

# Coupling Jets to their Surroundings: the Role of Entropy

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Cen A VLP collaboration

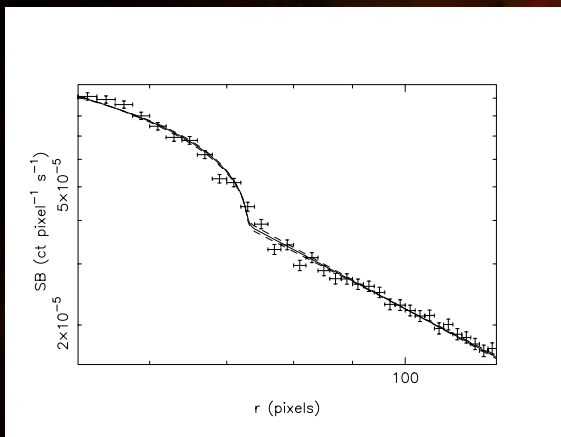
# Jet of Cygnus A

Powerful, near FR II radio source (Carilli & Barthel 1994); radio luminosity  $\approx 7 \times 10^{44}$  erg s $^{-1}$ ;  
 $z = 0.056$ ;  $D_L = 250$  Mpc; scale = 1.088 kpc/arcsec

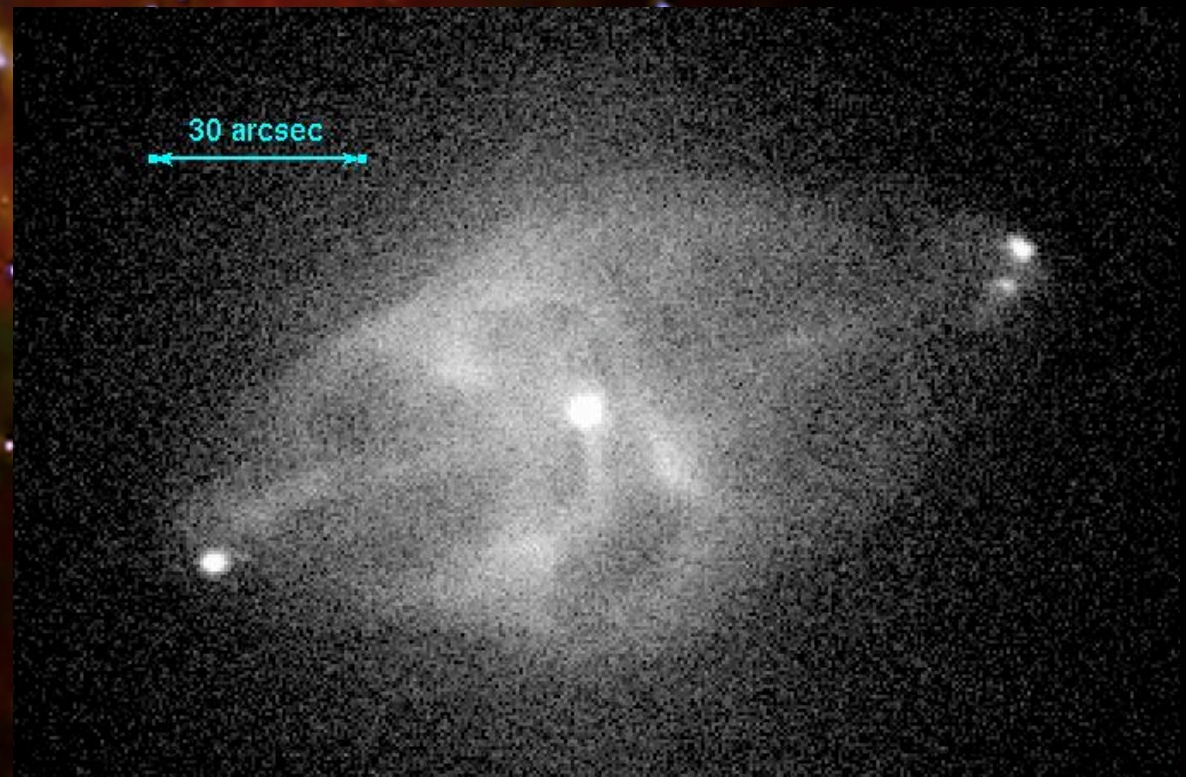
Hosted by cluster central galaxy

Smith et al (2002), Chandra shows AGN, jets, radio hotspots, cocoon shock, etc

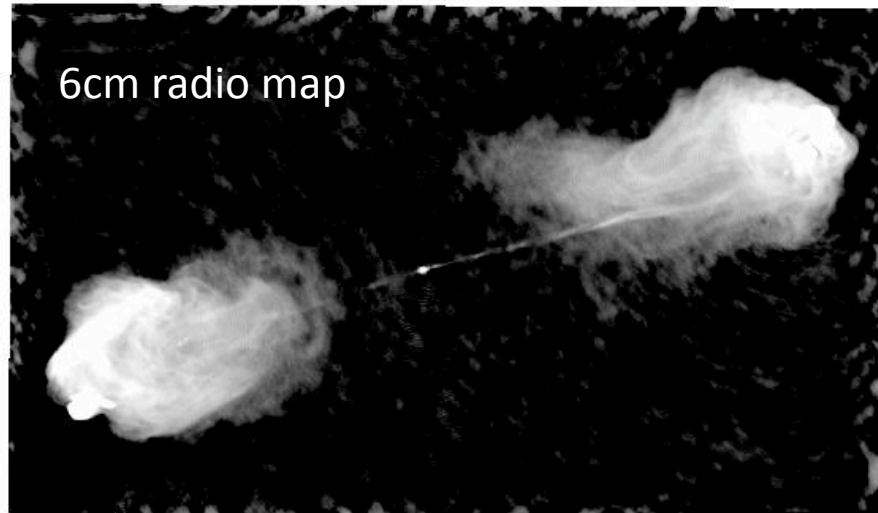
