

# (Inverse) Higgs Phenomenology at LHC

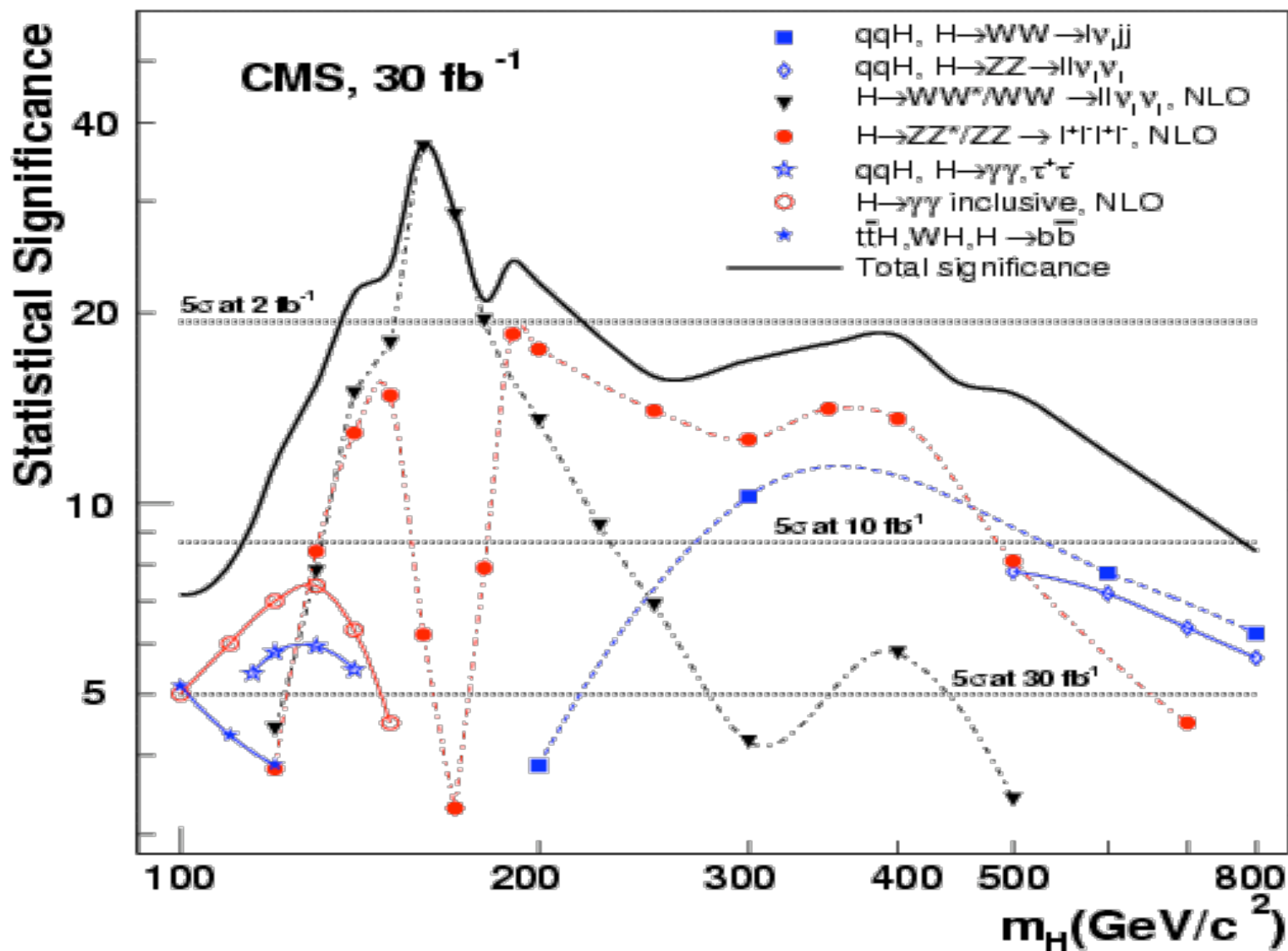
Kazuhiro Tobe (Michigan State University)

“LHC Inverse Problem” workshop @ Ann Arbor

This talk is based on collaborations and discussions with  
Alexander Belyaev, Chuan-Ren Chen,  
Daisuke Nomura and C.-P. Yuan

