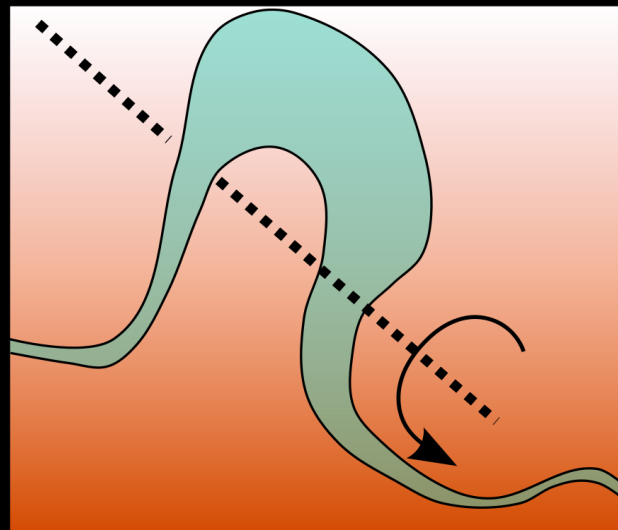
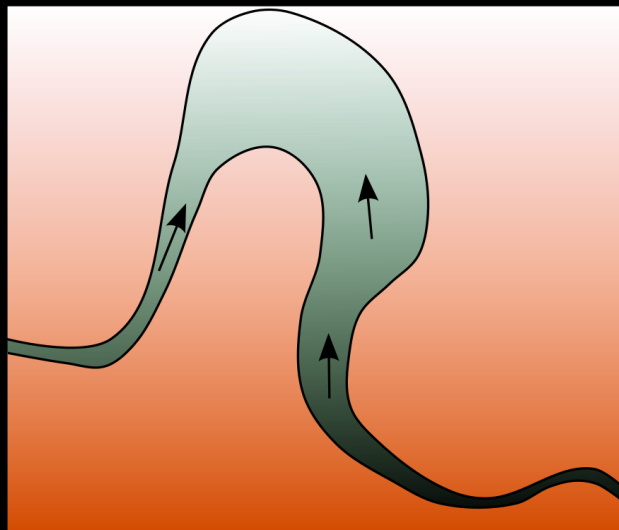
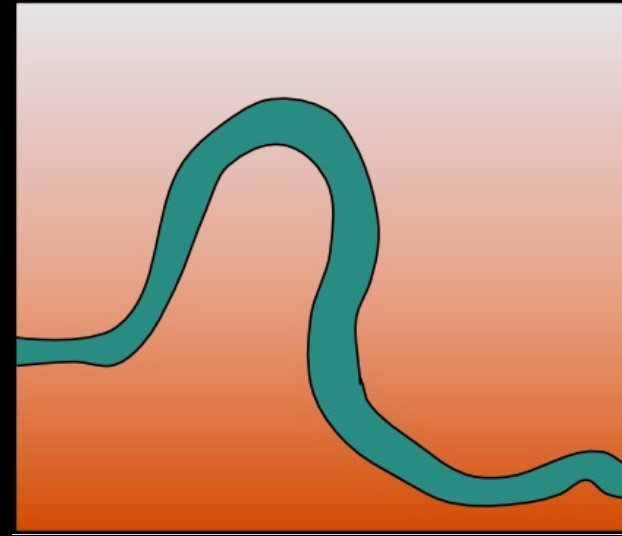
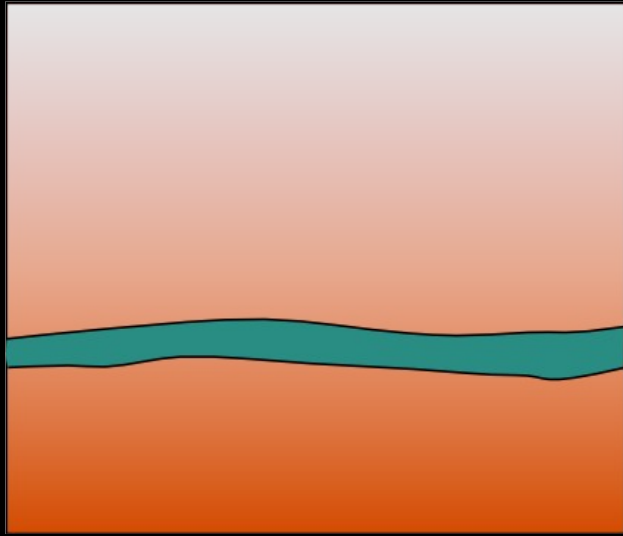
- However, this neglects damping of the waves by plasma effects, e.g. transit-time damping, cyclotron resonance, non-linear Landau damping...
- For example during a merger when the turbulent motions are high, Brunetti & Lazarian (2007) estimate a cutoff of the power spectrum of the magnetosonic waves:

