

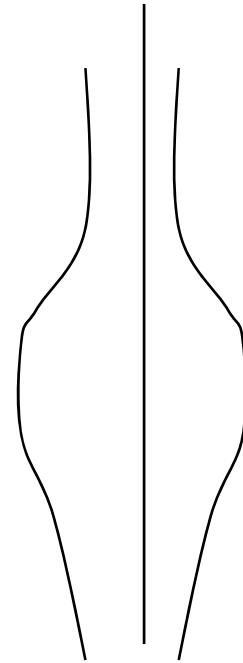
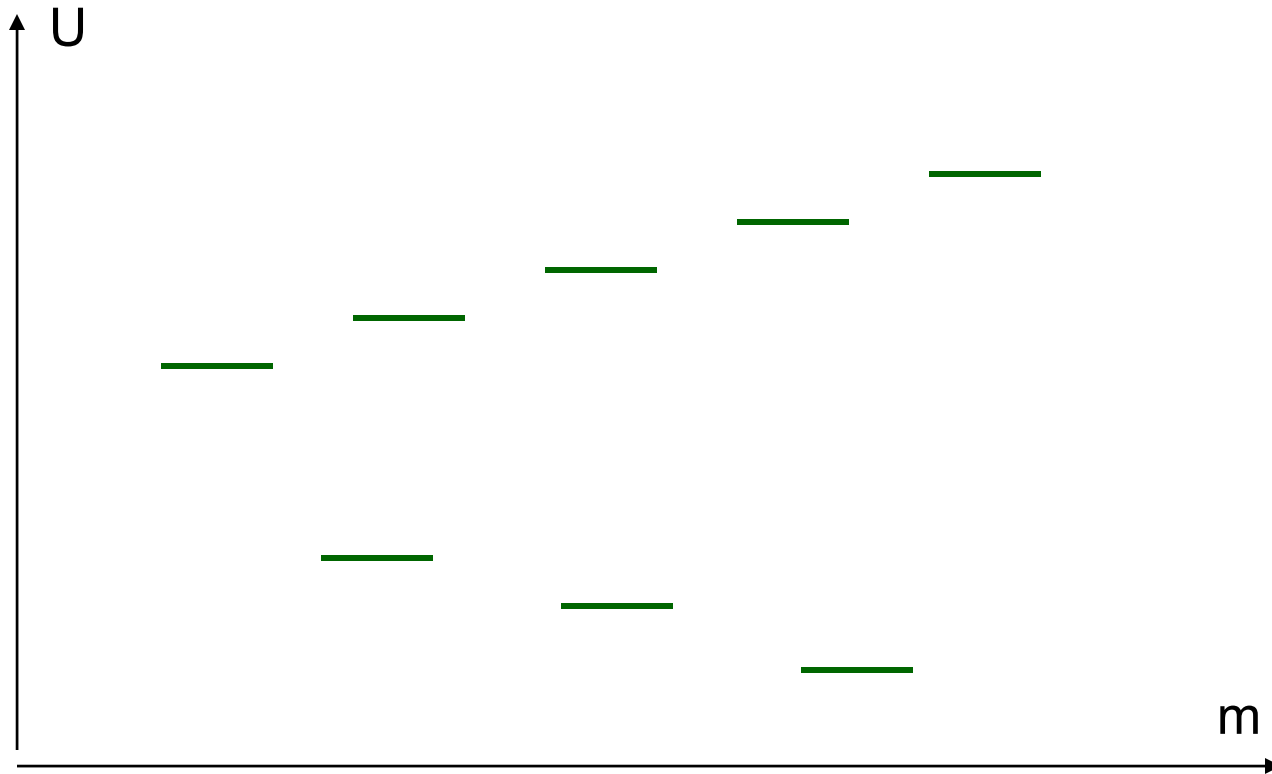
Traps for precision measurement

Long coherence times= good.

Traps for precision measurement

Magnetic trap (Ioffe-Pritchard) for neutrals

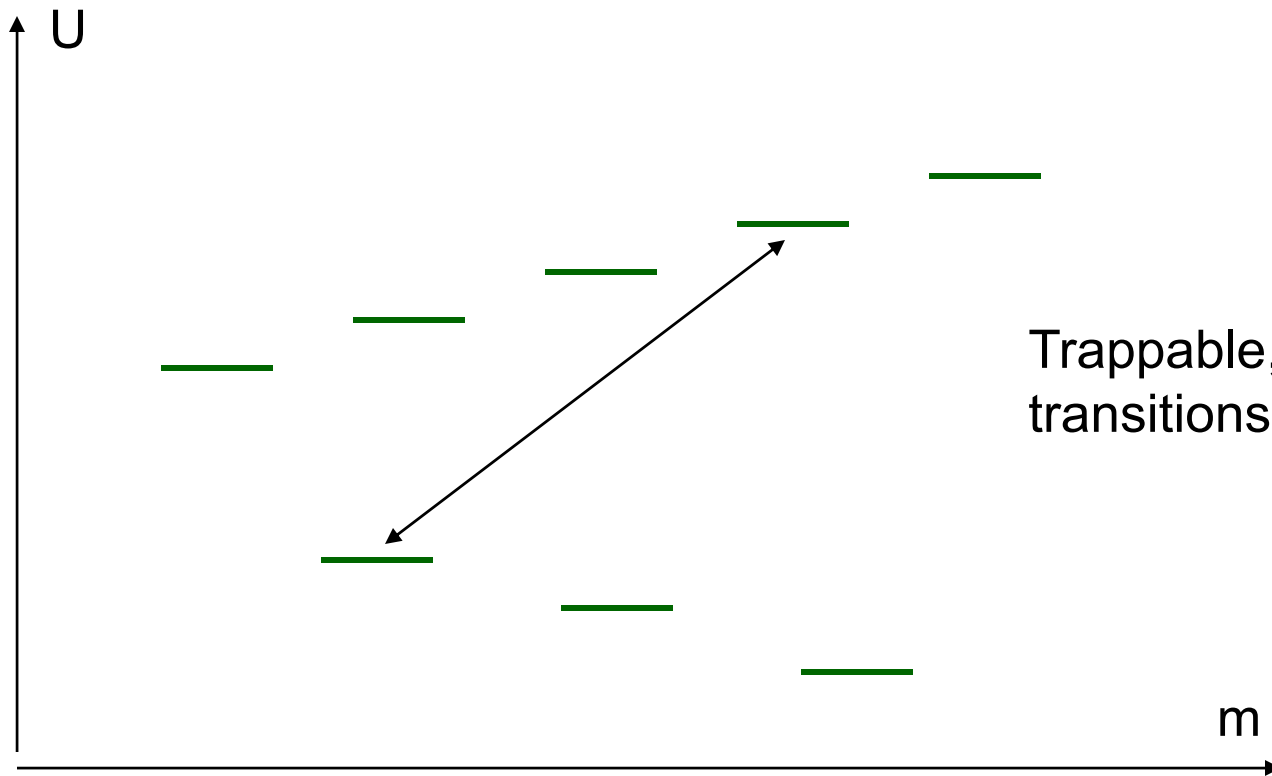
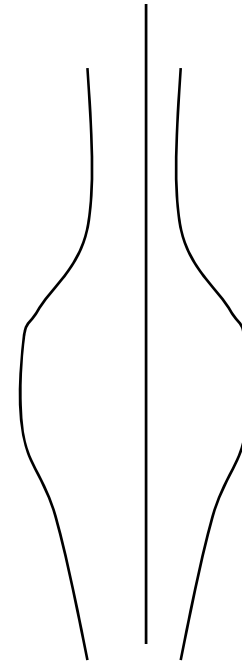
$$\text{Force} = \text{grad} (|B| \mu)$$



Traps for precision measurement

Magnetic trap (Ioffe-Pritchard) for neutrals

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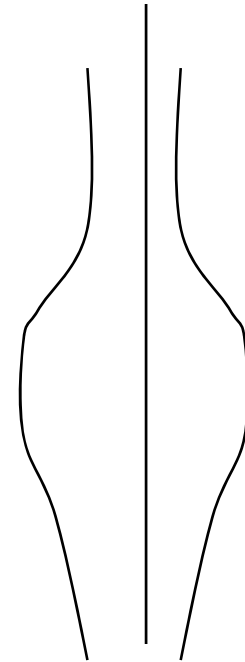
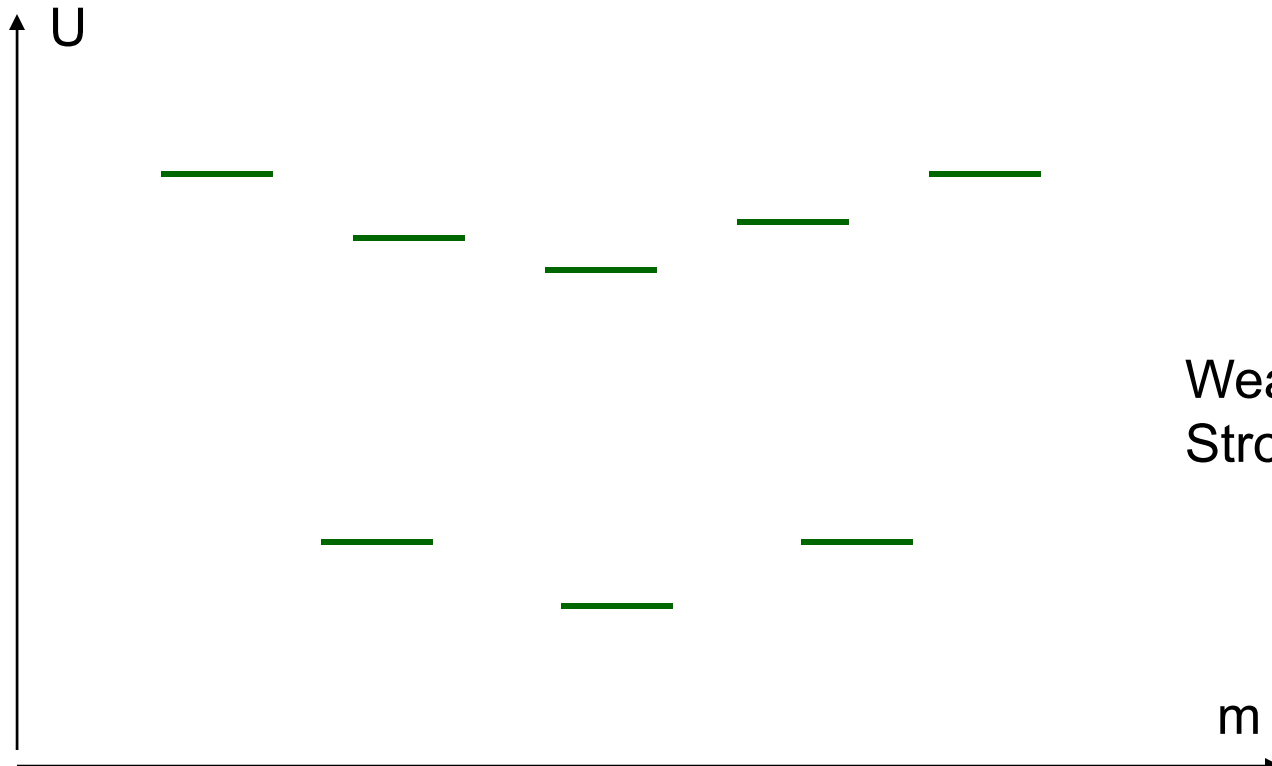


Trappable, B-field insensitive transitions exist, in most species.

Traps for precision measurement

Electrostatic trap (Ioffe-Pritchard) for neutrals

$$\text{Force} = \text{grad} (|E| d(E))$$

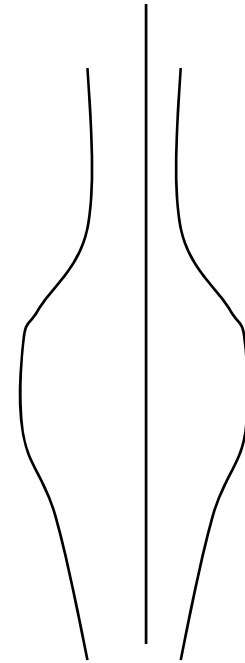
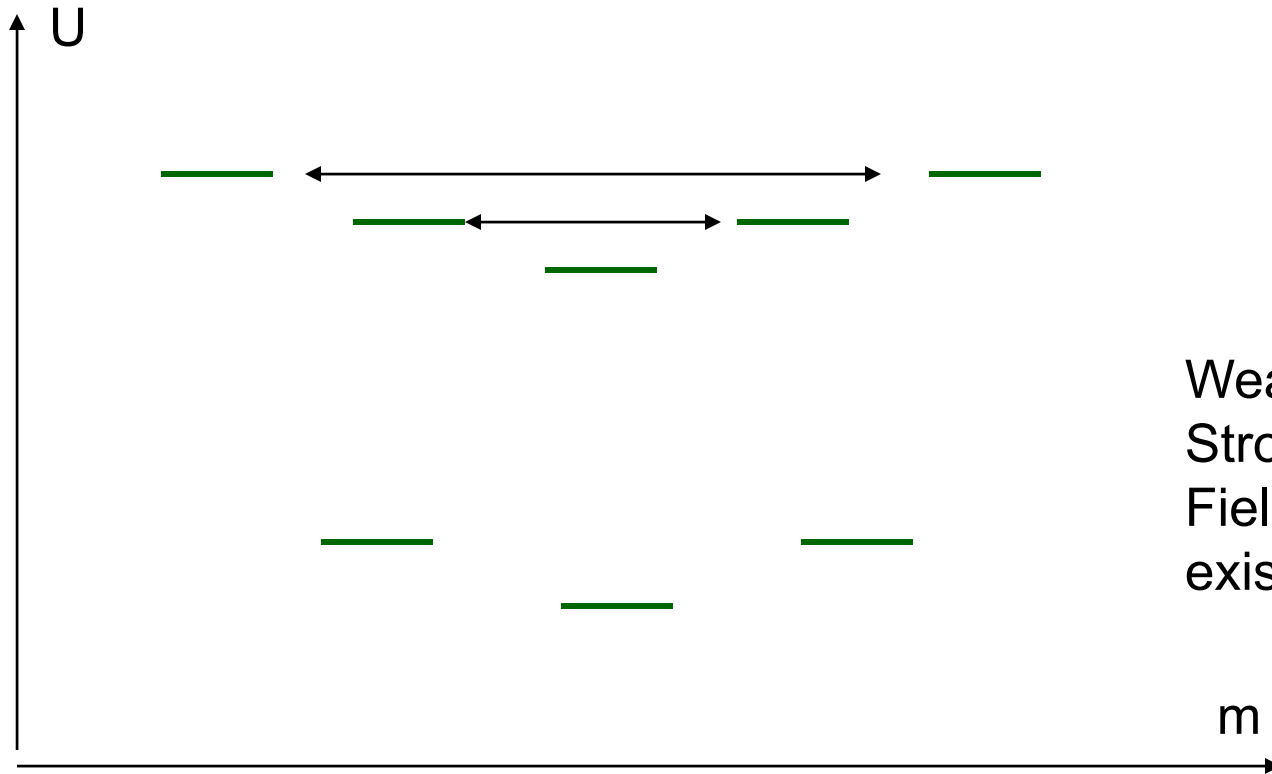


Weak, for atoms.
Strongish, for molecules.