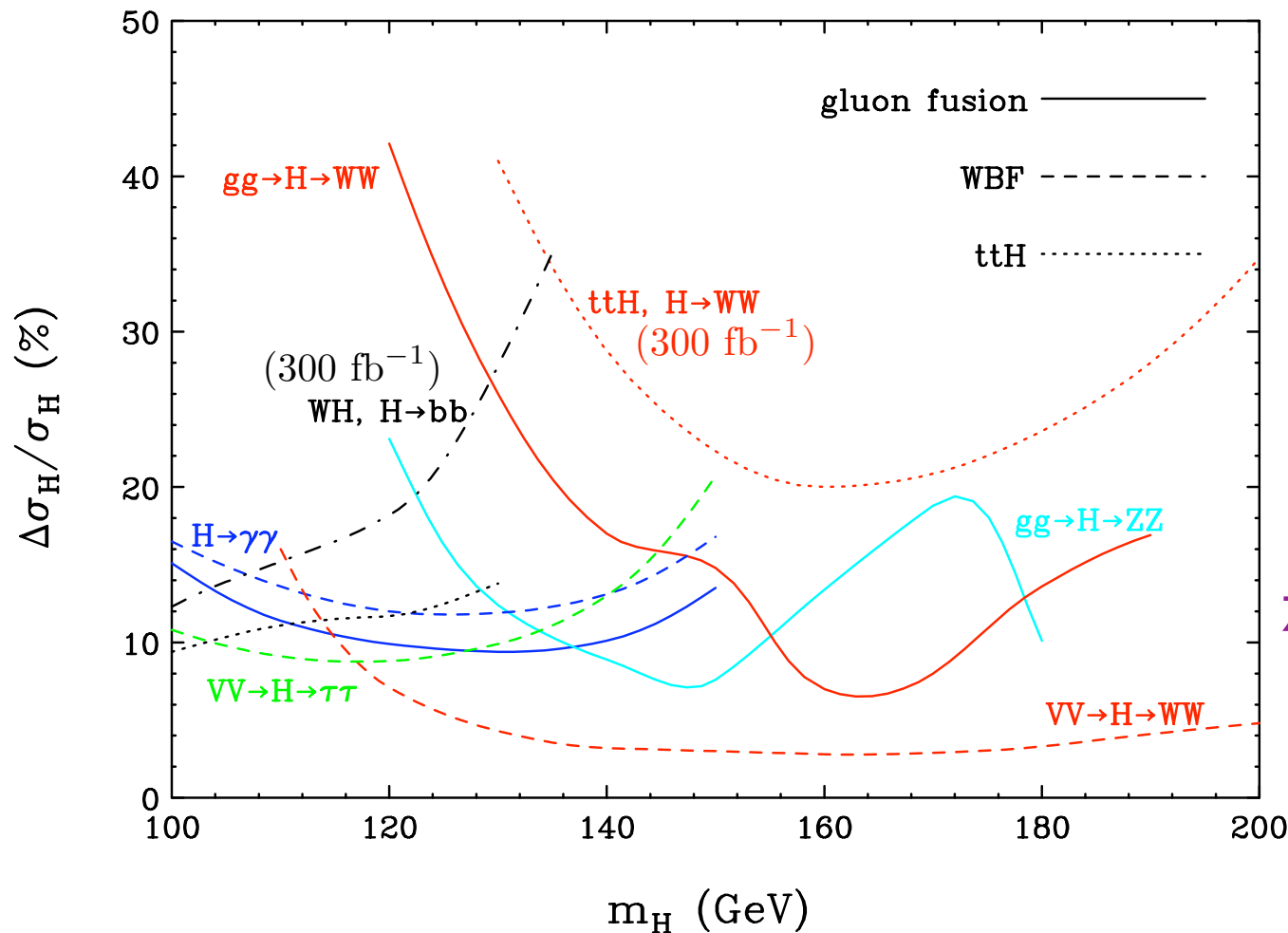
# ◆ Introduction

LHC experiments have the great potential to discover a Higgs boson in various channels.



In physics beyond the SM (SUSY, Little Higgs etc), we expect some non-SM Higgs couplings.

Can we extract any information on new physics from Higgs data?



Expected relative error on the determination of  $\sigma \times \text{BR}$  for various Higgs search channels at the LHC with  $200 \text{ fb}^{-1}$  of data

Zeppenfeld hep-ph/0203123

Zeppenfeld et al, hep-ph/0002036

Duhrssen et al, hep-ph/0406323

# Outline

## ◆ Introduction

## ◆ Higgs interactions in new physics (NP) models

MSSM, Littlest Higgs model with T-parity

## ◆ Constraint on NP parameters from Higgs data

$\chi^2$  analysis using Higgs  $\sigma \times \text{BR}$  data

- Littlest Higgs model with T-parity

new particle mass scale  $f$

- MSSM  $\beta, \alpha, \dots$

(very preliminary)

## ◆ Summary

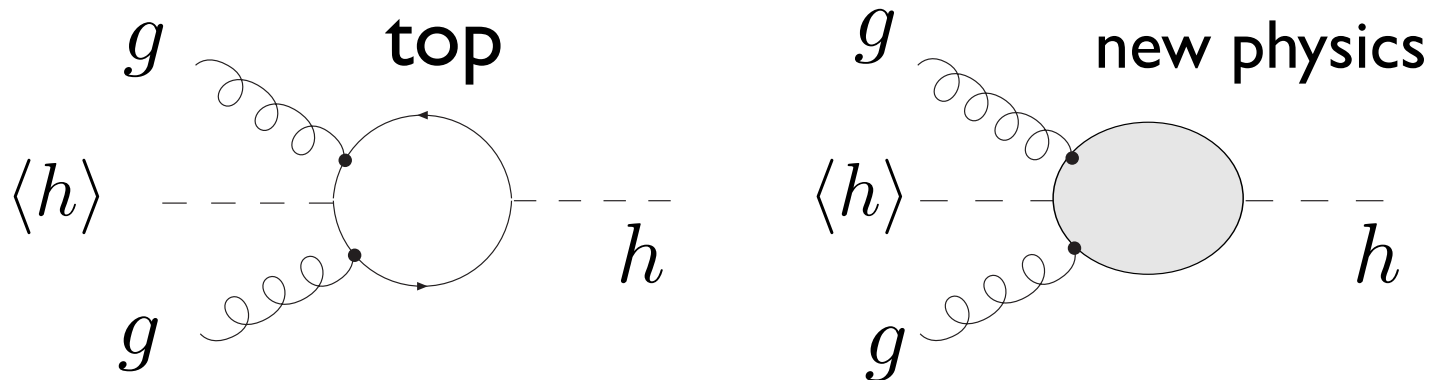
# ◆ Higgs interactions in new physics models

- Solution to “naturalness problem” of Higgs mass

If there is “new physics” which cancels the  $\Lambda^2$  induced by top,



the “new Physics” can affect gluon fusion process.