$$\lambda_{cutoff} = \frac{1}{3} \text{kpc} \gg r_g \sim 10^{-9} \text{kpc}$$

(C:) CR diffusion negligibly small ?

- After the merger has relaxed, the turbulence level decreases significantly and things get even worse, i.e. damping significant at even larger wavelengths.
- At this two effects are important:
 - i) streaming of CRs generates the necessary scattering (Achterberg 1981, Foote&Kulsrud 1979, Yan&Lazarian 2008)
 - ii) mirroring from bent magnetic fields lines

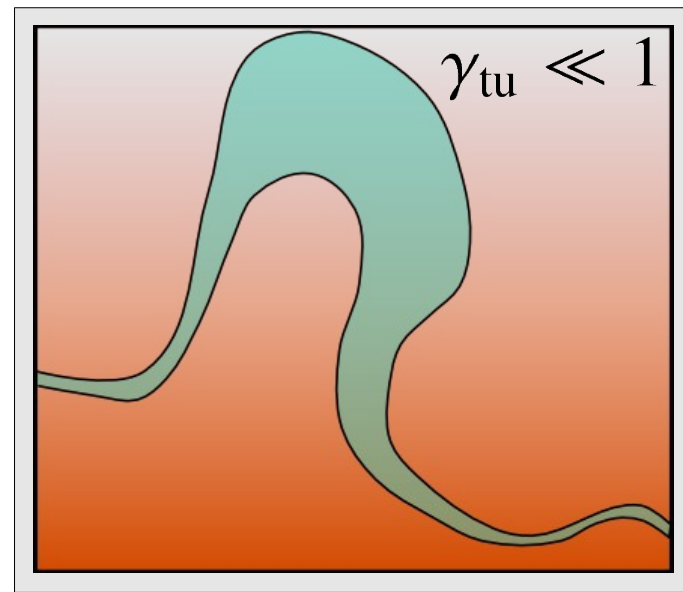
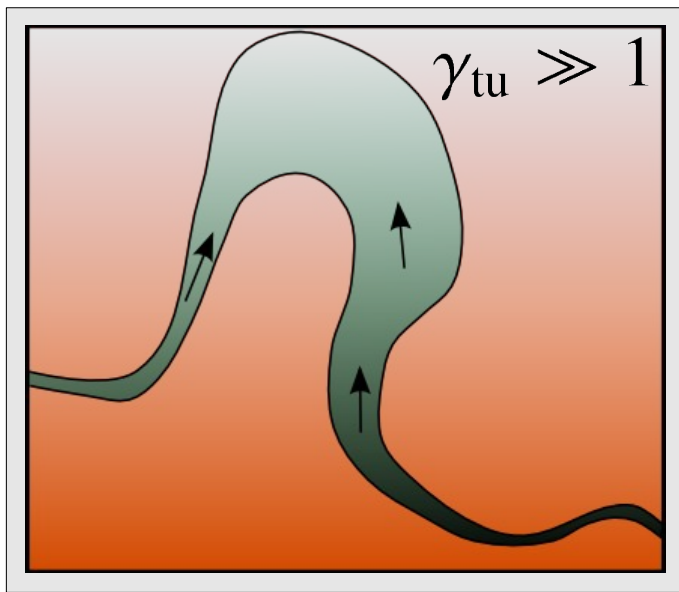
Rolling Flux-tubes