Traps for precision measurement

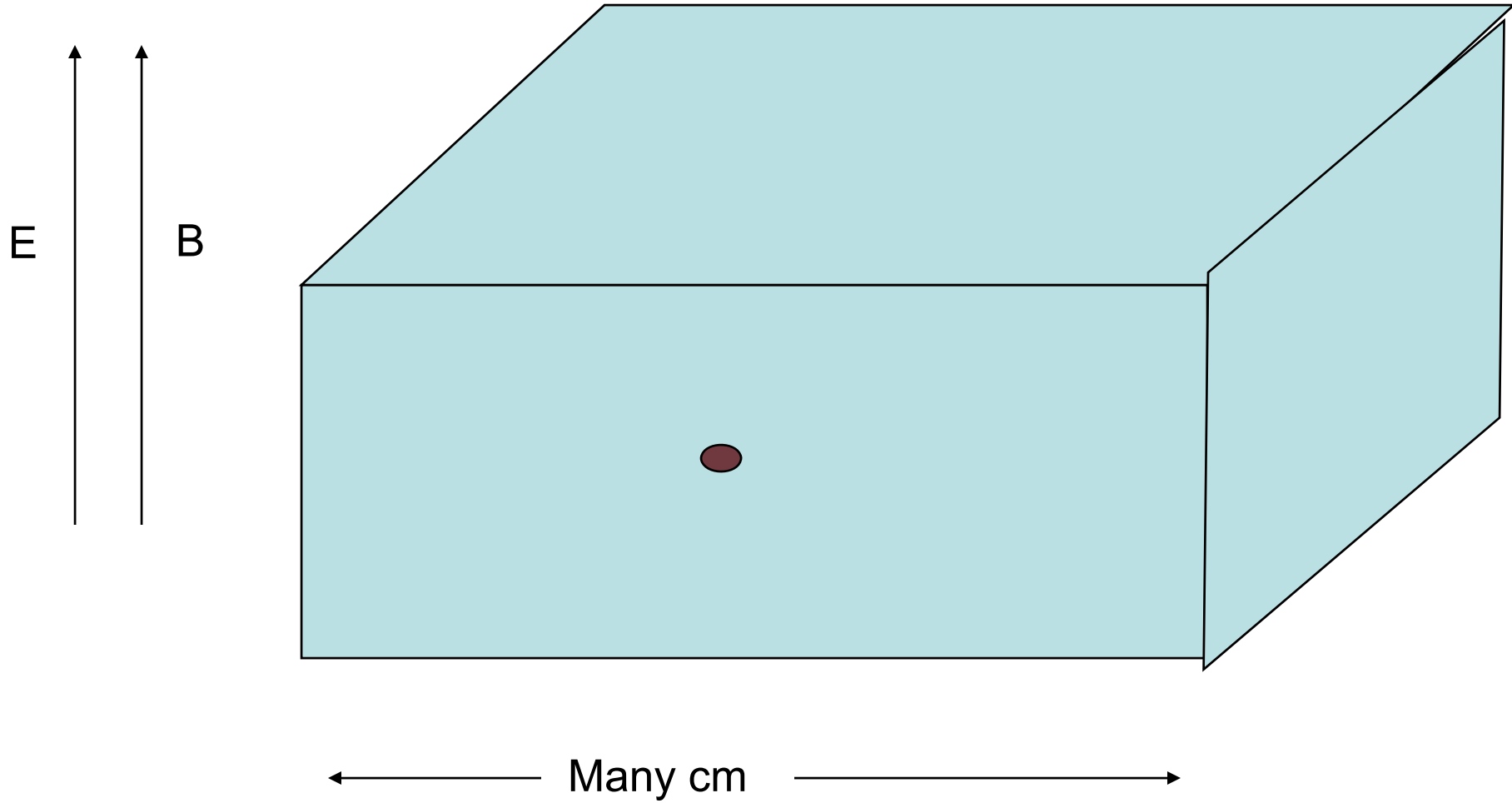
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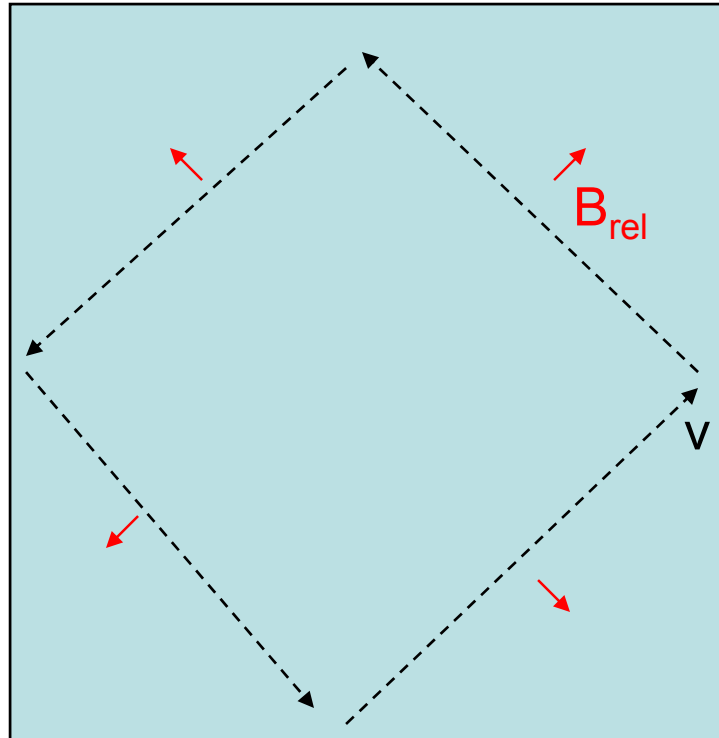
Weak, for atoms.
Strongish, for molecules.
Field insensitive transitions exist.

Neutron-in-a-box (literally)



B_0 , E_0 , point up out of the screen

Neutron motion partially transforms strong electric field into B-field.

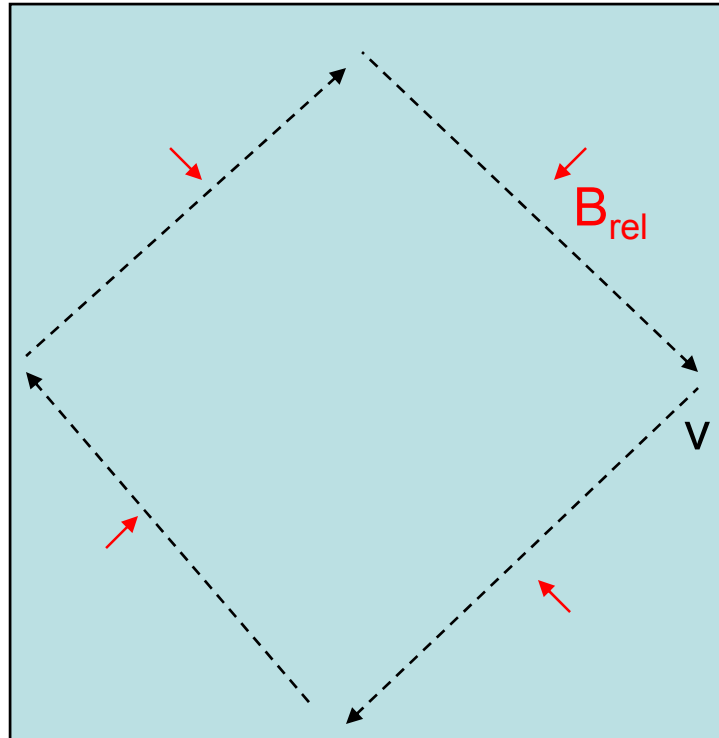


Go back to this case:
No dirt (no spatial
gradient in B) means
no systematic. But,
what about dephasing?

Enclosed area of neutron trajectory means
enclosed area of B -vector in time. A shift in phase between
 $m=1/2$ and $m=-1/2$ levels!

B_0, E_0 , point up out of the screen

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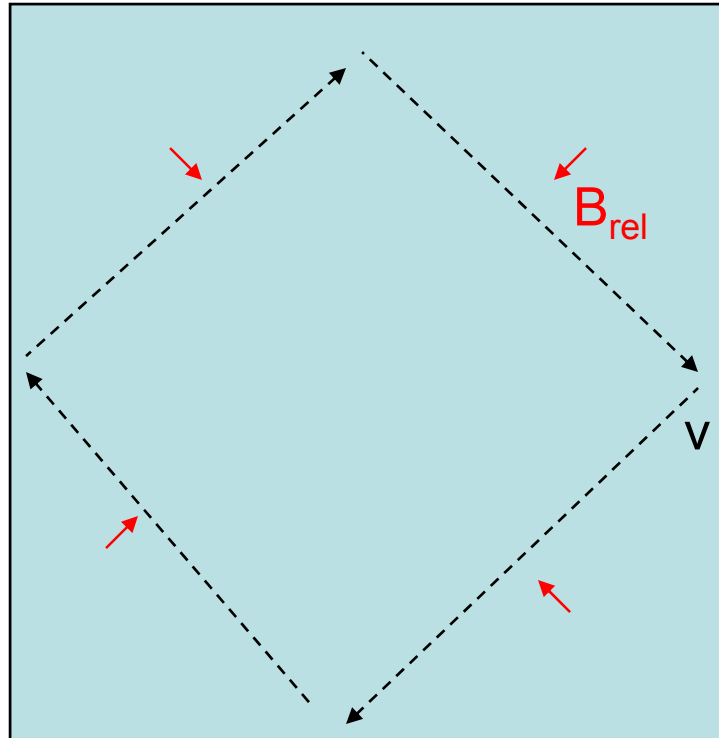
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Thermal distribution of trajectories means this effect as no net sign.

Enclosed area of neutron trajectory means enclosed area of B-vector in time. A shift in phase between $m=1/2$ and $m=-1/2$ levels!

OK for a box. What about trapped particles!?

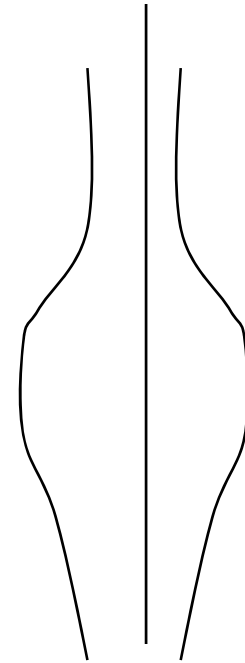
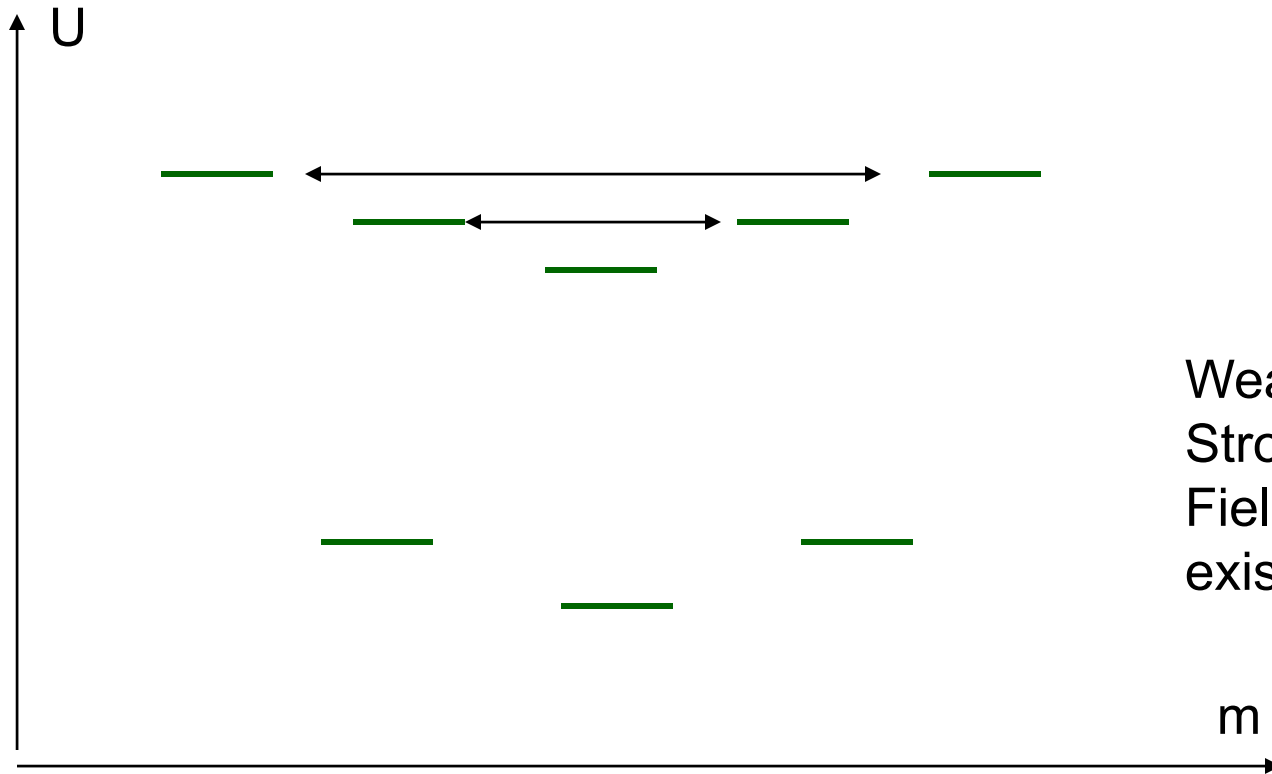
Dephasing due Berry's phase arising from random motion in inhomogeneous trapping fields.