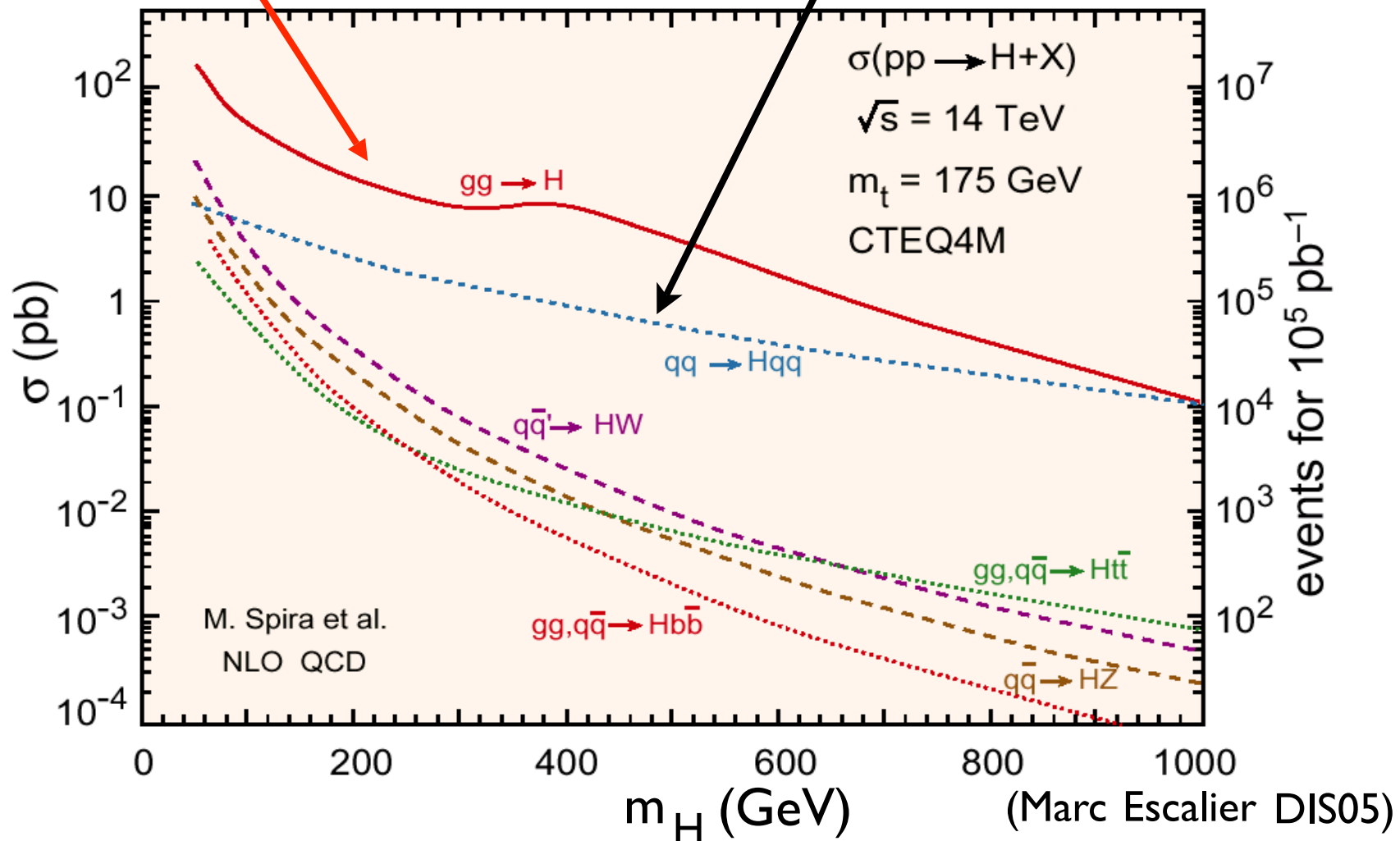
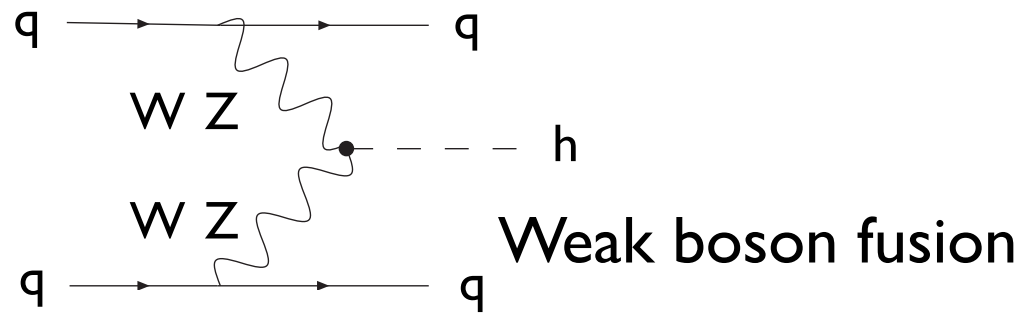
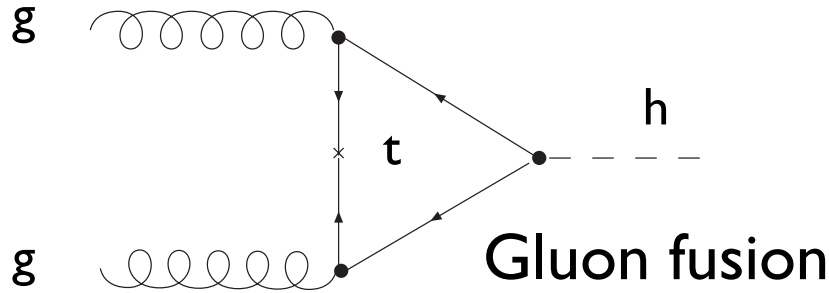


MSSM --- light stop (gluophobic scenario)