Chandra 0.5 – 7 keV



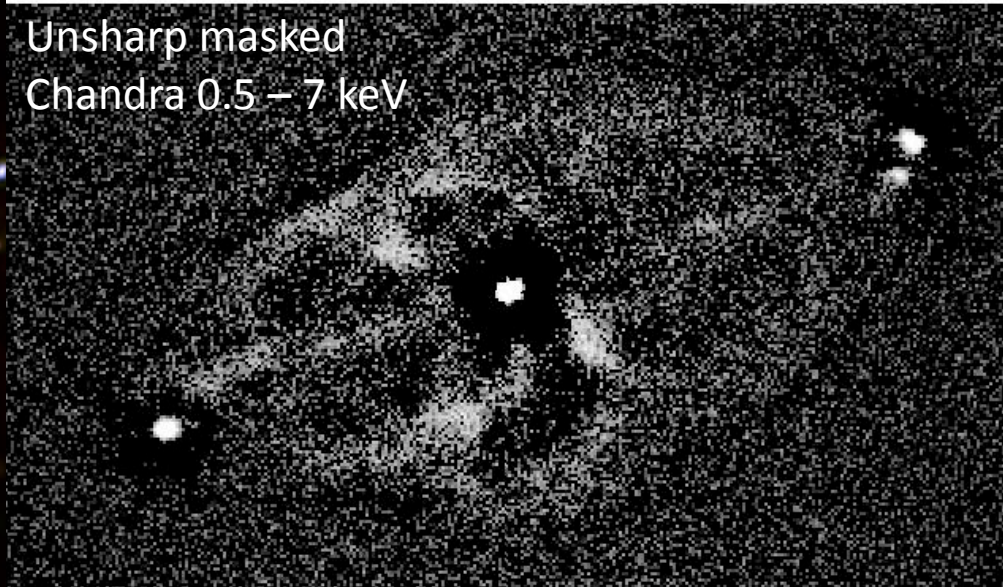
SW shock: Mach 1.37,  
 $r \approx 40$  kpc, age  $\approx 1.6 \times 10^7$  yr,  
mean power  $\approx 4 \times 10^{45}$  erg s $^{-1}$



# X-ray vs radio jet



Unsharp masked  
Chandra 0.5 – 7 keV



X-ray jet is resolved by Chandra

- not coincident with radio jet  
(Steenbrugge & Blundell 2007)

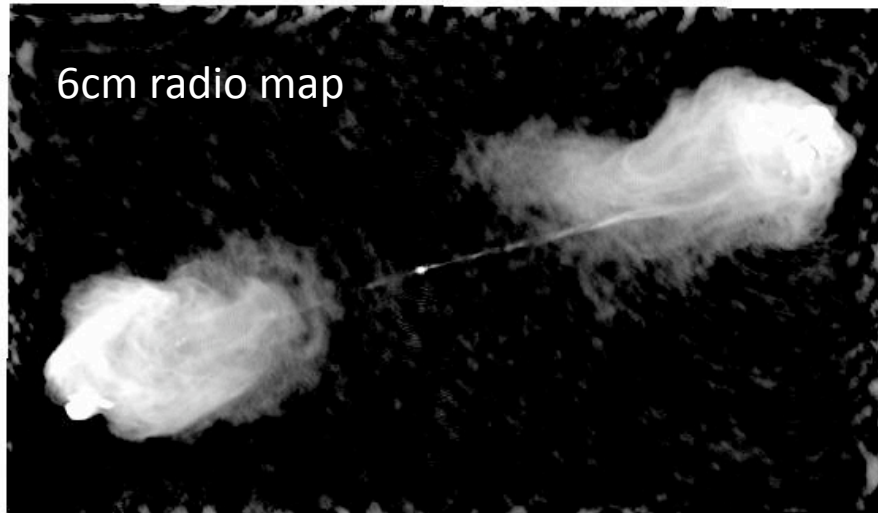
Symmetry (also: eastern jet receding)  
=> jet X-ray emission not Doppler  
boosted

X-ray spectrum for 11"×5.7" region at  
eastern end => power law (photon index  
 $1.69 \pm 0.26$ , 90%).