T. ENBLIN

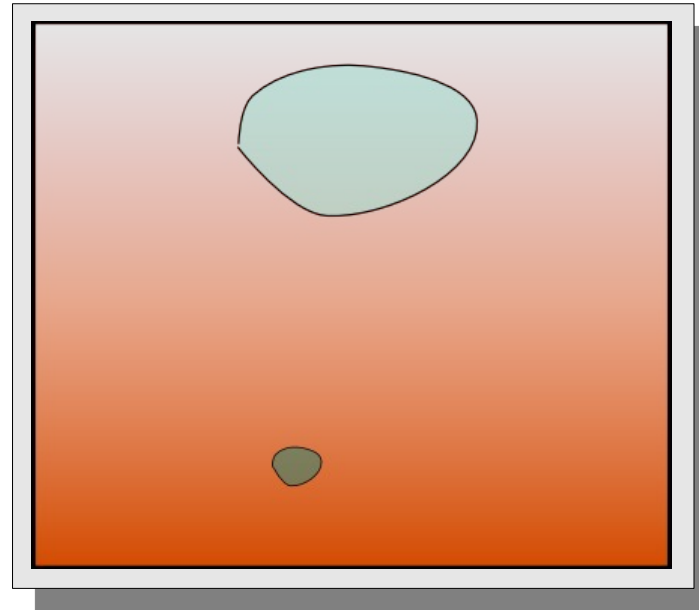
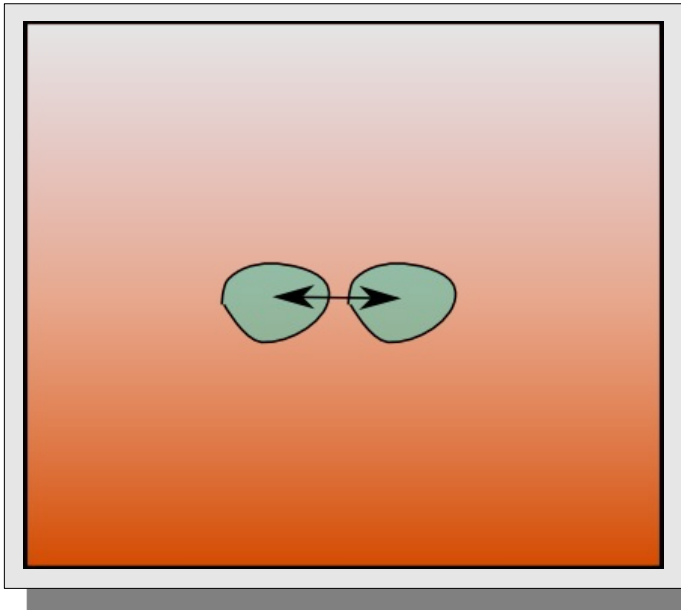
turbulent to streaming ratio

$$\gamma_{tu} = \frac{u_{tu}}{u_{st}}$$



CR profile due to advection

$$\eta(r) = \left(\frac{P(r)}{P_0} \right)^{\frac{3}{5}}$$



CR profile

$$f_{CR} : x_1 \rightarrow x_2 \quad \rho_{CR}(x_1) \rightarrow \rho_{CR}(x_2) = \rho_{CR}(x_1) \left(\frac{P_2}{P_1} \right)^{1/\gamma}$$

$$\frac{\rho_{CR}(x_2)}{\rho_{CR}(x_1)} = \left(\frac{P_2}{P_1} \right)^{1/\gamma}$$

CR transport model

$$\frac{\partial \rho_{CR}}{\partial t} = \vec{\nabla} \cdot (\rho_{CR} \vec{u}_{CR})$$