Potentially big problem, except if atoms/molecules are "cold enough".

Traps for precision measurement

Electrostatic trap (Ioffe-Pritchard) for neutrals

$$\text{Force} = \text{grad} (|E| d(E))$$

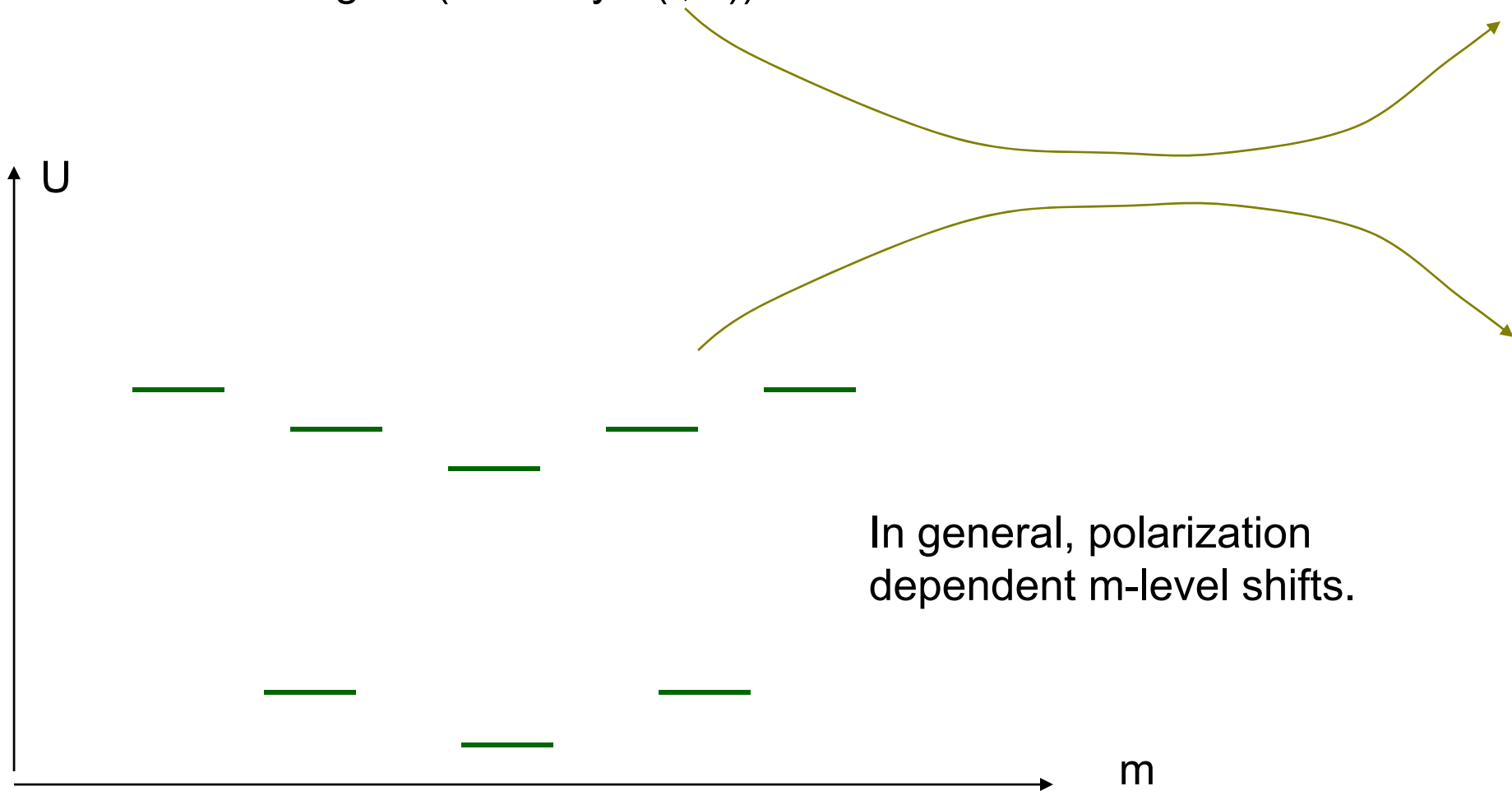


Weak, for atoms.
Strongish, for molecules.
Field insensitive transitions exist.

Traps for precision measurement

Optical dipole traps for neutrals

$$\text{Force} = \text{grad} (\text{Intensity } P(I, \nu))$$



In general, polarization dependent m -level shifts.

Traps for precision measurement

Paul trap for ions

Force = $\mathbf{E} e$ (big) Field can be much more spatially homogenous.

At trap center $\langle \mathbf{E} \rangle = 0$

But you can have a rotating bias field.

Symmetry arguments constrain systematics.

Traps for precision measurement

Others: Penning trap. TOP trap. Ioffe-Pritchard trap variants.

Eric's Tips for Better Precision Measurement with Atoms and Molecules (and with other things)

Eric's Tips for Better Precision Measurement with Atoms and Molecules (and other things)

1. In general, frequencies are the easiest thing to measure with precision. From mHz to EHz, you can get clocks stable to 10^{-14} or better, absolute accuracy to 10^{-13} . Take advantage! Try to turn the quantity you want to measure into a frequency.

- a. Voltage: Josephson junction oscillation frequency.
- b. Magnetic field: Zeeman splitting
- c. Electric field: Stark shift.
- d. optical intensity: ac Stark shift.
- e. mass: cyclotron frequency
- f. capacitance: resonance of LC circuit.
- g. distance: resonance of a fabry-perot laser cavity.
- h. force: ???
- i. etc

Eric's Tips for Better Precision Measurement with Atoms and Molecules (and other things)

2. D.C is where precision measurements go to die. Get as far away as you can!

Example. Lens-Thiring effect

Eric's Tips for Better Precision Measurement with Atoms and Molecules (and other things)

3. If you want to measure a very small oscillating field, use heterodyne detection (as e.g. alternative to “photon counters.”)

Eric's Tips for Better Precision Measurement with Atoms and Molecules (and other things)

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Corollary: for quantum mechanical effects, add an “offset amplitude.”

Eric's Tips for Better Precision Measurement with Atoms and Molecules (and other things)

4. Experimenters (and numerical experimenters): if you want to understand if some imperfection in your experiment is causing you problems, don't make it better: make it worse!

Eric's Tips for Better Precision Measurement with Atoms and Molecules (and other things)

5. You are going to die someday.

Eric's Tips for Better Precision Measurement with Atoms and Molecules (and other things)

5. You are going to die someday. Stop smoking.

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Eric's Tips for Better Precision Measurement with Atoms and Molecules (and other things)

5. You are going to die someday. Stop smoking. Wear a seatbelt. Most important: Have a sensible data collection strategy. It's like wearing a seatbelt in the lab!

Eric's Tips for Better Precision Measurement with Atoms and Molecules (and with other things)

6. What field provides your quantization axis?

Eric's Tips for Better Precision Measurement with Atoms and Molecules (and with other things)

7. \hbar is a small number. But N_A is a big number!

Is $N_A \hbar < 1$?

Is $N_A^2 \hbar < 1$?

"The Casimir-Polder Force"

The force experienced by an atom near a surface, arising from spatial patterns in the fluctuations in the E&M field.
With implications for anomalous gravity stuff.

John Obrecht

Rob Wild [Dave Harber]

Thanks, Colleen Gillepsie, Giacomo Roati

NSF, NIST

Other experiments:

Hinds

Westbrook/Aspect

Vuletic, Shimizu, Ketterle...

Theory: London, Casimir,

Polder, Lifshitz

More recently, Eberlein,

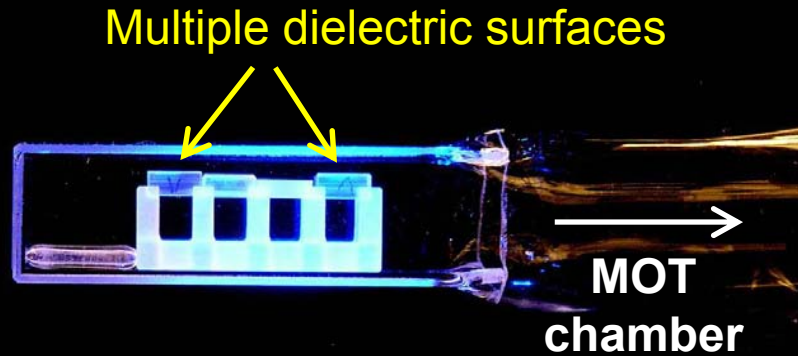
Henkel

We acknowledge

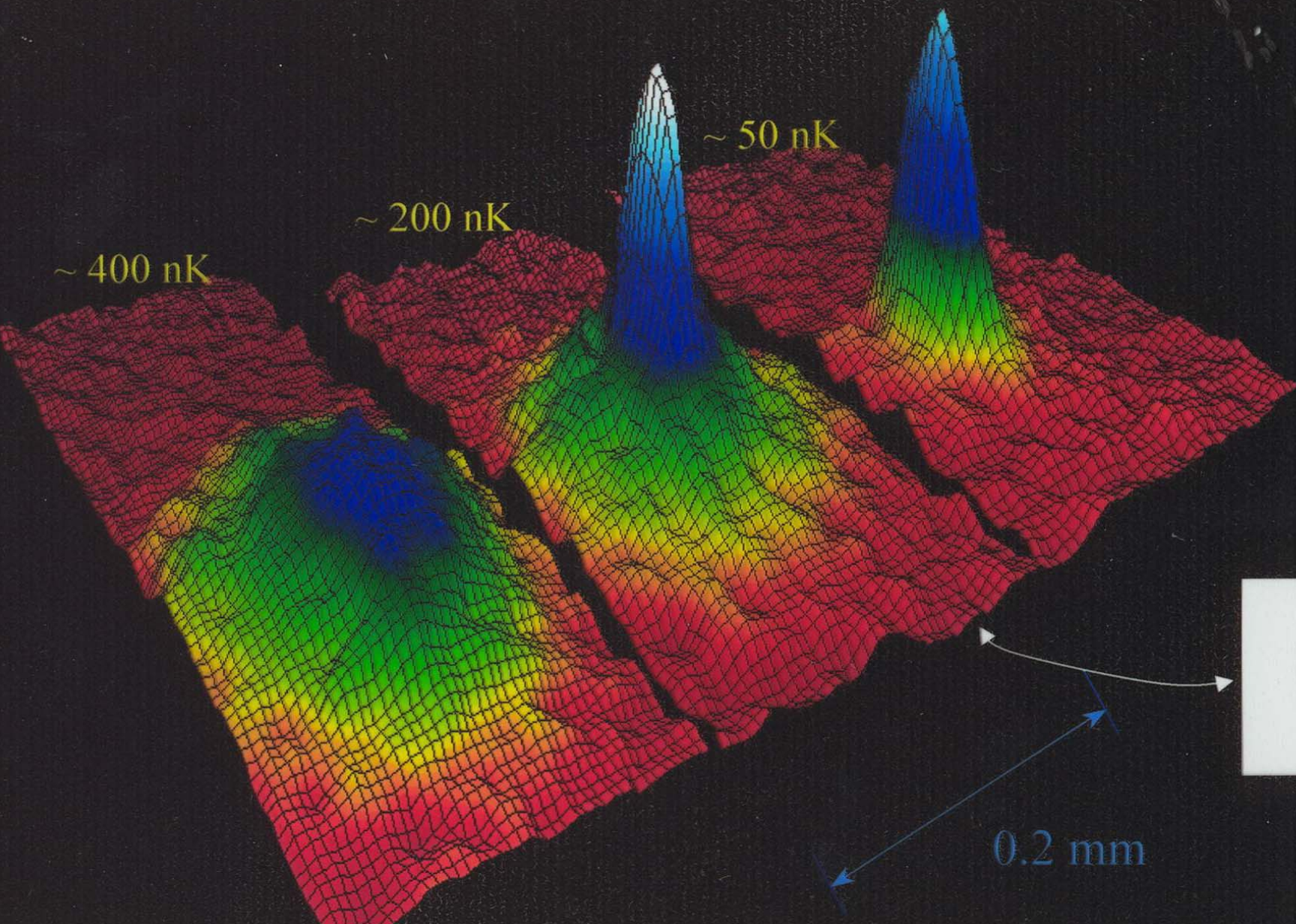
"I Tre Trentini":