Little Higgs --- fermionic partner of top etc

It can affect the SM like Higgs boson search at LHC

# Note: SM Higgs production at LHC



- Higgs sector can be modified

## MSSM ---- two Higgs doublet model

tree level couplings of the lightest Higgs is modified

e.g.  $\frac{g_{hdd}}{g_{hdd}^{\text{SM}}} = -\frac{\sin \alpha}{\cos \beta}$  it can be larger one  
 $g_{hdd}/g_{hdd}^{\text{SM}} \simeq 1.5$  for  $m_h$  max scenario

Higgs production and decay can be very different from the SM prediction

## Higgs interactions

gauge bosons and fermions

$$\mathcal{L}_{eff} = -g_{hVV} h V_{\mu}^{\dagger} V^{\mu} - g_{hff} h \bar{f} f$$

$$g_{hVV}/g_{hVV}^{\text{SM}} = \gamma_V(A, B, \dots), \quad g_{hff}/g_{hff}^{\text{SM}} = \gamma_f(A, B, \dots)$$

gluon and photon

$$\Gamma_{gg}/\Gamma_{gg}^{\text{SM}} = 1 + \delta_{gg}(A, B, \dots), \quad \Gamma_{\gamma\gamma}/\Gamma_{\gamma\gamma}^{\text{SM}} = 1 + \delta_{\gamma\gamma}(A, B, \dots)$$

A,B,... are parameters of the new physics model

Higgs production and decay data constrain (hopefully measure)  
 the parameters of the theory (A,B,...)

# Example I: MSSM

## MSSM lightest Higgs boson couplings

- gauge boson

$$\frac{g_{hVV}}{g_{hVV}^{\text{SM}}} = \sin(\beta - \alpha) \text{ for } V = W, Z$$

$$\tan \beta = v_u/v_d$$

$\alpha$  : Higgs mixing angle

- fermion

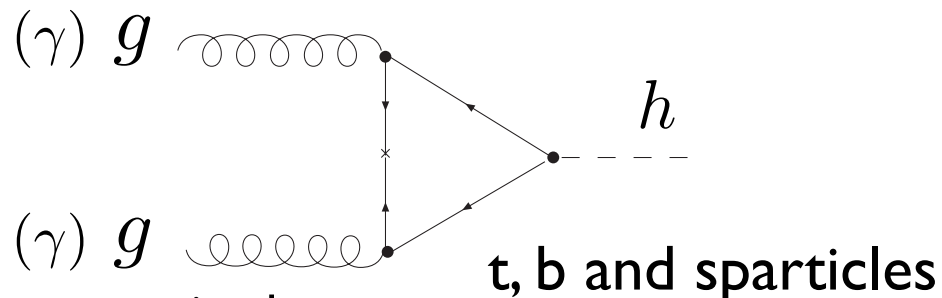
$$\frac{g_{h uu}}{g_{h uu}^{\text{SM}}} = \frac{\cos \alpha}{\sin \beta} \qquad \frac{g_{h dd}}{g_{h dd}^{\text{SM}}} = -\frac{\sin \alpha}{\cos \beta}$$

$$\frac{g_{h bb}}{g_{h bb}^{\text{SM}}} = -\frac{\sin \alpha}{\cos \beta} (1 + \delta_b) \text{ for bottom} \qquad (b \rightarrow \tau \text{ for tau})$$

$$\frac{g_{h tt}}{g_{h tt}^{\text{SM}}} = \frac{\cos \alpha}{\sin \beta} (1 + \delta_t) \text{ for top}$$

- gluon and photon

$$\frac{\Gamma_{gg}}{\Gamma_{gg}^{\text{SM}}} = \frac{\sigma_{gg}}{\sigma_{gg}^{\text{SM}}} = 1 + \delta_{gg}$$



Similarly, photon-Higgs coupling is parametrized

$\alpha, \beta, \delta_p$  ( $p = b, t, \tau, gg$  and  $\gamma\gamma$ ) are (effective) parameters for Higgs couplings



## Example 2: Littlest Higgs model with T-parity

- $SU(5)/SO(5)$  non-linear sigma model

Arkani-Hamed, Cohen, Katz, Nelson hep-ph/0206021

$$SU(5) \xrightarrow{f} SO(5) \quad \mathbf{1}_0 \oplus \mathbf{3}_0 \oplus \mathbf{2}_{\pm 1/2} \oplus \mathbf{3}_{\pm 1}$$

$$[SU(2) \times U(1)]_1 \times [SU(2) \times U(1)]_2 \longrightarrow SU(2) \times U(1)$$

- Collective symmetry breaking mechanism

Large SM one-loop quadratic divergences in Higgs mass  
are canceled by new particles

- T-parity Cheng, Low hep-ph/0308199

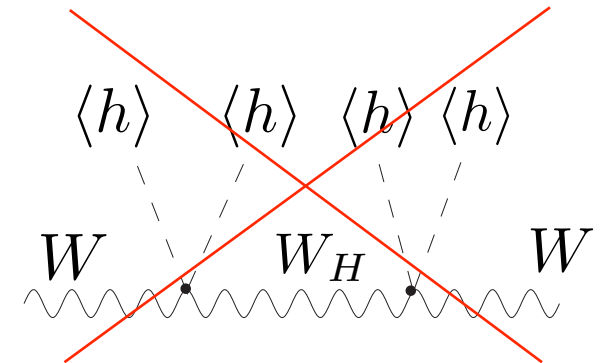
$$SU(2)_1 \times U(1)_1 \leftrightarrow SU(2)_2 \times U(1)_2$$