$$\vec{u}_{CR} = \vec{u}_{st} + \vec{u}_{ad} + \vec{u}_{di}$$

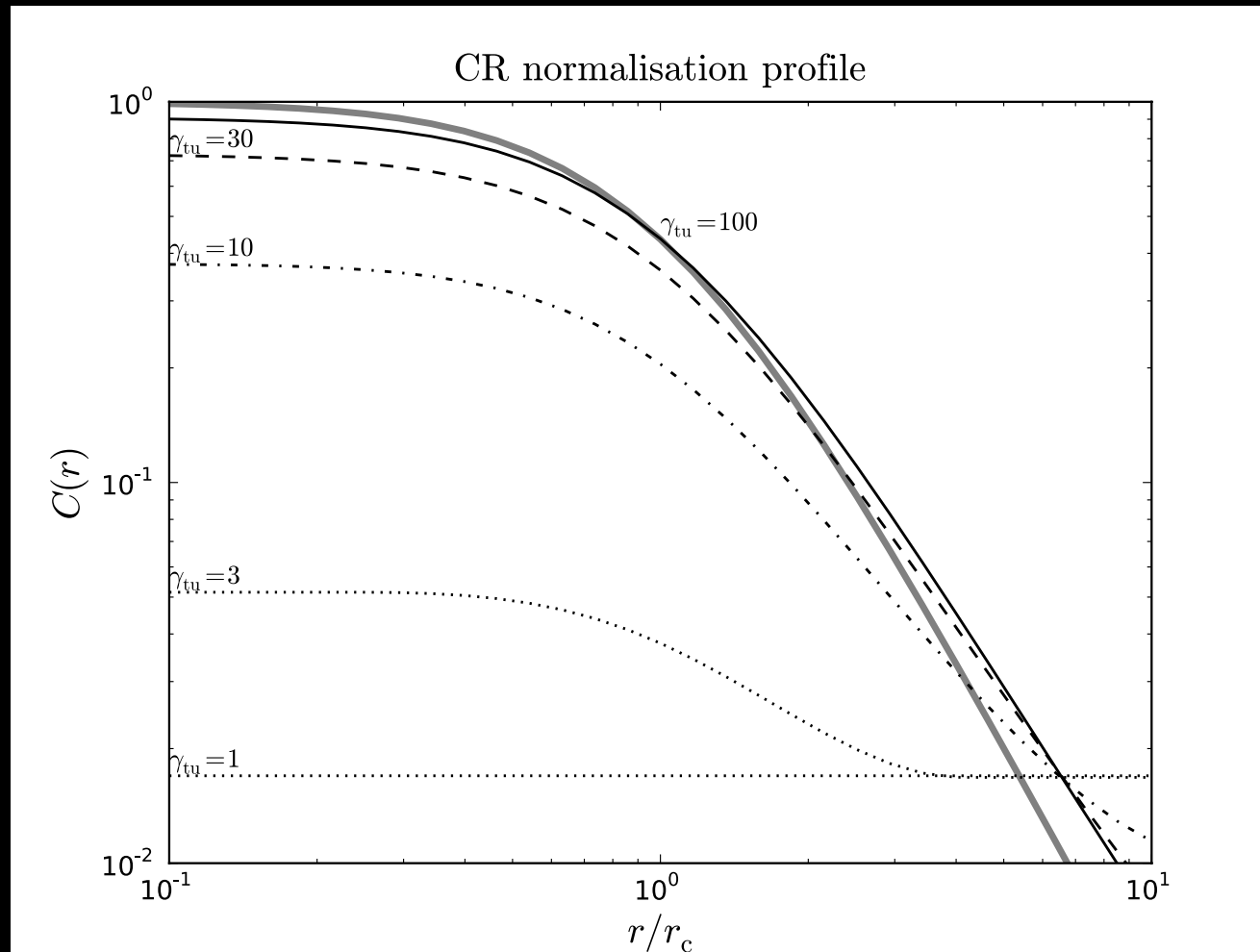
$$\vec{u}_{st} = -u_{st} \frac{\vec{\nabla} \rho_{CR}}{|\vec{\nabla} \rho_{CR}|}$$