Mauro Antezza, Lev Pitaevskii
and Sandro Stringari

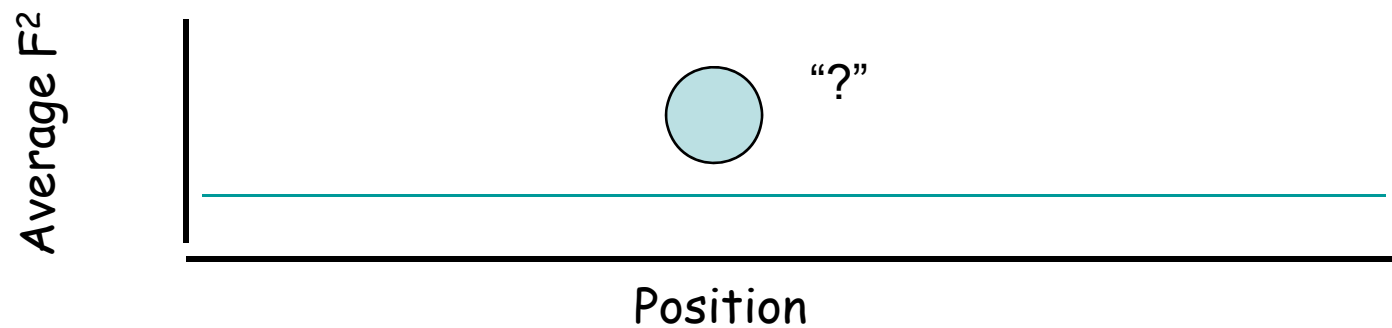
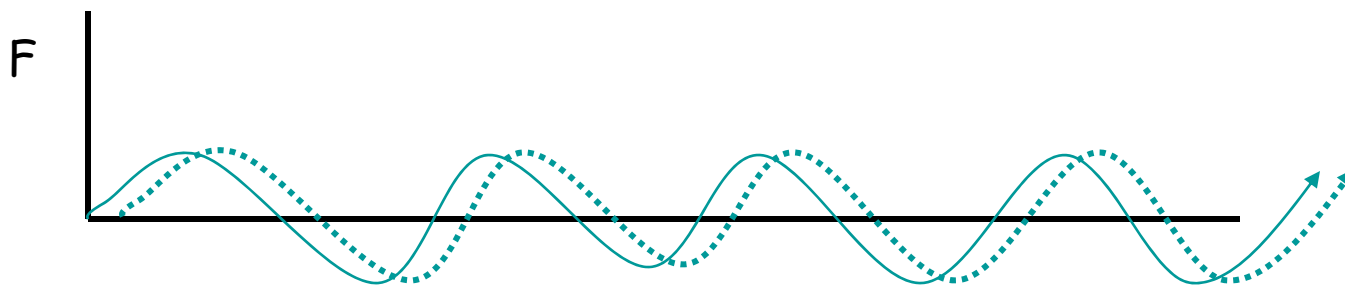
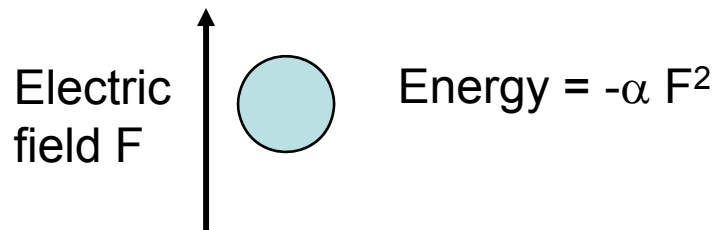
Casimir-Polder force near a dielectric surface



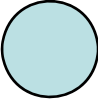
2 D velocity/density distributions

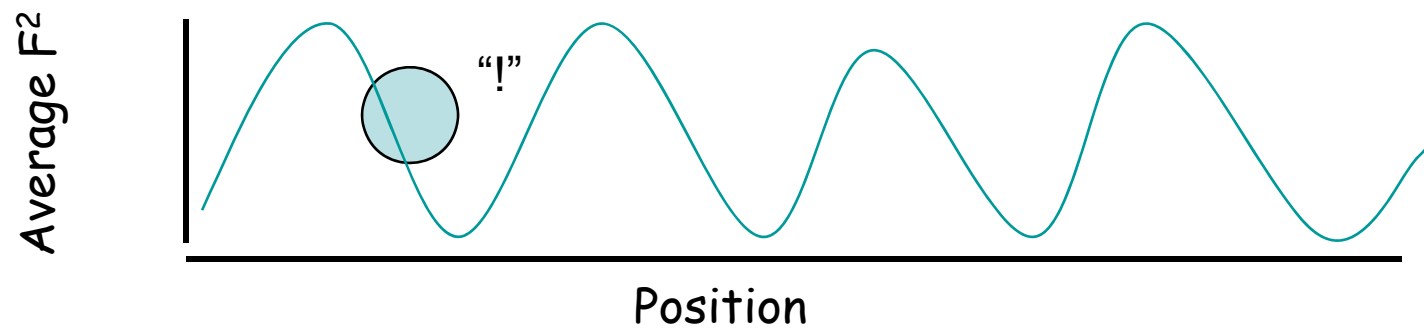
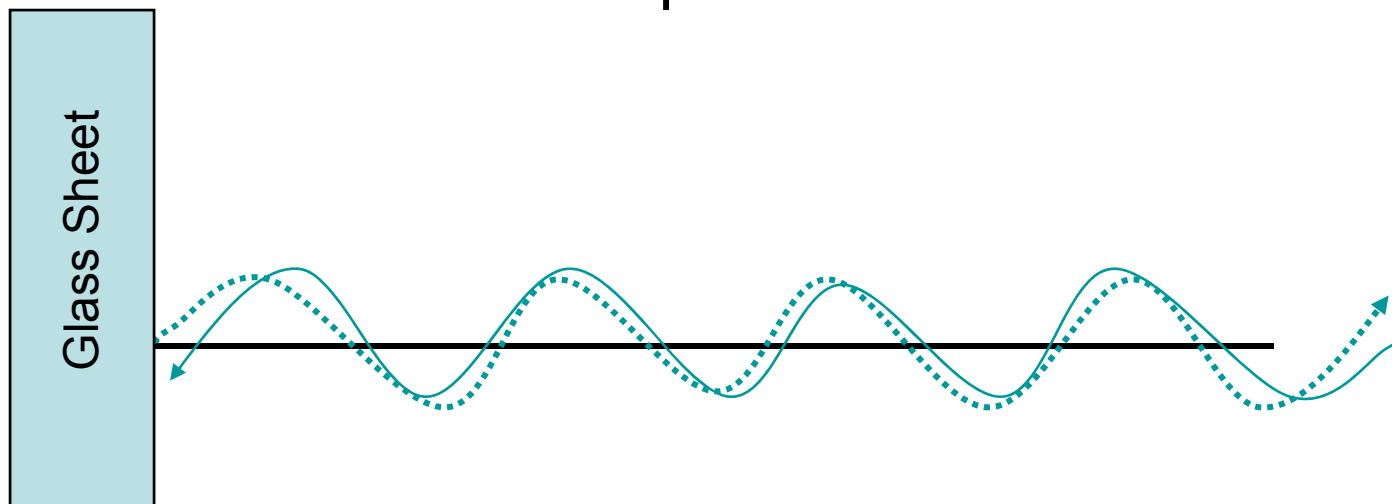


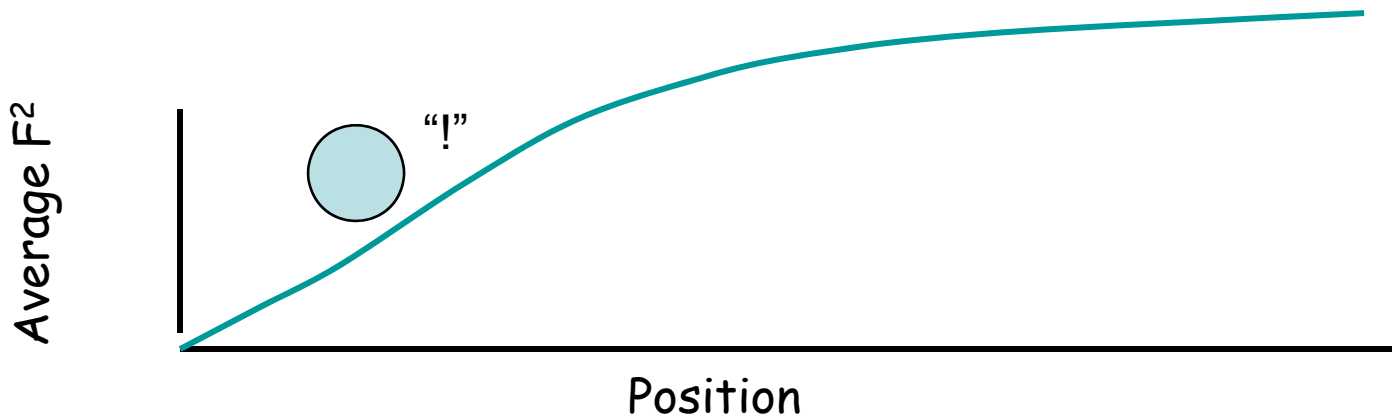
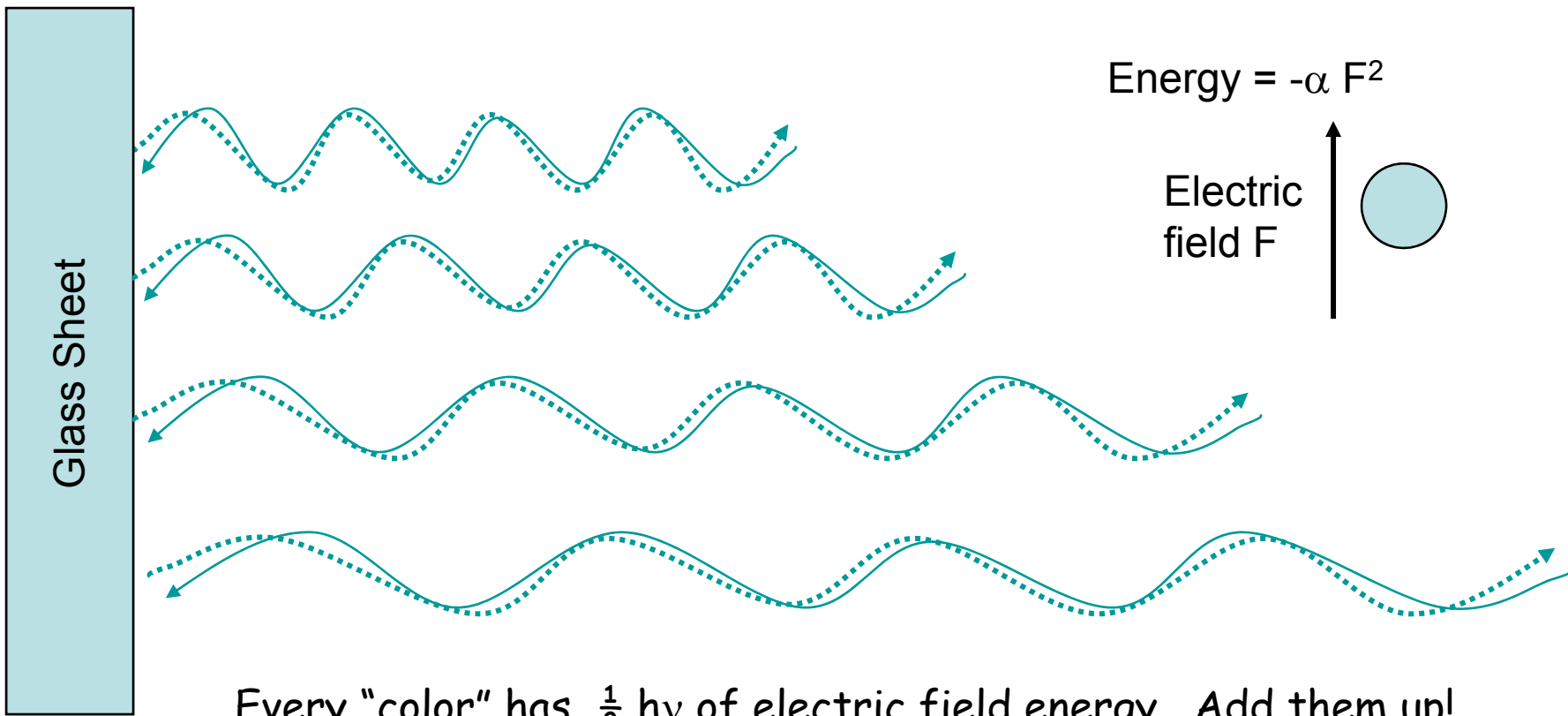
Use an atom.

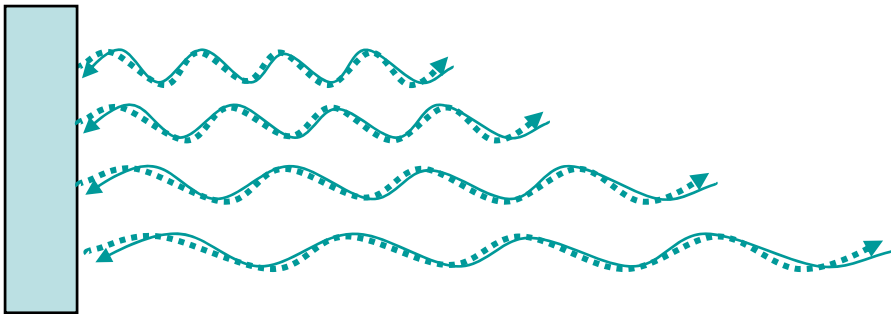


Use an atom.