SM particles  $\rightarrow$  +SM particles

$$(W_H, Z_H, A_H, \Phi) \rightarrow -(W_H, Z_H, A_H, \Phi)$$

(heavy gauge bosons and Higgs)

$$q_- \rightarrow -q_- \quad (\text{T-odd doublet fermions})$$



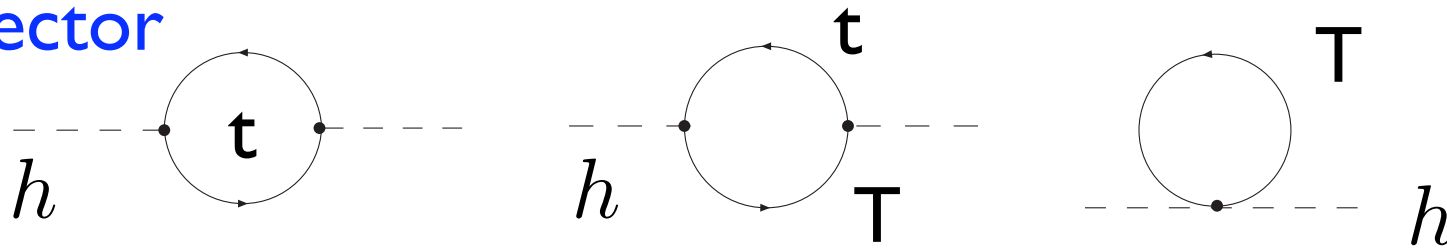
- \* Contributions to EW observables are loop suppressed.

The new particle scale  $f$  can be much lower than 1 TeV.

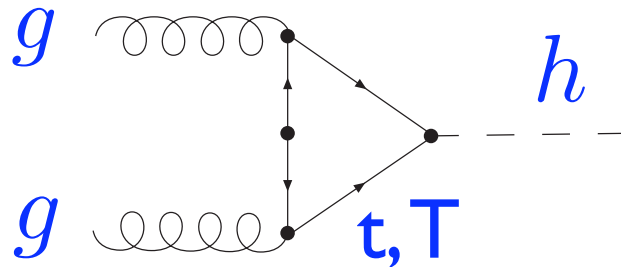
- \* The lightest T-odd particle can be a good dark matter candidate

# Gluon fusion process in Littlest Higgs model with T-parity

## Top sector



The large quadratic divergence induced by top is canceled by the heavy fermionic partner of the top-quark, T

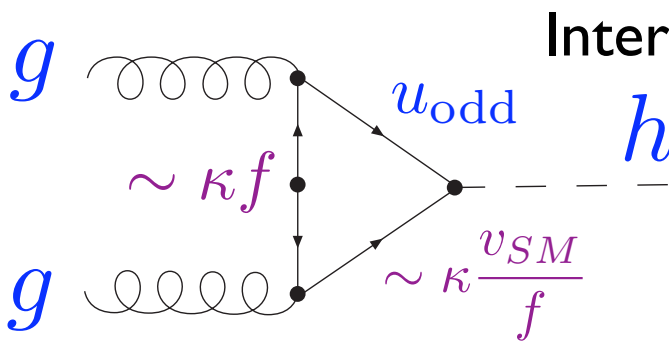


$$\frac{\delta\sigma_{gg\rightarrow h}(\text{top sector})}{\sigma_{gg\rightarrow h}(\text{top in SM})} \simeq -\frac{3}{2} \frac{v_{SM}^2}{f^2}$$

Han, Logan, McElrath, Wang hep-ph/0302188

## T-odd doublet quark sector

Low, hep-ph/0409025, Hubisz et al hep-ph/0411264



Interaction which generates T-odd doublet quarks induces the Higgs interaction

$$\frac{\delta\sigma_{gg\rightarrow h}(\text{T odd fermions})}{\sigma_{gg\rightarrow h}(\text{top in SM})} \simeq -\frac{3}{2} \frac{v_{SM}^2}{f^2}$$

Chen, Tobe, Yuan, hep-ph/0602211

# Gluon fusion

$$\frac{\delta\sigma_{gg\rightarrow h}}{\sigma_{gg\rightarrow h}^{\text{SM}}} \simeq -3\frac{v_{SM}^2}{f^2} \simeq \begin{cases} -37\% & \text{for } f = 700 \text{ GeV,} \\ -18\% & \text{for } f = 1000 \text{ GeV.} \end{cases}$$

it could be significantly suppressed

## Higgs couplings in Littlest Higgs model with T-parity

- gauge boson (V=W, Z)

$$\frac{g_{hVV}}{g_{hVV}^{\text{SM}}} \simeq 1 - \frac{1}{4}\frac{v_{SM}^2}{f^2} \simeq \begin{cases} 0.97 & \text{for } f = 700 \text{ GeV,} \\ 0.98 & \text{for } f = 1000 \text{ GeV.} \end{cases}$$

- fermion

$$\frac{g_{h uu}}{g_{h uu}^{\text{SM}}} \simeq 1 - \frac{3}{4}\frac{v_{SM}^2}{f^2} \simeq \begin{cases} 0.90 & \text{for } f = 700 \text{ GeV,} \\ 0.95 & \text{for } f = 1000 \text{ GeV.} \end{cases}$$

$$\frac{g_{h tt}}{g_{h tt}^{\text{SM}}} \simeq 1 - \frac{1}{2}\frac{v_{SM}^2}{f^2} \simeq \begin{cases} 0.94 & \text{for } f = 700 \text{ GeV,} \\ 0.97 & \text{for } f = 1000 \text{ GeV.} \end{cases} \quad (\lambda_1 = \lambda_2)$$

$$\frac{g_{h dd}}{g_{h dd}^{\text{SM}}} \simeq 1 - \frac{1}{4}\frac{v_{SM}^2}{f^2} \simeq \begin{cases} 0.97 & \text{for } f = 700 \text{ GeV,} \\ 0.99 & \text{for } f = 1000 \text{ GeV,} \end{cases} \quad \text{for Case A,}$$

$$\simeq 1 - \frac{5}{4}\frac{v_{SM}^2}{f^2} \simeq \begin{cases} 0.84 & \text{for } f = 700 \text{ GeV,} \\ 0.92 & \text{for } f = 1000 \text{ GeV,} \end{cases} \quad \text{for Case B.}$$

(We consider the same Yukawa structures in lepton sector, as in quark sector.)