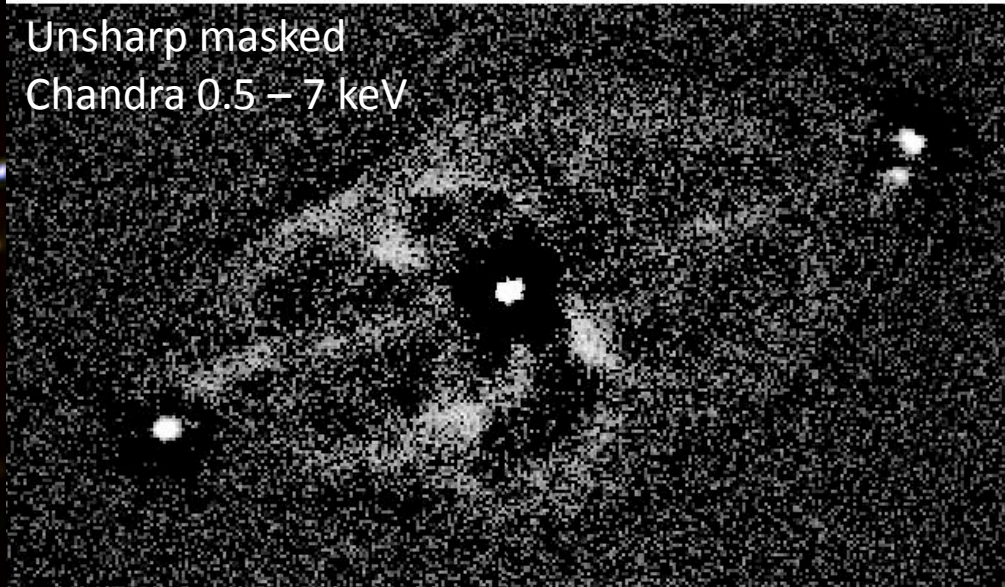
ICMB ( $100 < \gamma < 10,000$ ;  $p = 2.38$ )  
would require electron pressure  $> 10\times$   
surrounding gas pressure.

IC on beamed optical from AGN would  
require  $> 10^{46}$  erg  $s^{-1}$  beamed along jets

# X-ray vs radio jet



Unsharp masked  
Chandra 0.5 – 7 keV



Pressure equipartition in eastern lobe  
gives  $B \approx 55 \mu\text{G}$

– 1 keV synchrotron photons radiated by  
electrons with  $\gamma \approx 4 \times 10^7$

Synchrotron X-ray model requires  $\sim 10^{-5} \times$   
electrons as ICCMB model

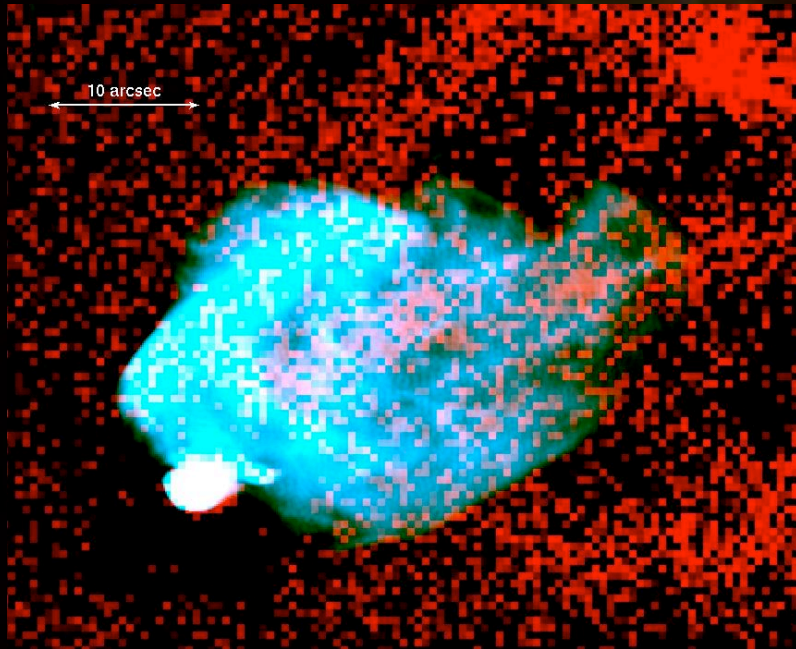
–  $t_{\text{synchrotron}} \approx 200 \text{ yr}$

– requires acceleration in situ

X-ray jet traces the path of the jet now

– radio does not

# X-ray vs radio jet



Unsharp masked 0.7 – 7 keV Chandra image (red)

Radio 6 cm (cyan)