Electric field F  Energy = $-\alpha F^2$

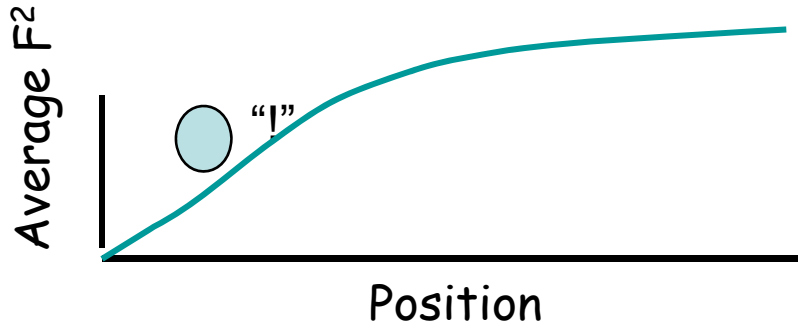
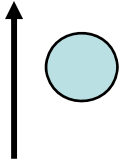






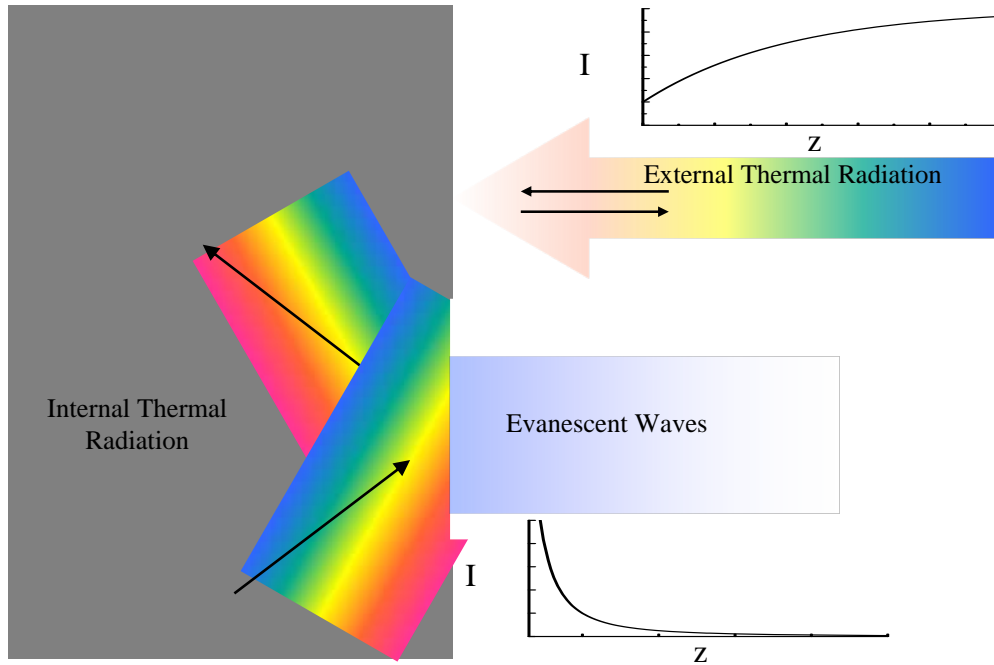
$$\text{Energy} = -\alpha F^2$$

Electric
field F



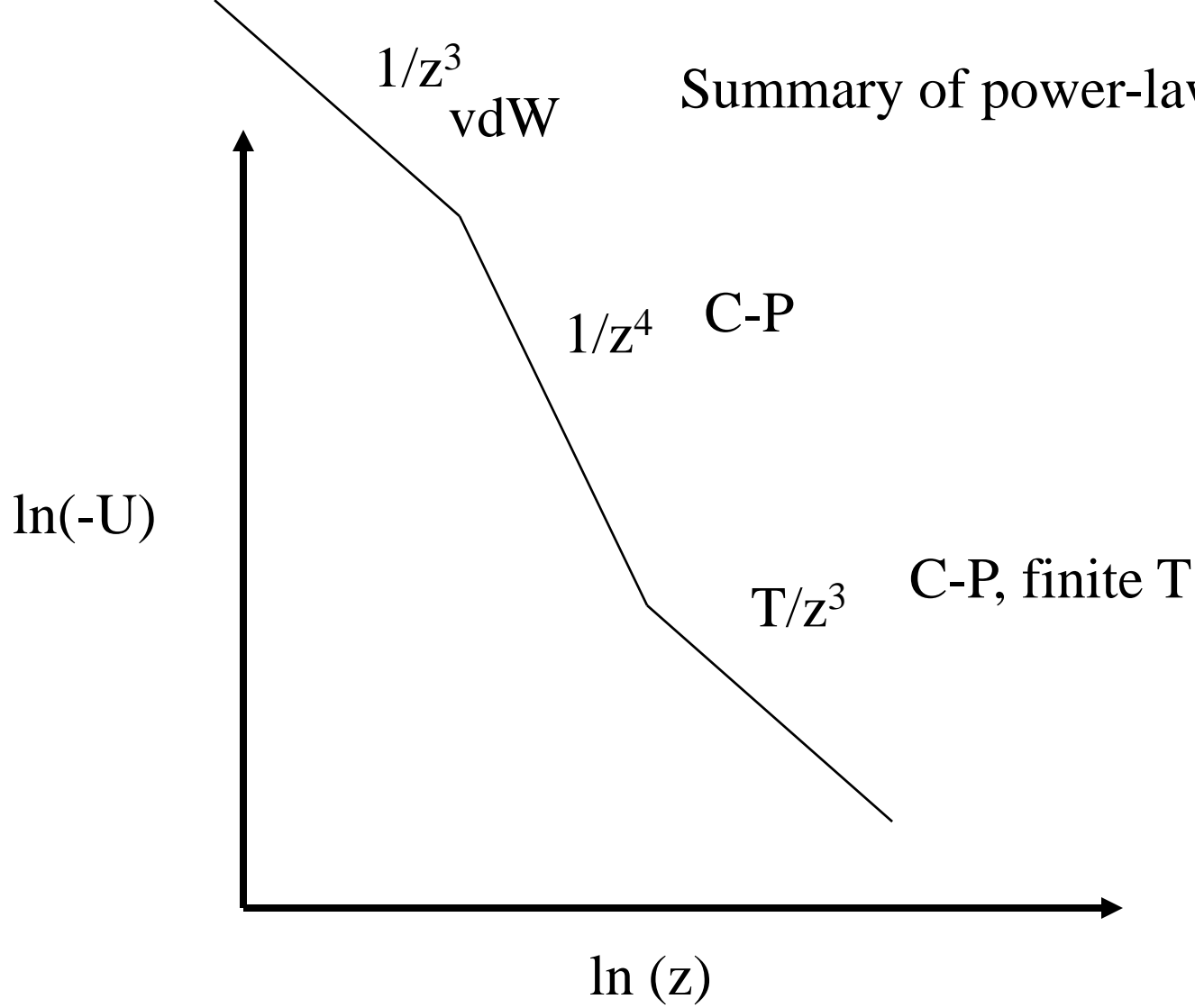
Uh-oh. Experiments say
Casimir-Polder force *pulls atom
in*, doesn't push atom out!

We have neglected evanescent waves,
and incident waves at grazing incidence!

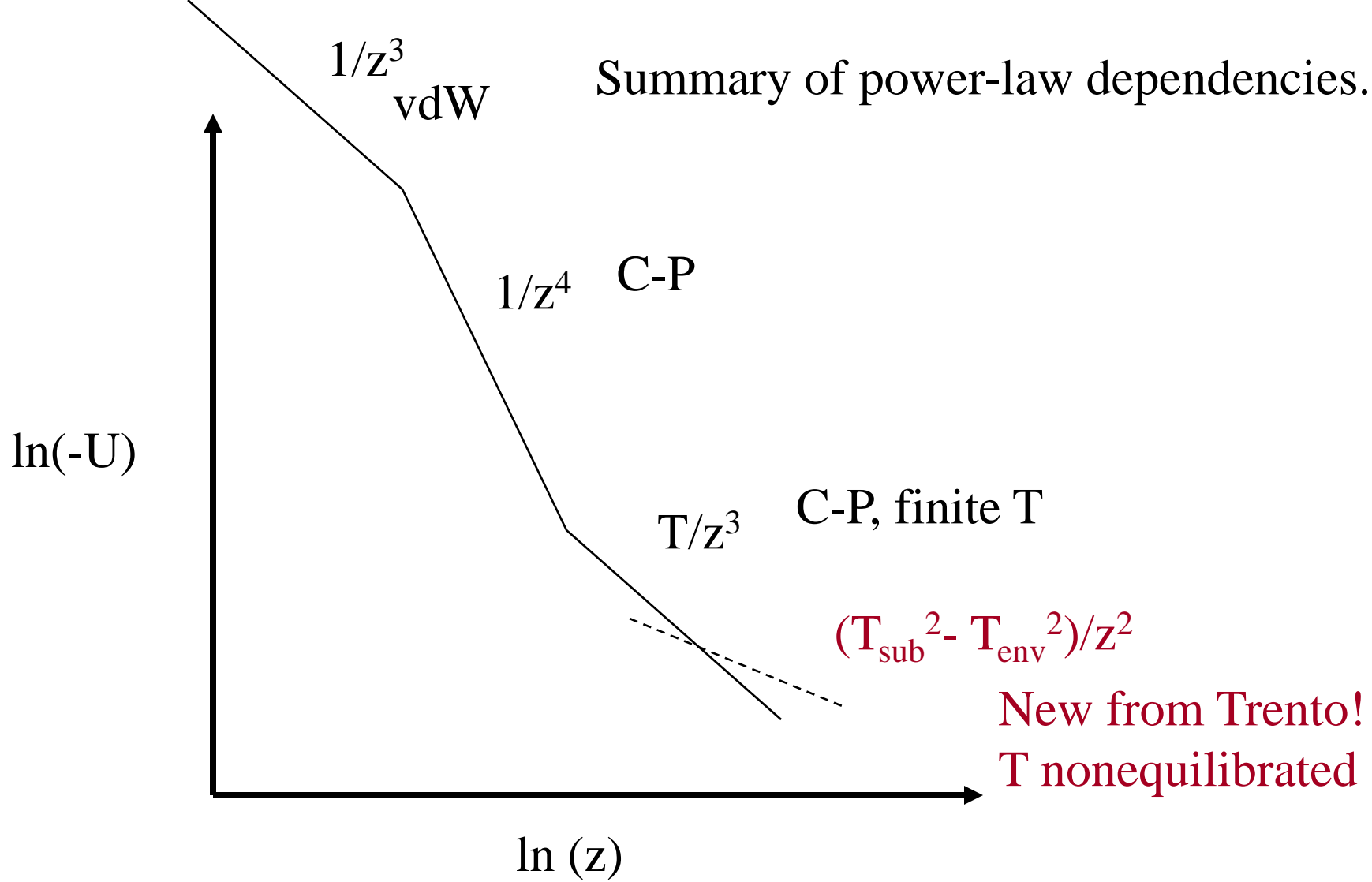


Impinging thermal
(and quantum) radiation
at normal incidence
contributes
with opposite sign to
force compared to
thermal evanescent
waves.

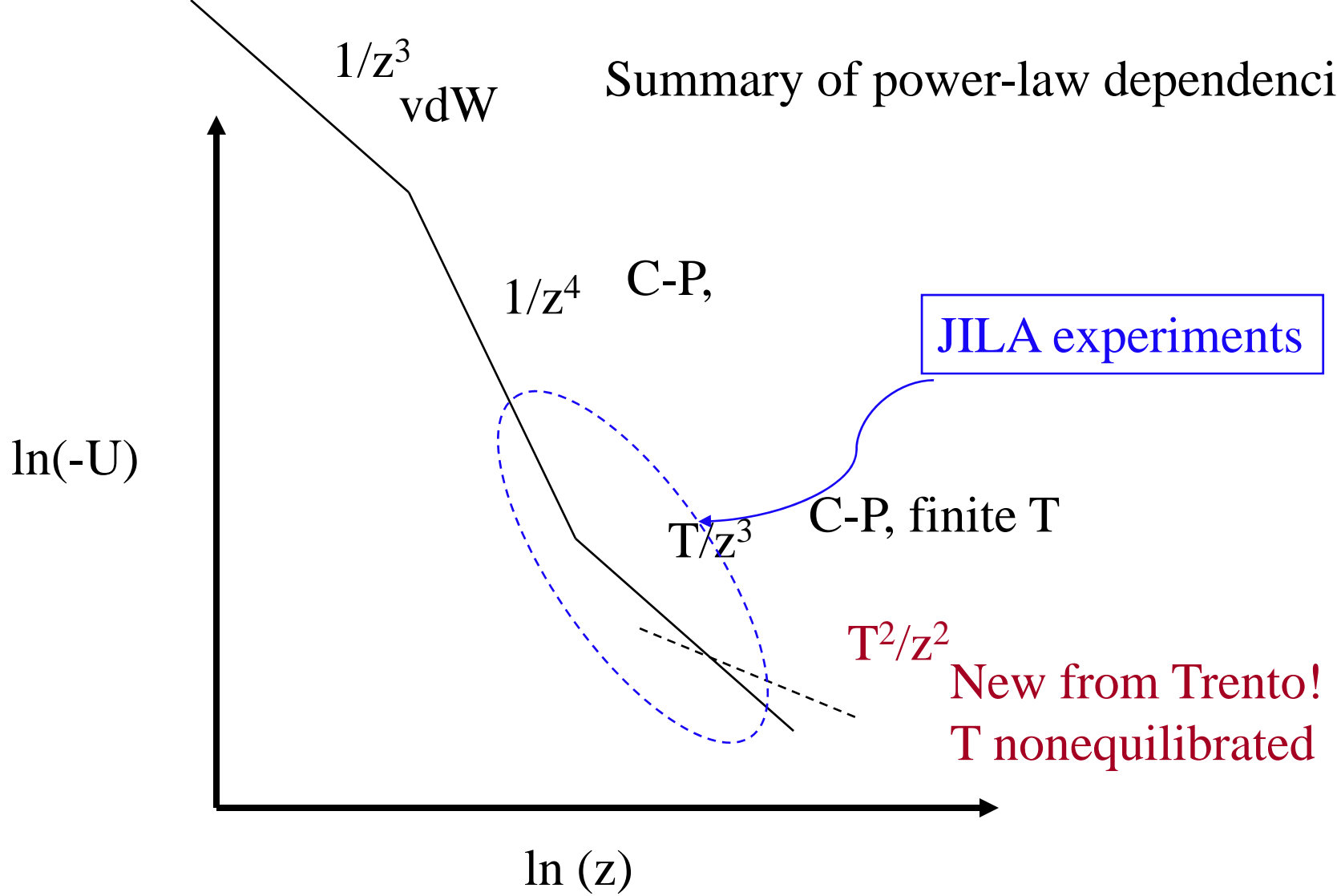
Summary of power-law dependencies.



Summary of power-law dependencies.



Summary of power-law dependencies.