Higgs couplings are function of the scale “f”

$$R_{\sigma(X)} = \frac{\sigma^{\text{LH}}(X)}{\sigma^{\text{SM}}(X)} \quad R_{\text{BR}(Y)} = \frac{\text{BR}^{\text{LH}}(Y)}{\text{BR}^{\text{SM}}(Y)}$$

$$R_{\sigma(X)} \times R_{\text{BR}(Y)} \quad \text{for } f = (600, 700, 1000) \text{ GeV}$$

$m_h = 120 \text{ GeV}$	$R_{\text{BR}(\gamma\gamma)}$	$R_{\text{BR}(\tau\tau)}$	$R_{\text{BR}(b\bar{b})}$	$R_{\text{BR}(VV)}$
$R_{\sigma(gg)}$ (Case A)	0.57, 0.68, 0.84	0.56, 0.67, 0.83	—	0.55, 0.66, 0.83
(Case B)	0.81, 0.86, 0.93	0.51, 0.63, 0.81	—	0.78, 0.84, 0.92
$R_{\sigma(VV)}$ (Case A)	0.97, 0.98, 0.99	0.95, 0.96, 0.98	—	0.94, 0.96, 0.98
(Case B)	1.34, 1.22, 1.09	0.84, 0.89, 0.95	—	1.30, 1.19, 1.08
$R_{\sigma(t\bar{t}h)}$ (Case A)	—	0.87, 0.90, 0.95	0.87, 0.90, 0.95	—
(Case B)	—	0.77, 0.83, 0.92	0.77, 0.83, 0.92	—
$R_{\sigma(Vh)}$ (Case A)	0.97, 0.98, 0.99	—	0.95, 0.96, 0.98	—
(Case B)	1.34, 1.22, 1.09	—	0.84, 0.89, 0.95	—
$m_h = 200 \text{ GeV}$	$R_{\text{BR}(\gamma\gamma)}$	$R_{\text{BR}(\tau\tau)}$	$R_{\text{BR}(b\bar{b})}$	$R_{\text{BR}(VV)}$
$R_{\sigma(gg)}$ (Case A)	—	—	—	0.55, 0.67, 0.83
(Case B)	—	—	—	0.56, 0.67, 0.83
$R_{\sigma(VV)}$ (Case A)	—	—	—	0.90, 0.94, 0.97
(Case B)	—	—	—	0.90, 0.94, 0.97

- Higgs production via gluon fusion is suppressed.
- $\gamma\gamma$ ,  $VV$  decay modes via weak boson fusion can be enhanced in small Higgs mass region in Case B.

# ◆ Constraint on NP parameters from Higgs data (preliminary)

- Suppose LHC experiments find a Higgs in various channels and measure  $\sigma \times \text{BR}$
- The following errors on the determination of  $\sigma_h = \sigma \times \text{BR}$  for various Higgs search channels are assumed: [Zeppenfeld, hep-ph/0203123](#)

$\Delta\sigma_h/\sigma_h(\%)$	$\gamma\gamma$ (decay)	$\tau\tau$	$bb$	$VV$
$gg$ (production)	9.9, 9.4	-	-	23.1, 12.4 ( $ZZ^*$ ) 42.1, 26.0 ( $WW^*$ )
$qqh$	12.0, 11.9	8.8, 9.9	-	7.1, 4.3 ( $WW^*$ )
$tth$	-	17.0, 23.3	11.4, 14.3	-, 42.0 ( $WW^*$ )
$Wh$	-	-	19.0, 25.0	

for  $m_h = 120$  and  $130$  GeV, respectively

$gg$	10%
$qqh$	4%
$tth$	15%
$WH$	7%

Theoretical QCD and PDF uncertainties on the various Higgs production channels.