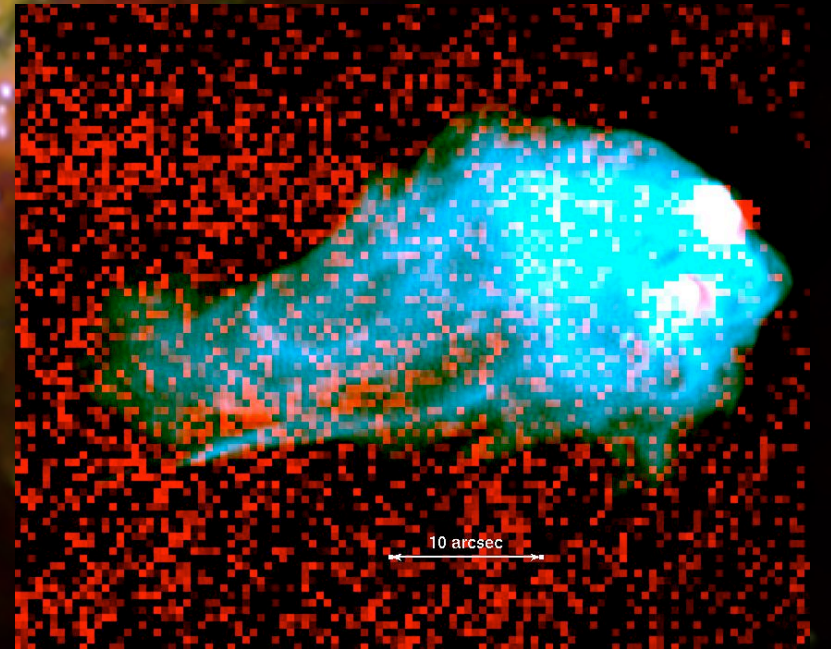
Excess hotspot pressure due to jet ram pressure

SSC model (Harris et al 1994) ->  $B_{\text{hotspot}} \approx 246 \mu\text{G}$ ;

Wilson (2003) ->  $B_{\text{hotspot}} \approx 150 \mu\text{G}$

$$\rho_{\text{hotspot}}/\rho_{\text{jet}} \approx (B_{\text{hotspot}}/B_{\text{jet}})^2 \approx 10 - 20$$

Jet radius in eastern lobe  $\approx 3 \text{ kpc}$



# Jet Flow Model

(Laing & Bridle 2002)

Proper density of jet rest mass,  $\rho$ , rate of mass flow through jet:

$$\dot{M} = \gamma \beta c A \rho$$

Power through jet ( $h$  = enthalpy per unit volume,  
 $h = p + e = \Gamma p / (\Gamma - 1)$ , for pressure  $p$ ):

$$P = (\gamma - 1) \dot{M} c^2 + h A c \beta \gamma^2$$

Momentum flux:

$$\Pi = (P/c + \dot{M} c) \beta$$

Have estimates for: power (from shock),  $P$ , area,  $A$ , and pressure,  $p$ .

Jet comes to (near) halt in hotspot => hotspot pressure  $\approx$  ram pressure, so  $\Pi \approx p_{\text{ram}} A$ , also known ( $p_{\text{ram}} = p_{\text{hotspot}}$ ; hotspot width matches jet)

Power + momentum equations give:

$$\frac{P}{p A c} = \left( \frac{p_{\text{ram}} / p}{\gamma + 1} + \frac{\Gamma}{\Gamma - 1} \right) \beta \gamma$$

– solve for jet speed, mass flow rate, etc

# Cyg A Jet Flow Model

For  $P_{\text{jet}} = 2 \times 10^{45} \text{ erg s}^{-1}$ , jet pressure  $p = 2.4 \times 10^{-10} \text{ erg cm}^{-3}$ ,  $r_{\text{jet}} = 3.1 \text{ kpc}$ , ratio of specific heats,  $\Gamma = 5/3$  (see below),  $p_{\text{ram}}/p = 20$ , get:

$\beta_{\text{jet}} = 0.079$ ; mass flow rate through jet =  $9 M_{\odot} \text{ yr}^{-1}$ ;  $n_{e,\text{jet}} = 4.4 \times 10^{-4} \text{ cm}^{-3}$ ; for pressure balance,  $kT_{\text{jet}} = 175 \text{ keV}$  (hence gas is non-relativistic).

Note a) 1-dimensional flow model, b) no correction for projection.