$$\vec{u}_{ad} = -\kappa_{tu} \vec{\nabla} \ln \frac{\rho_{CR}}{\eta}$$

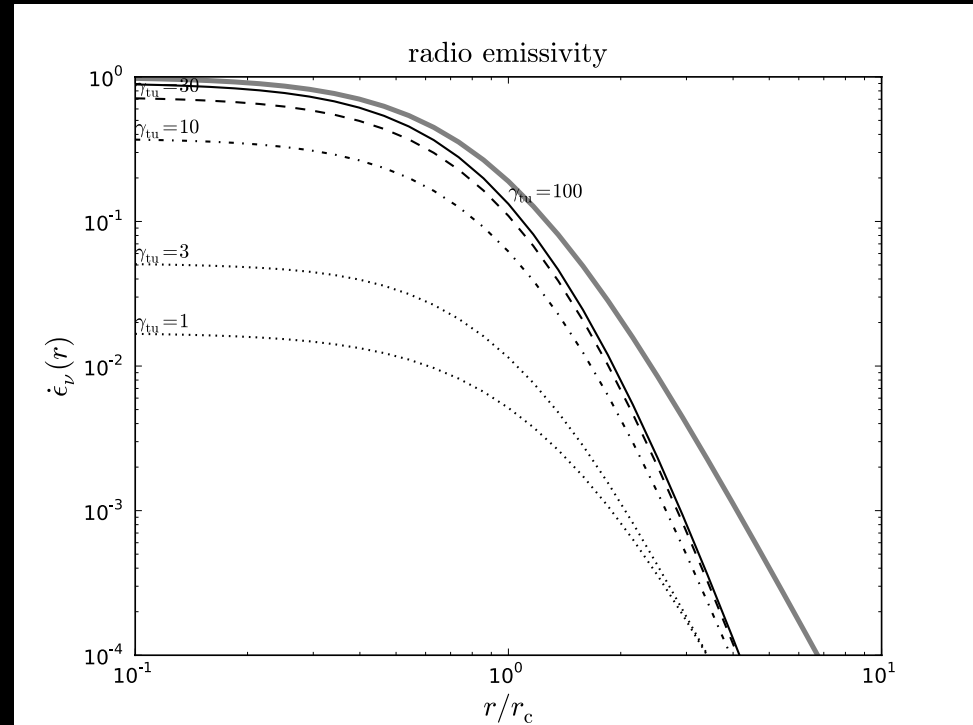
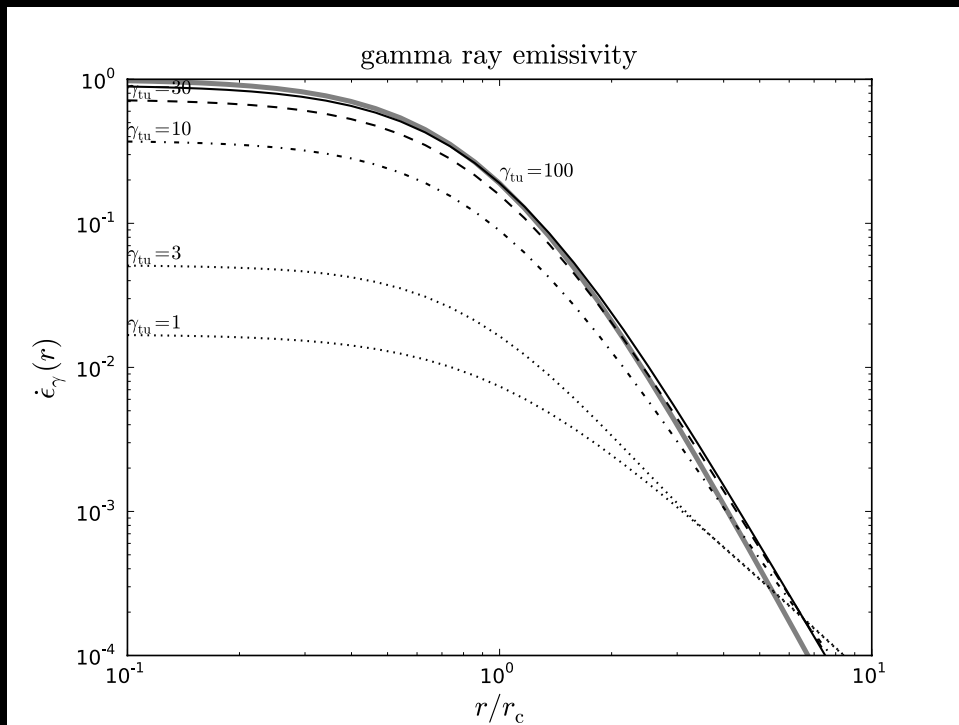
$$\kappa_{tu} = \frac{L_{tu} u_{tu}}{3}$$

$$\vec{u}_{di} = -\kappa_{di} \vec{\nabla} \ln \rho_{CR}$$

CR Spatial Distribution



Gamma-ray, Synchrotron Emissivities



Summary

- A. Not yet in a position to rule out acceleration of CR protons at structure formation shocks
- B. no problem with seeding the magnetic field
- C. Assumption that CR diffusion in GCs is negligibly small is most likely an oversimplification. CR transport in ICM needs important revision, with essential input from both theory and simulations.
- D. Radio halo emission powered by hadronic interactions switch-off timescales depend on streaming of CRs and magnetic field lines topology after the merger has settled