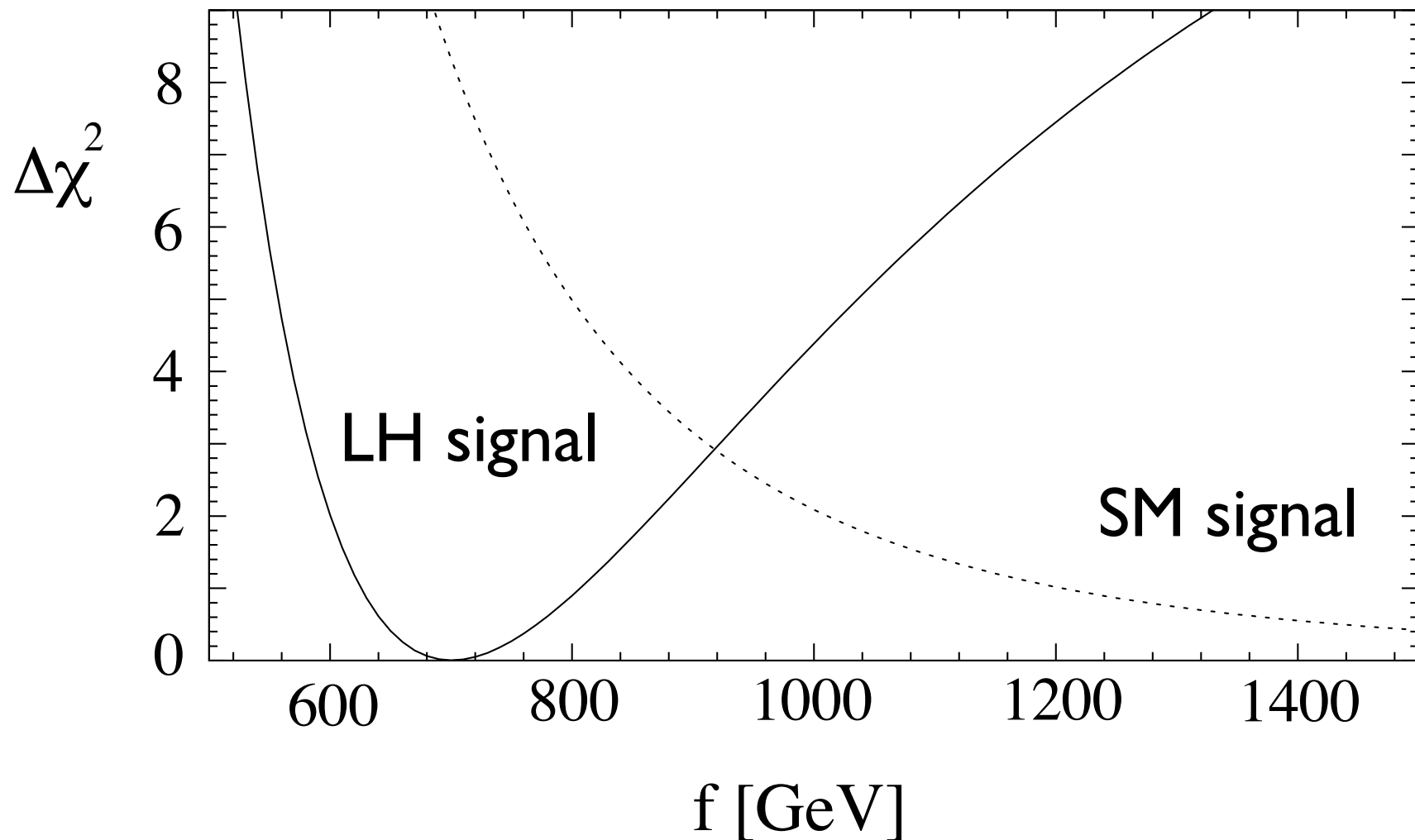
[Duhrssen et al, hep-ph/0406323](#)

[Anastasiou et al, hep-ph/0509014](#)

- $\chi^2$  analysis on  $\sigma_h = \sigma \times \text{BR}$
- Constraints on parameters of the models