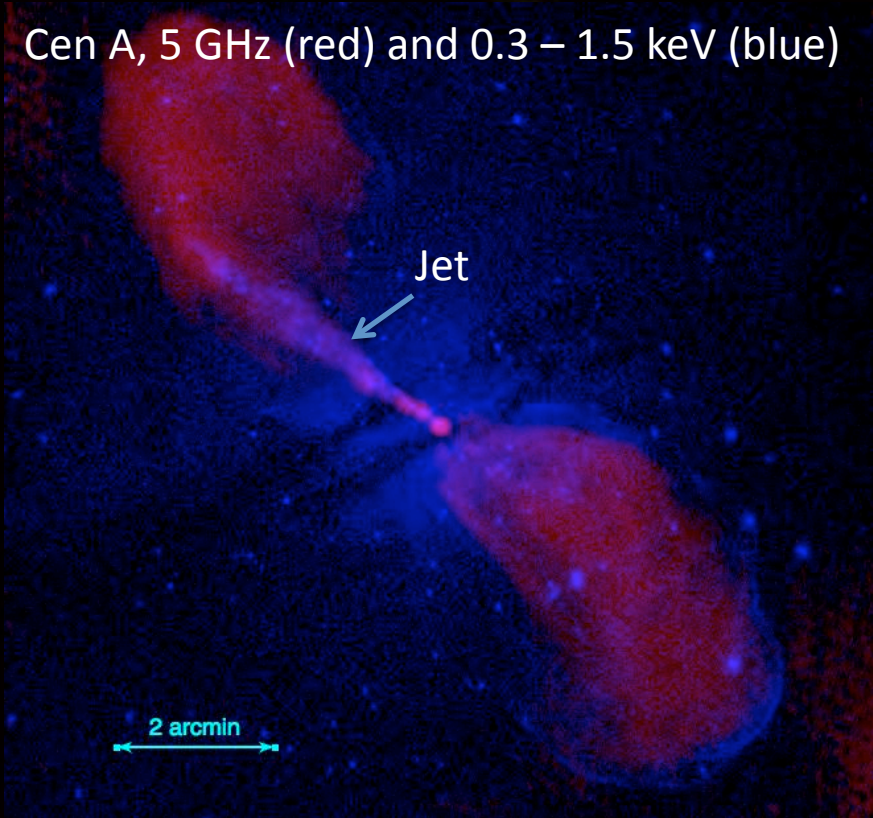
As above, with  $p_{\text{ram}}/p = 10$  and  $\Gamma = 13/9$ :

$\beta_{\text{jet}} = 0.12$ ; mass flow rate through jet =  $3 M_{\odot} \text{ yr}^{-1}$ ;  $n_{e,\text{jet}} = 9.5 \times 10^{-5} \text{ cm}^{-3}$ ; for pressure balance,  $kT_{\text{jet}} = 810 \text{ keV}$  (electrons relativistic, protons not).

Increasing  $P/(pAc)$  or decreasing  $p_{\text{ram}}/p$  will increase flow speed.

Large ram pressure  $\Rightarrow$  large mass flux  $\Rightarrow$  *significant entrainment by the jet.*

Cen A, 5 GHz (red) and 0.3 – 1.5 keV (blue)



## Flow Model for Cen A Jet

Steady, near 1-d flow  $\Rightarrow v = v(R)$ .

Area of cross section,  $A(R)$ , and external pressure,  $p(R)$ , known ( $kT \approx 0.55$  keV,  $n_e \sim r^{-1.26}$ ) – equate to internal pressure (cf. Laing & Bridle 2002).

Mass flow through jet:

$$\dot{M} = \gamma \beta c A \rho$$

Entrainment rate,  $\alpha$ , per unit volume:

$$\dot{M} \Big|_1^2 = \int_1^2 \alpha A dR$$

Assume constant power along jet:  
(enthalpy  $h = \Gamma p / (\Gamma - 1)$ , with  $\Gamma = 13/9$  here).

$$P = (\gamma - 1) \dot{M} c^2 + h A c \beta \gamma^2$$

Momentum flux

$$\Pi = (P/c + \dot{M} c) \beta$$

is affected by buoyancy

$$\Pi \Big|_1^2 = \int_1^2 \frac{dp}{dR} A dR$$



# Flow Parameters

Fiducial values: Jet power,  $P = 6 \times 10^{42} \text{ erg s}^{-1}$  (Croston et al 2009).

Initial speed,  $\beta = 0.7$  (radio knots move at  $\approx 0.5c$  near inner end; Hardcastle et al 2003).

Stellar mass loss: star density =  $f \times$  gravitating mass density (hydrostatic equilibrium), with  $f \approx 1$  at  $R = 100 \text{ arcsec}$  (1.8 kpc; consistent with photometry) and  $\alpha = f \rho_{\text{grav}} / \tau$ , with  $\tau = 10^{12} \text{ yr}$  (Faber & Gallagher 1976).

Other entrainment – let  $f$  vary.