# Littlest Higgs model with T-parity

$$m_h = 120 \text{ GeV}$$

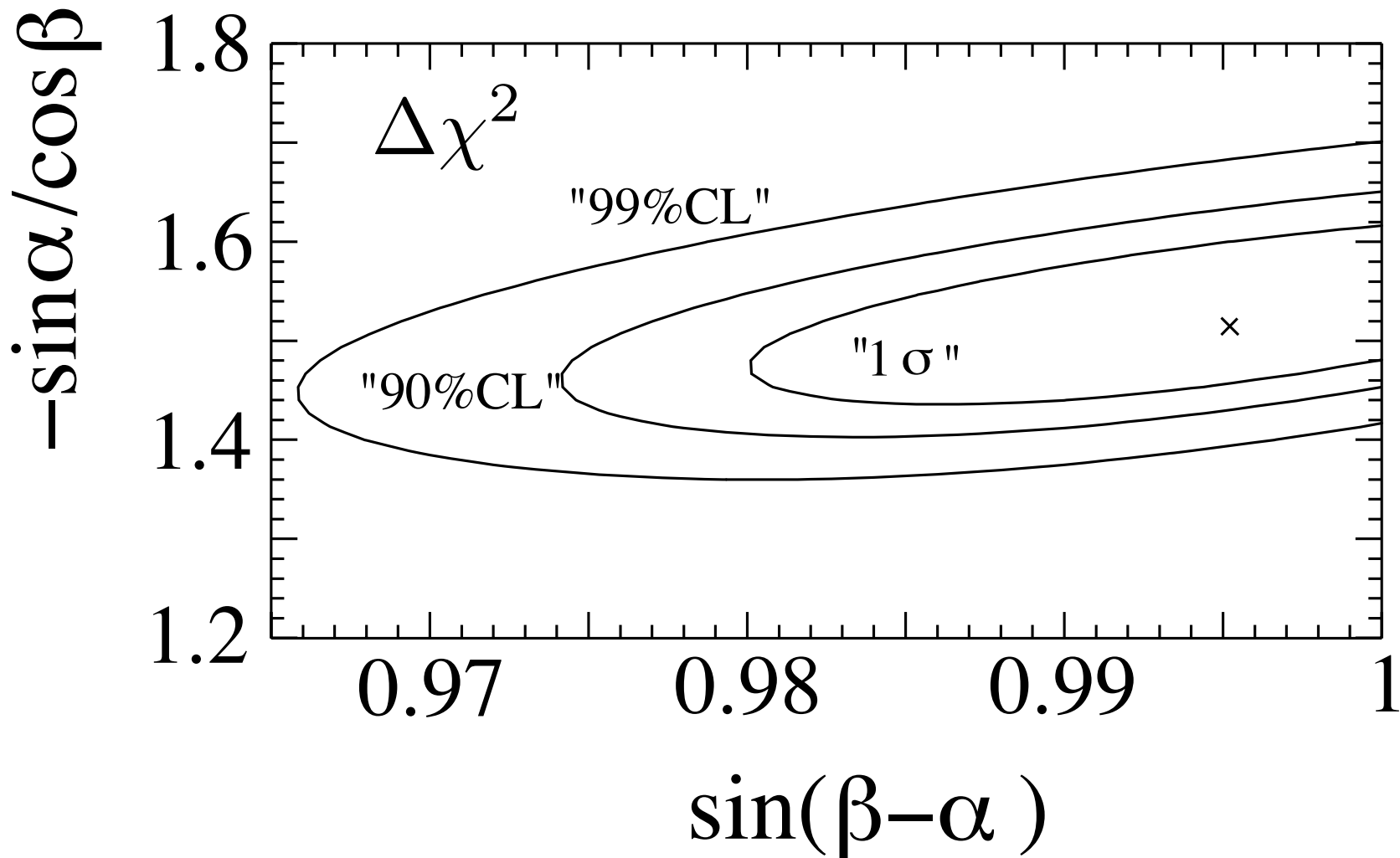


LH signal : Littlest Higgs in Case A with  $f = 700$  GeV

MSSM

For simplicity,

$$\delta_b = \delta_t = \delta_{gg} = \delta_{\gamma\gamma} = 0$$



MSSM signal: mh max scenario with

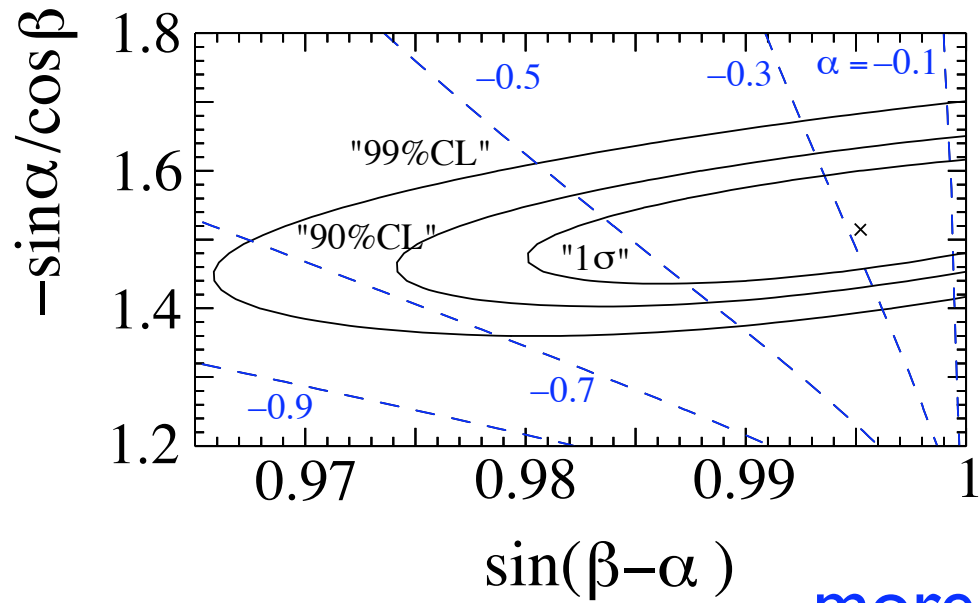
$$\tan\beta = 8 \text{ and } m_A = 250 \text{ GeV}$$

$\sigma_h/\sigma_h^{\text{SM}}$	$\gamma\gamma$ (decay)	$\tau\tau$	$b\bar{b}$	$VV$
$gg$ (production)	0.52	-	-	0.54 ( $ZZ^*$ ) 0.55 ( $WW^*$ )
$qqh$	0.54	1.3	-	0.57 ( $WW^*$ )
$tth$	-	1.3	1.3	0.56 ( $WW^*$ )
$Wh$	-	-	1.27	

FeynHiggs

Table 1:  $\sigma \times \text{BR}/(\sigma \times \text{BR})_{SM}$  in  $m_h$  max scenario with  $\tan\beta = 8$  and  $m_A = 250$  GeV.

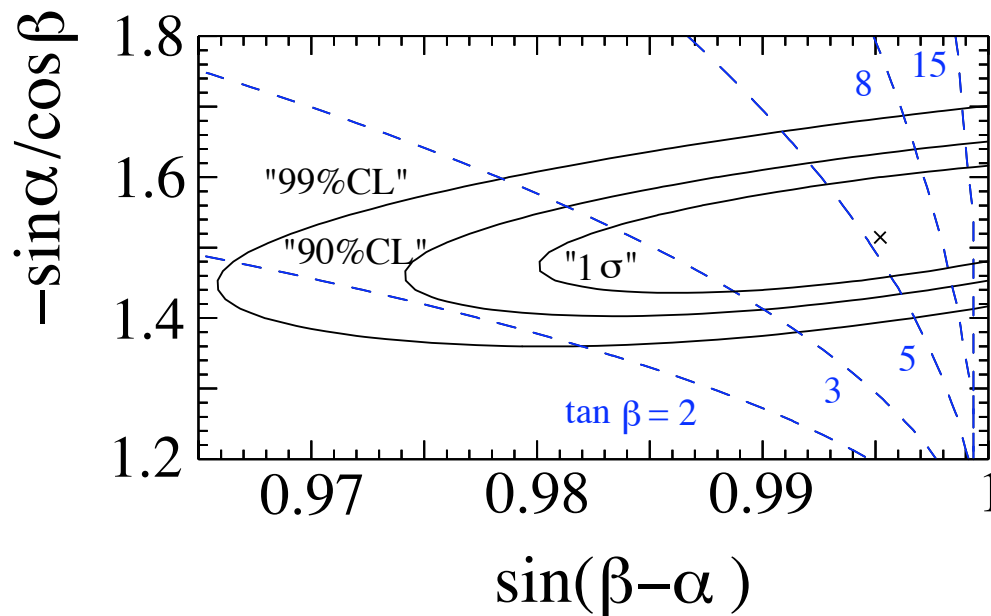




$\tan\beta$  and  $\alpha$



more fundamental parameters ?



It's difficult to  
constrain  $\tan\beta$

## ◆ Summary

- LHC experiments have the great potential to discover a Higgs in various channels.
- We expect non-SM Higgs interaction in interesting new physics model (MSSM, Little Higgs etc).
- Higgs production and decay data might provide interesting information on new physics models.

Higgs physics at the LHC can be indirect search for new physics