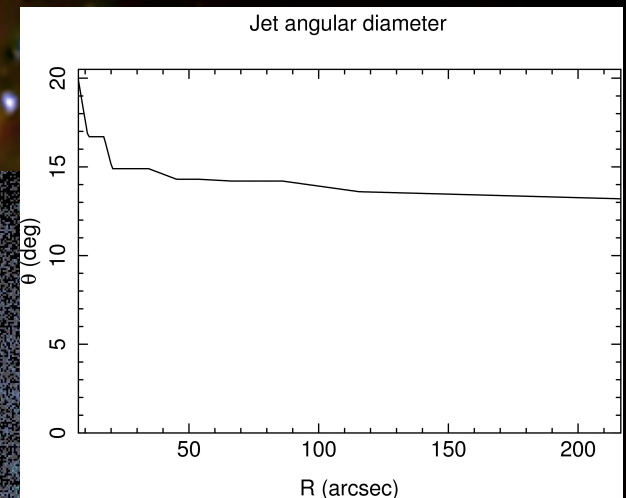
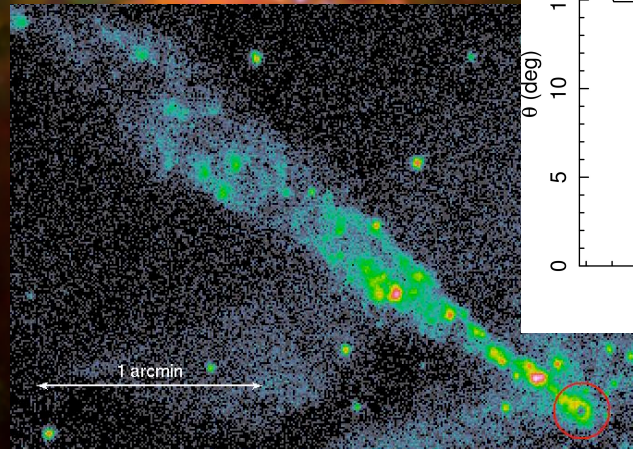
Radio and X-ray measurements of jet width agree.

No dissipation (constant  $\dot{M}$ )

P eqn: speed must decrease

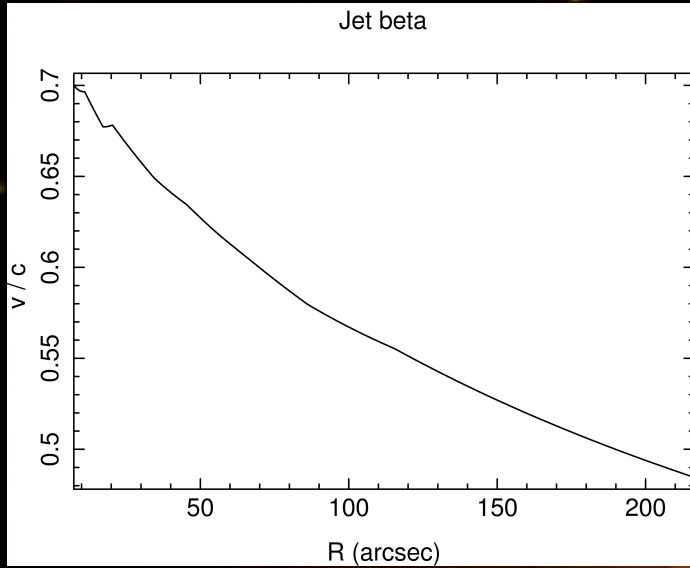
$\Pi$  eqn: speed must increase

**=> Flow is dissipative**



Flow is over-determined: adjust  $f$  to make 3 equations consistent.

# Fiducial Model

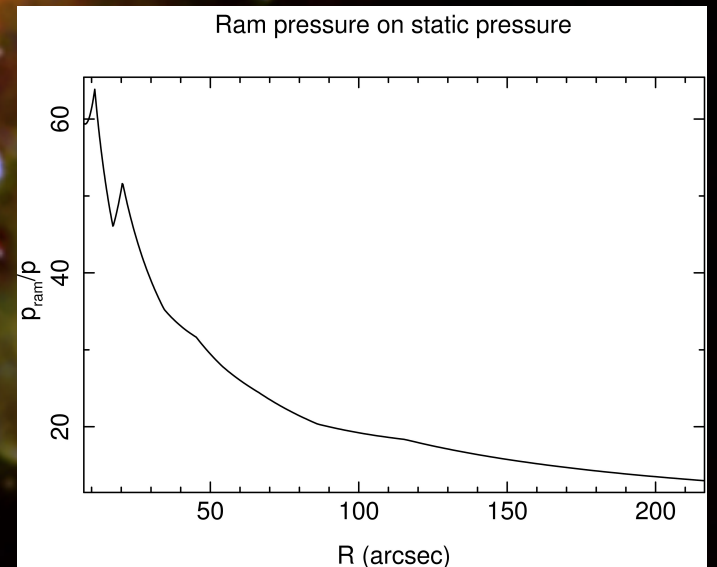
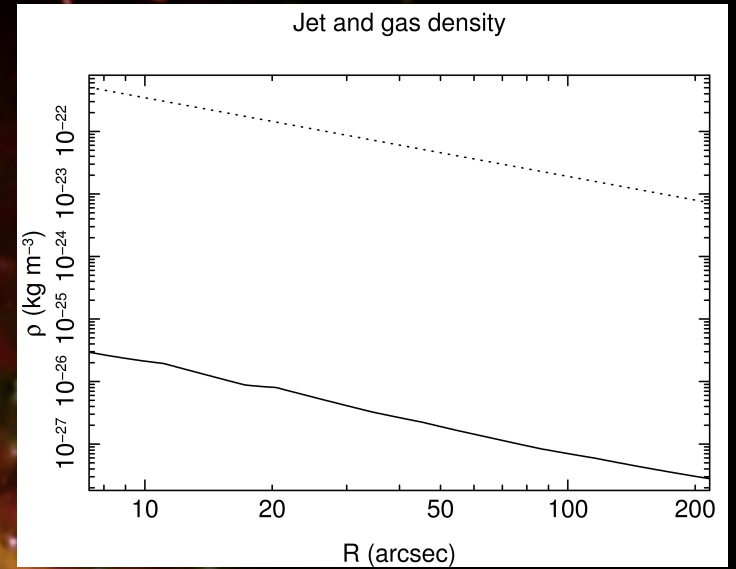
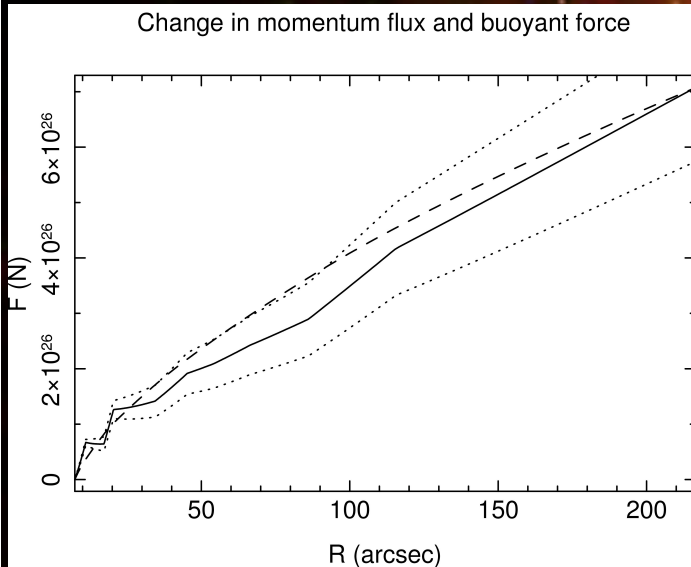


Solution determined by 2 equations, energy & mass flow.

Choose  $f$  to make  $\Delta\Pi =$  buoyant force

$\Rightarrow f = 0.60$

Insensitive to flow parameters



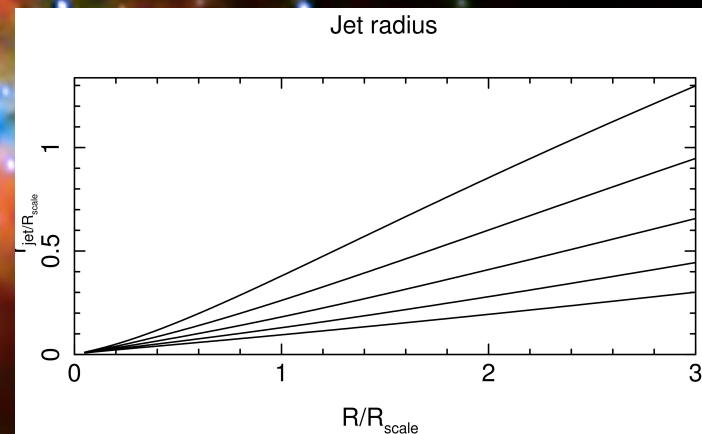
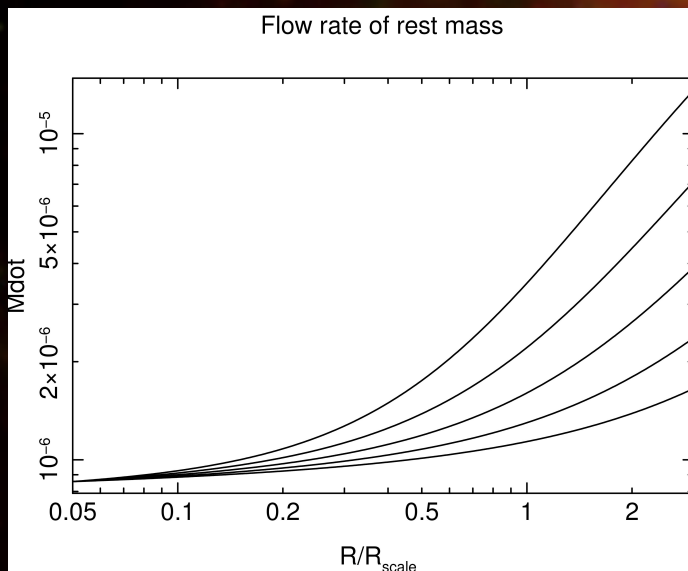
# Effect of Environment

Dissipation due mass entrainment makes a jet unstable:

larger cross section => more entrainment => more dissipation => jet broadens

Model: 
$$\frac{d\dot{M}}{dR} = \alpha A; \quad \frac{d\Pi}{dR} = -A \frac{dp}{dR}; \quad P = \text{constant}$$

Same power, initial speed, mass injection ( $\alpha$ ) as fiducial model, but bounding pressure is scaled by down factors of up to 2:



Moderate asymmetry in environment can make the difference between a jet and a lobe.

# Effect of Environment

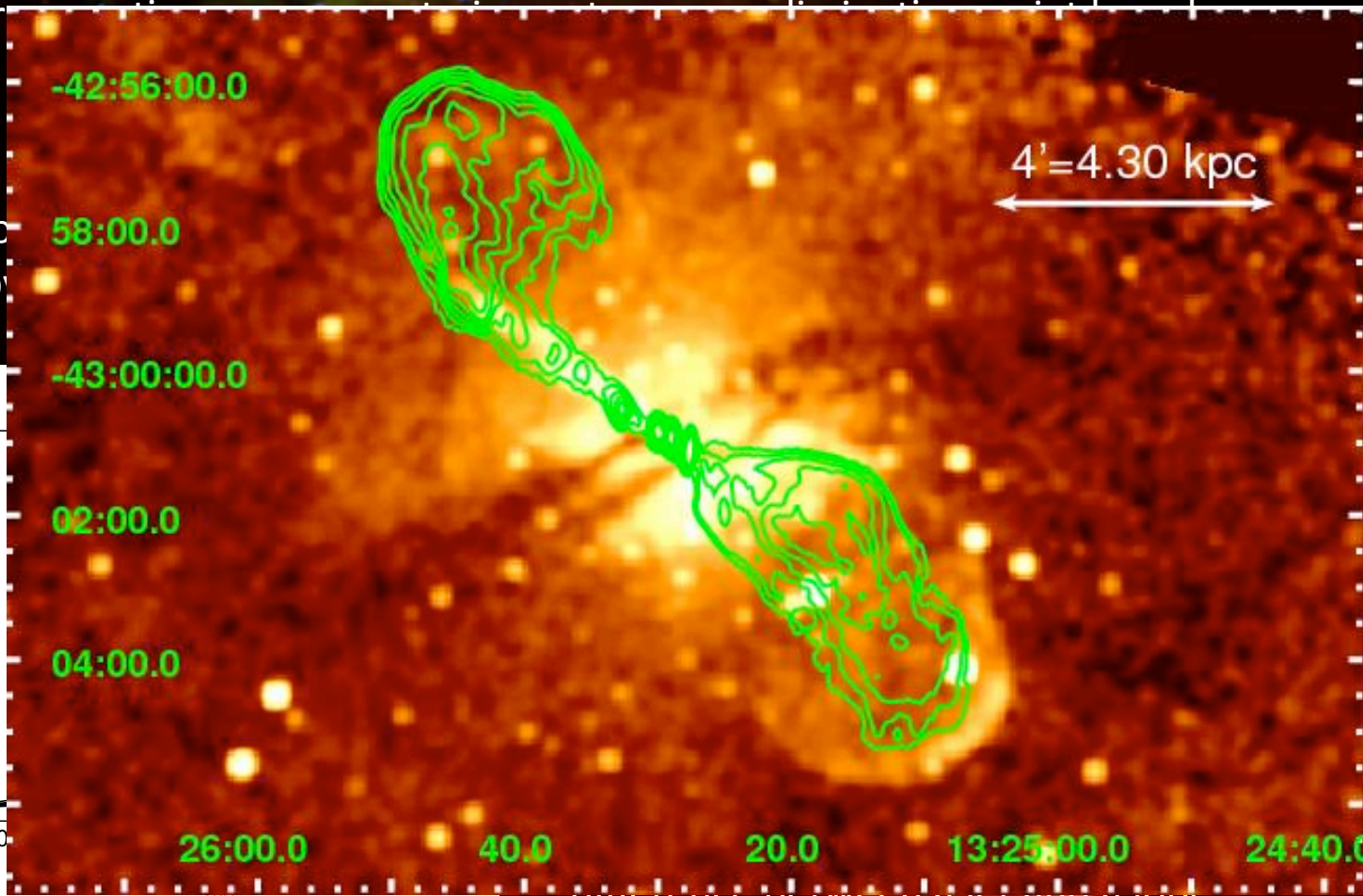
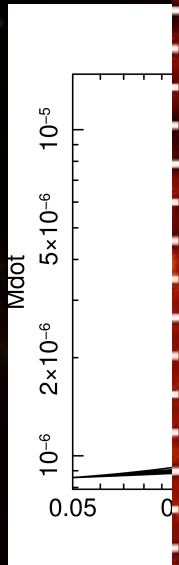
Dissipation due mass entrainment makes a jet unstable:

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Model:

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difference between a jet and a lobe:

# Conclusions

- Appears to be substantial entrainment in the jets of both Cyg A (FR II) and Cen A (FR I)
- Dissipation due to entrainment places the Cen A jet close to the margin for rapid inflation
- - fate of the jet is sensitive to environmental influence
- Details of jet physics affect the site and manner of energy deposition by AGN jets