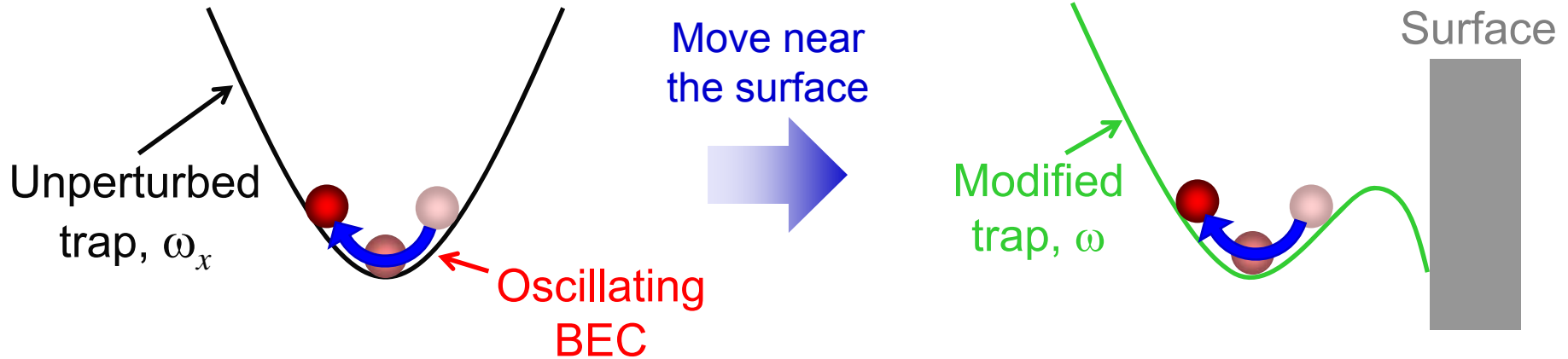
Measuring atom-surface forces

Use trapped BEC as a mechanical oscillator

Measure changes in dipole oscillation frequency

Negative curvature
attractive potential

Trap frequency decrease



Express trap frequency changes as normalized frequency shifts:

$$\frac{\omega_x - \omega}{\omega_x} \approx -\frac{1}{2\omega_x^2 m} \frac{d^2 U}{dx^2}$$

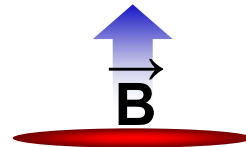


Actual experiment cycle

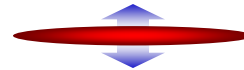
Create BEC



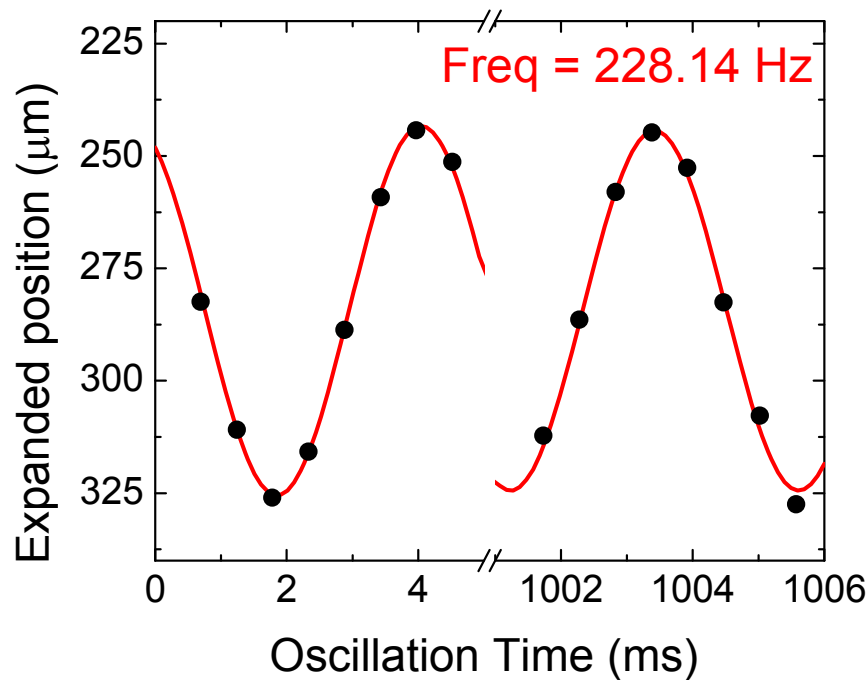
Move near surface



Excite dipole mode



Wait,
image,
repeat

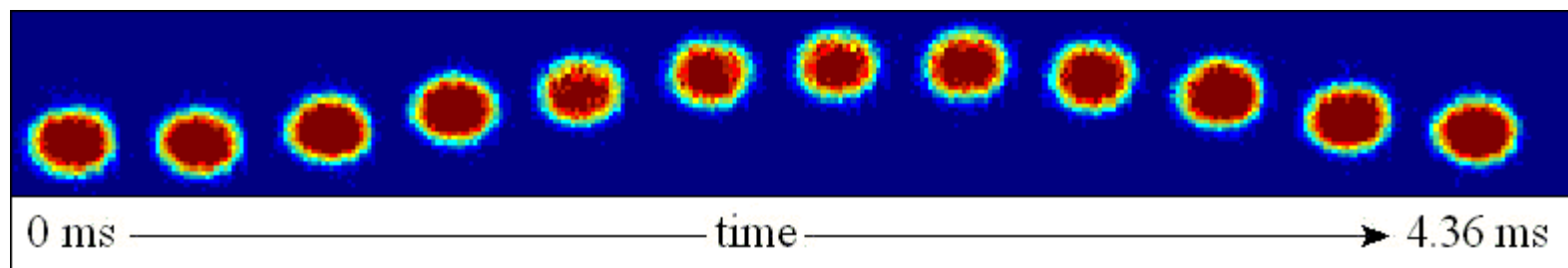
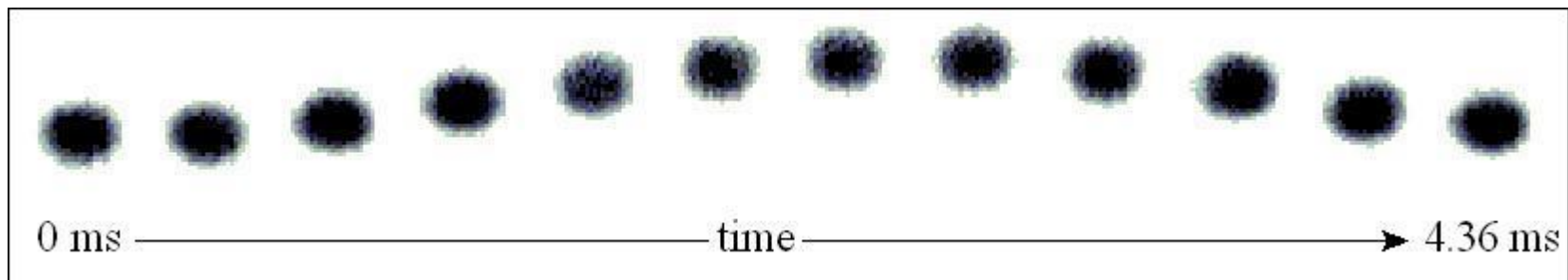


Dipole mode oscillation:

Damping time **~ 10 seconds**

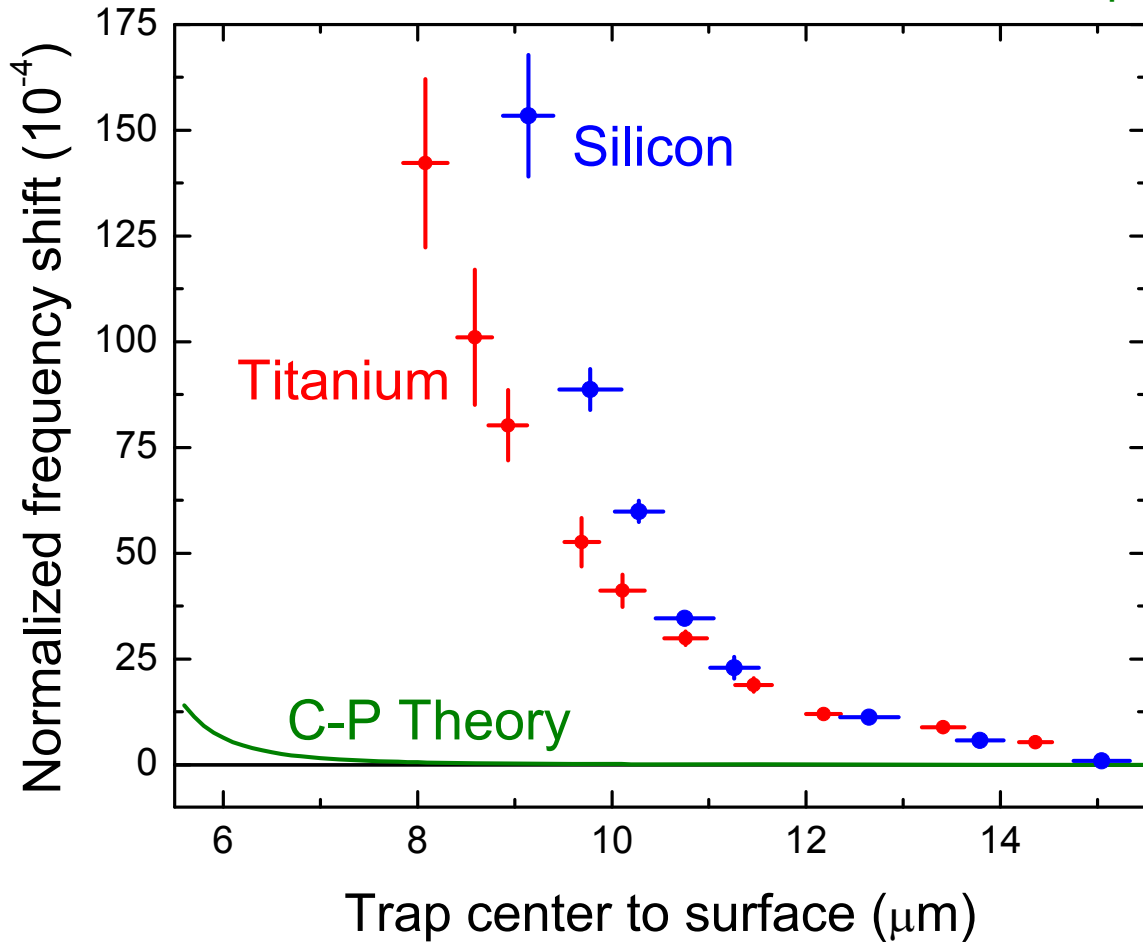
Frequency resolution **~ 10 mHz**

Normalized frequency shift
resolution **$\sim 4 \times 10^{-5}$**





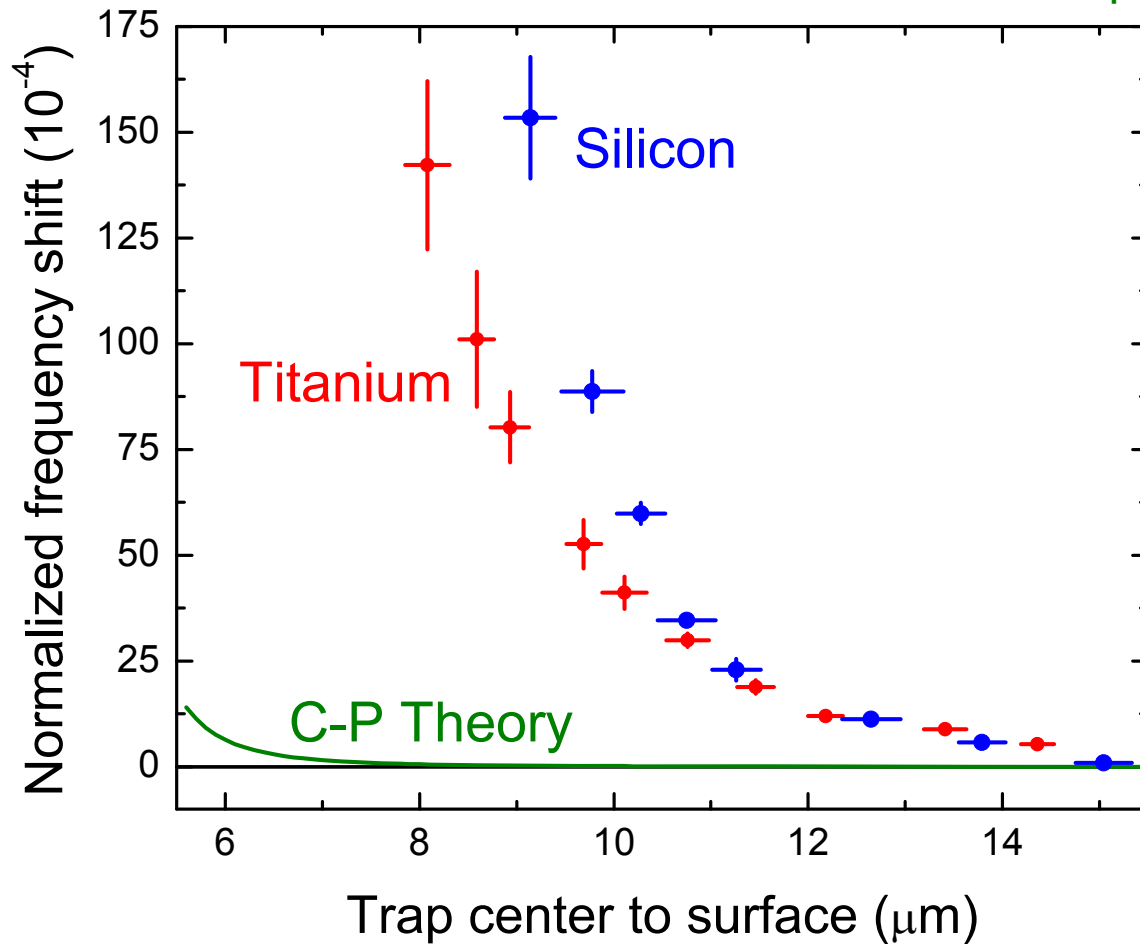
Our first Casimir-Polder measurement attempt:



Electric fields from alkali adsorbates

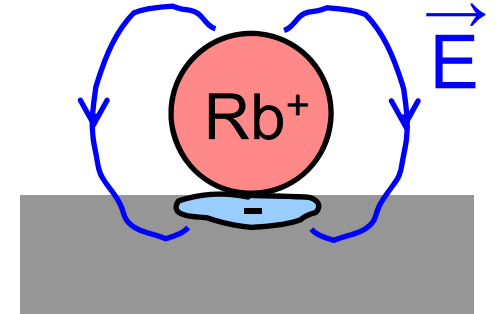
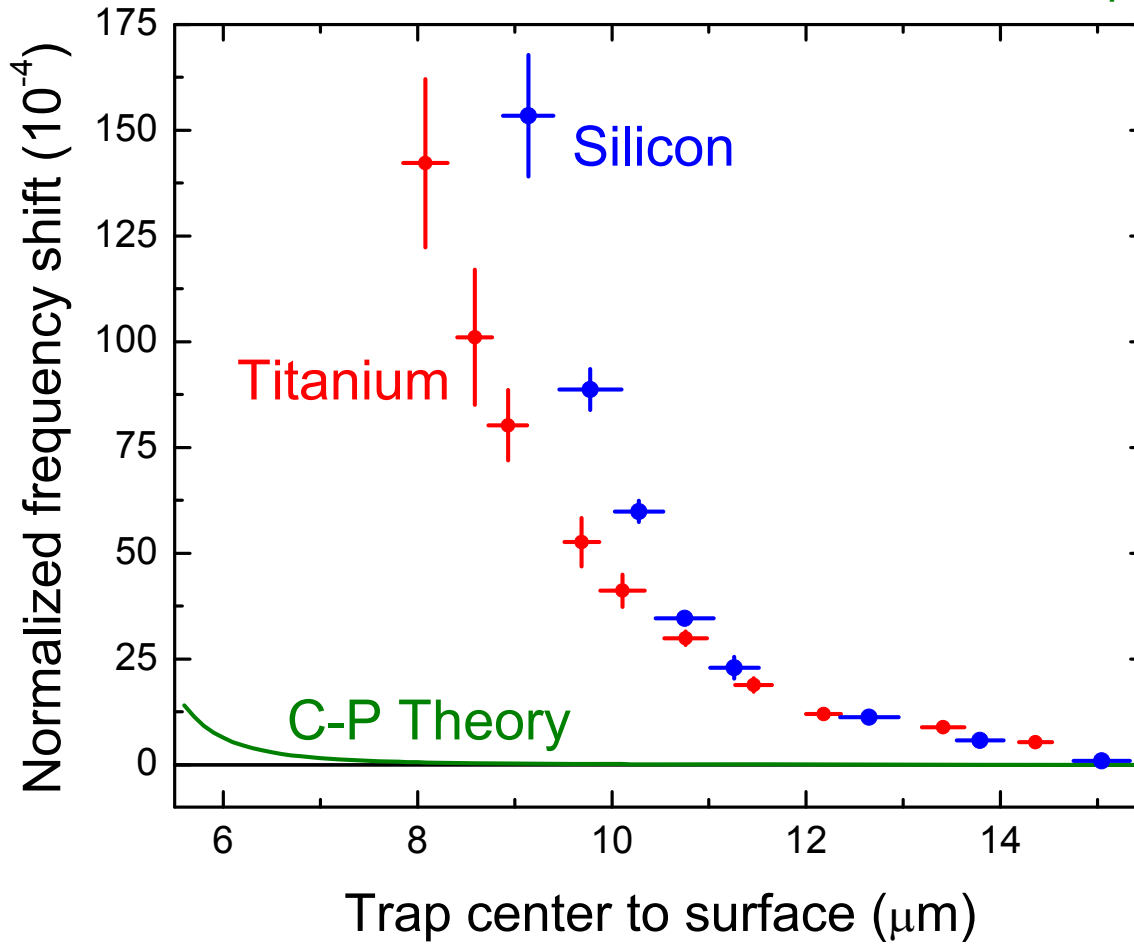


Our first Casimir-Polder measurement attempt:





Our first Casimir-Polder measurement attempt:



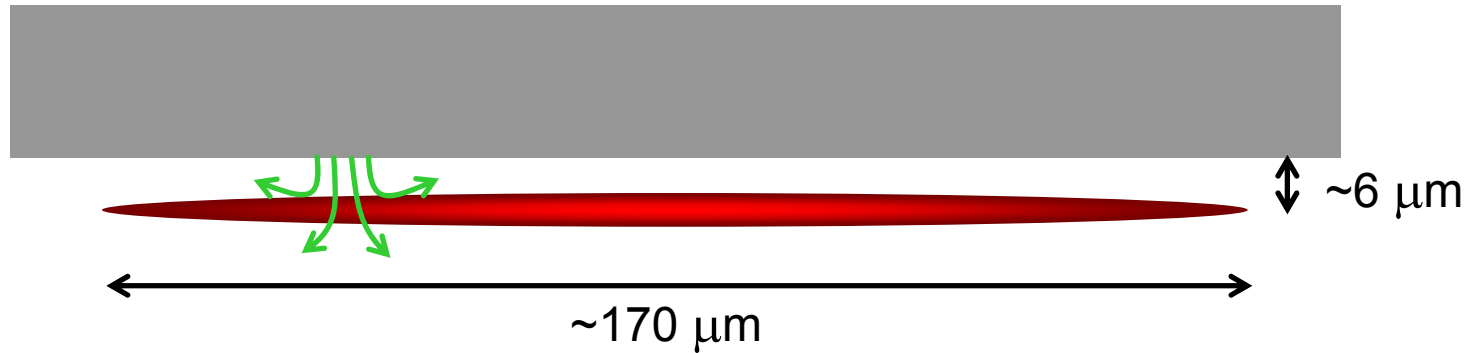
Electric fields from alkali adsorbates



Systematics

How can we put limits on forces from electric and magnetic surface contaminants?

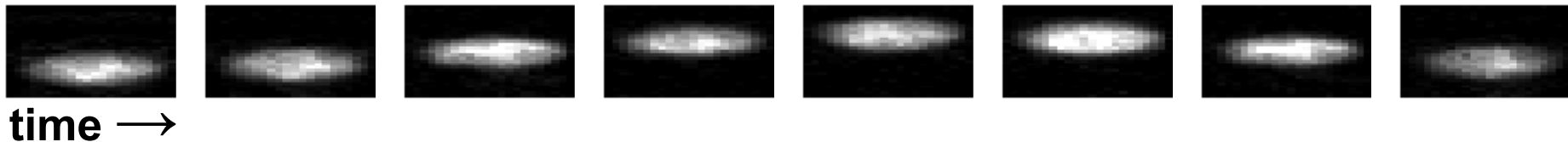
Electric or magnetic surface contaminants are typically localized
→ affect only part of BEC



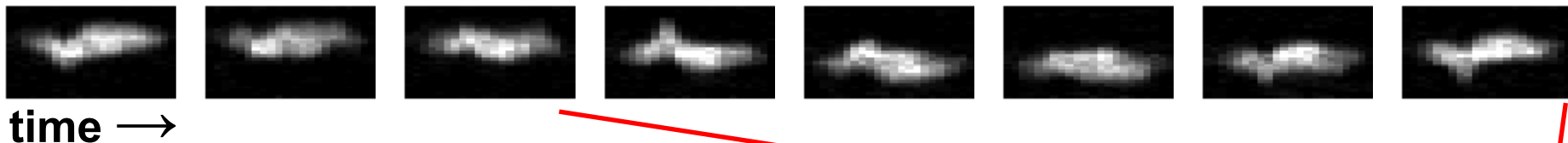
Detect spatially inhomogeneous forces by measuring the normalized frequency shift along BEC

FFS across BEC

Normal BEC oscillation:



BEC oscillation near “contaminated” surface region:

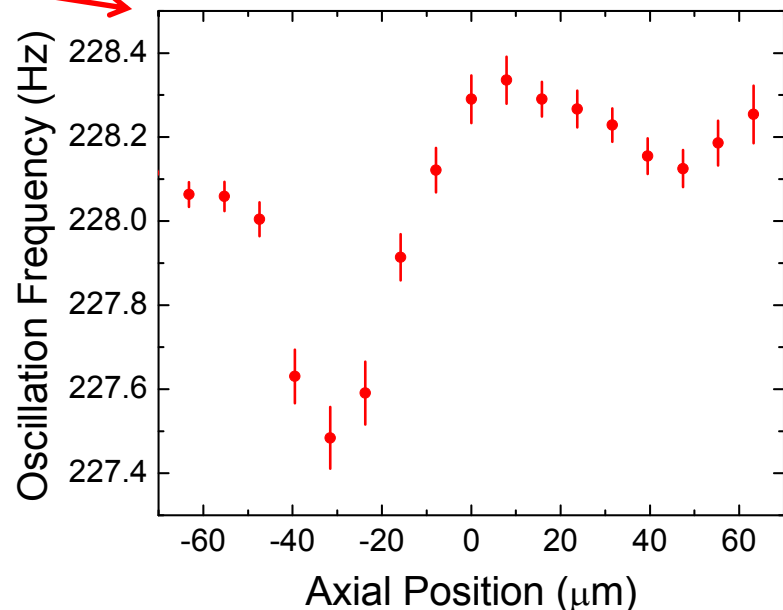


Analyze the oscillation frequency along the BEC:

If spatial variation $>$ statistical uncertainty

→ Significant spatial inhomogeneity

∞ Spatial variation of the oscillation frequency provides limit on spatially inhomogeneous forces



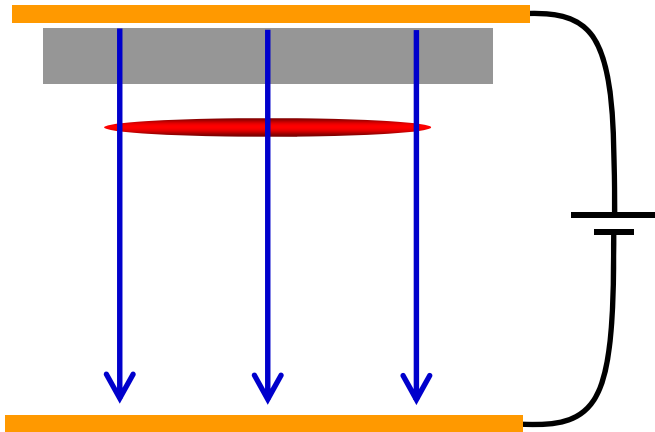


Uniform fields

What about electric & magnetic fields uniform across BEC?

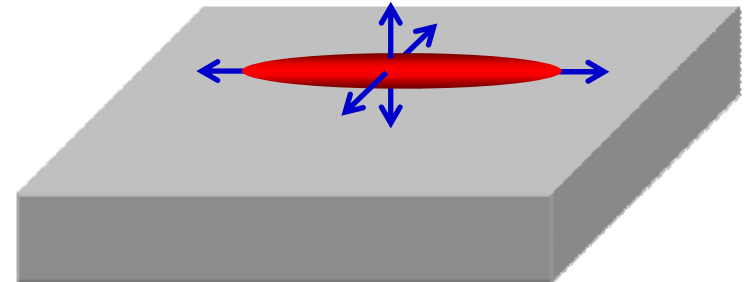
To detect electric fields:

- Use our electric field measurement techniques



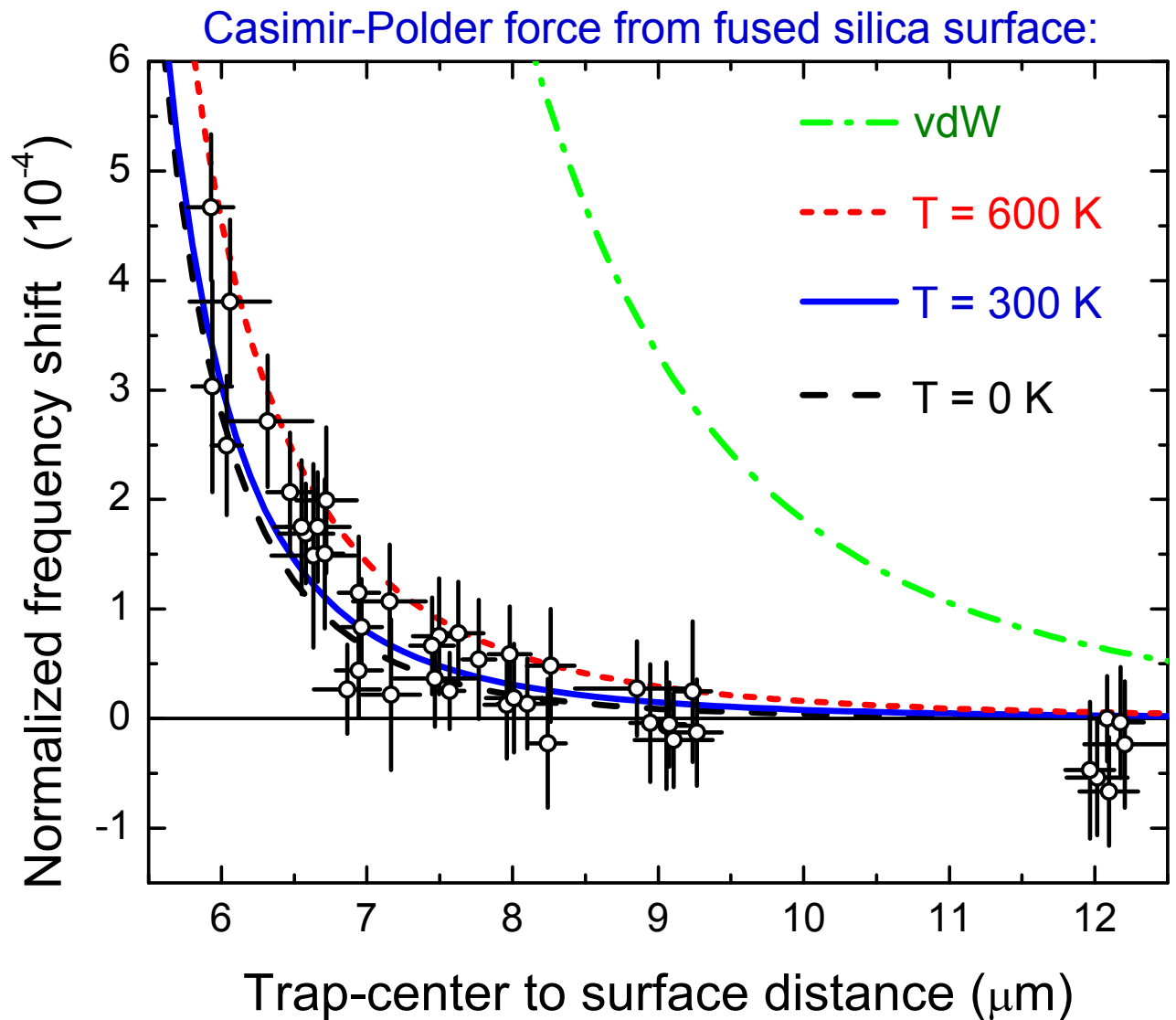
To detect magnetic fields:

- Magnetic distortions modify the trapping potential in multiple directions
- Measure trap frequencies in directions parallel to surface
- Detect center-of-mass position deviations

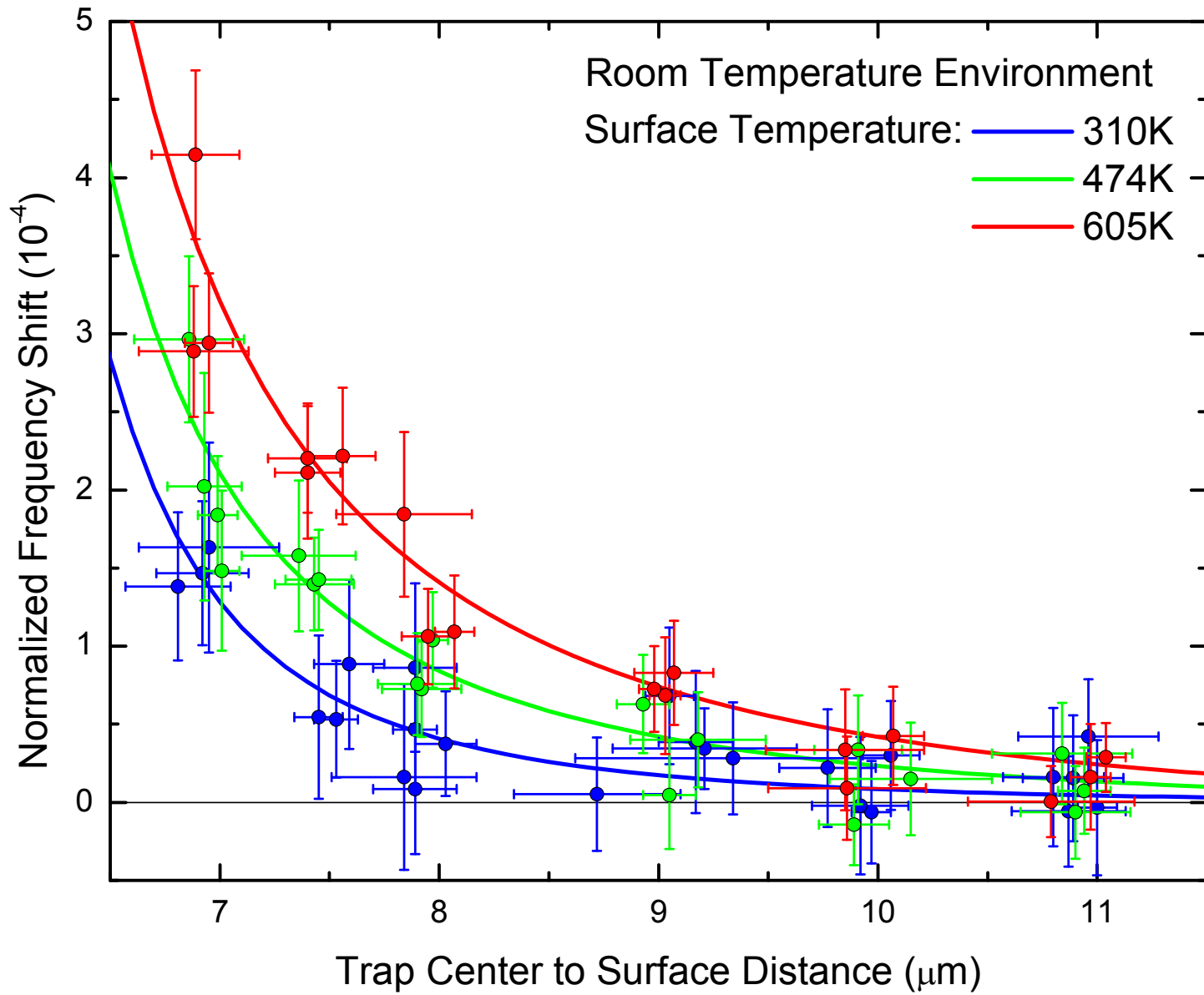


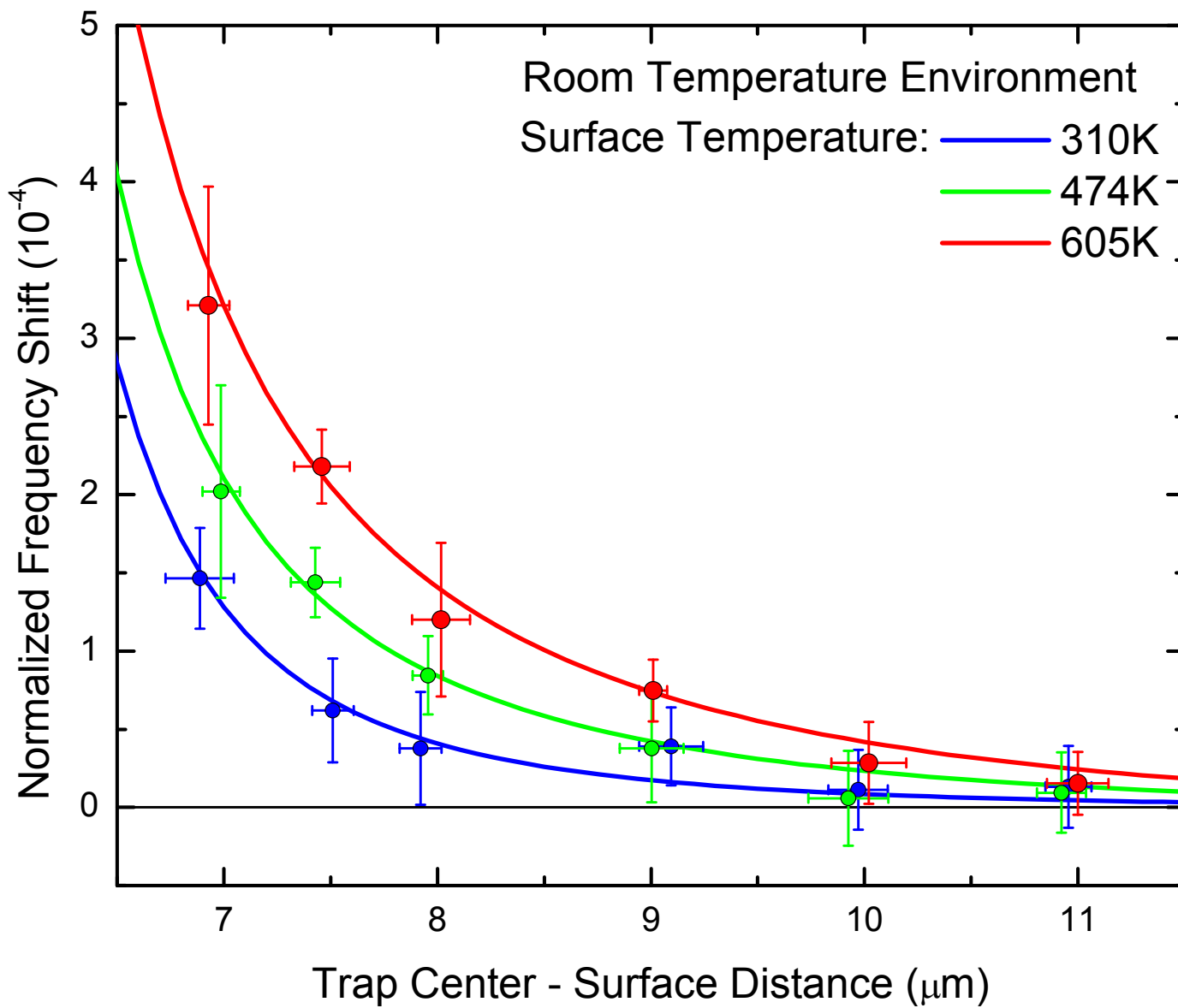


- Data from two different surface locations
- Error bars include statistical and systematic error
- Our data is in agreement with C-P force
- Resolution is not sufficient to discern the temperature correction



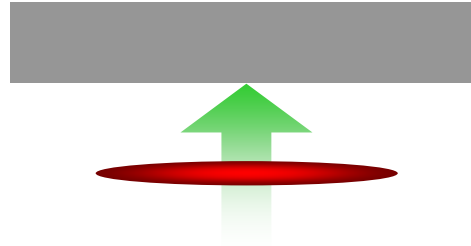
Our measurement





Yukawa-type forces?:

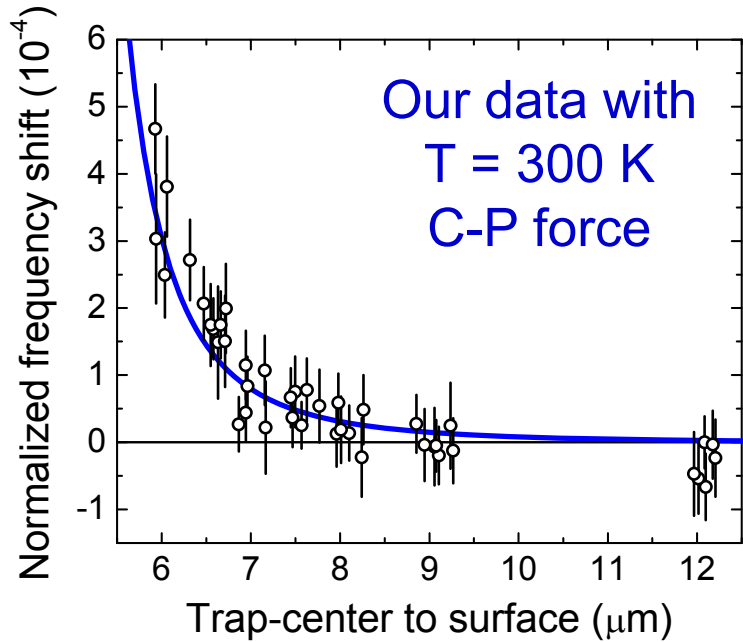
- Exotic force limits from our C-P measurement



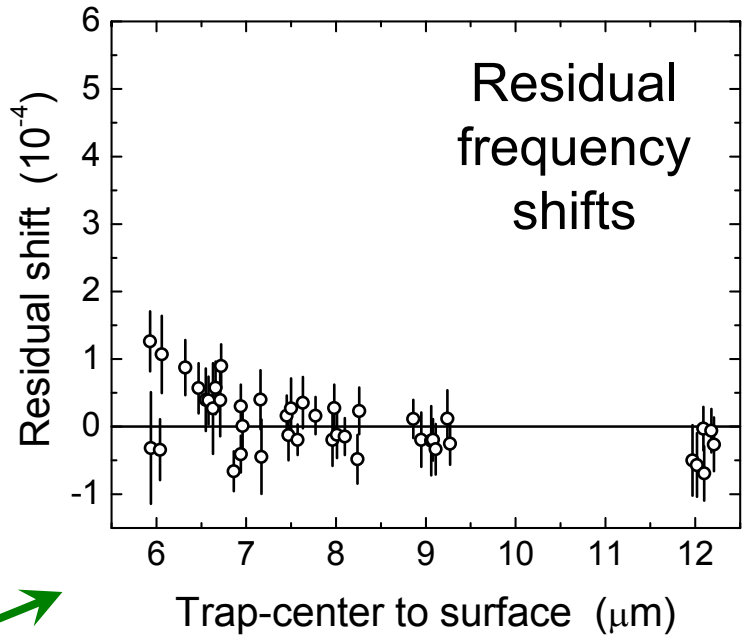
$$U = \int_{\text{Substrate}} \frac{Gm_{\text{Rb}} \rho_{\text{substrate}} dV}{r} \left(1 + \alpha e^{-r/\lambda} \right)$$



The *absence* of forces in addition to C-P force allows us to obtain limits from our data:



Subtract off
C-P force

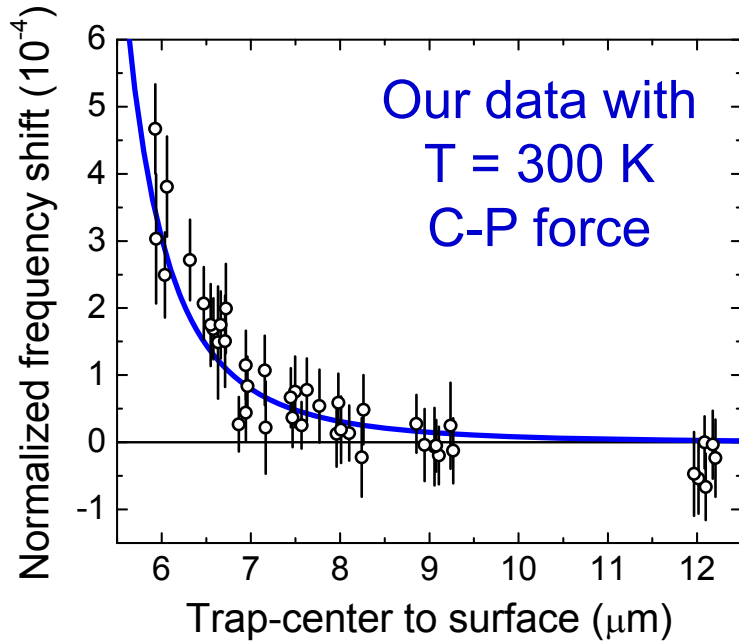


Use residuals to obtain a limit on the presence of additional forces

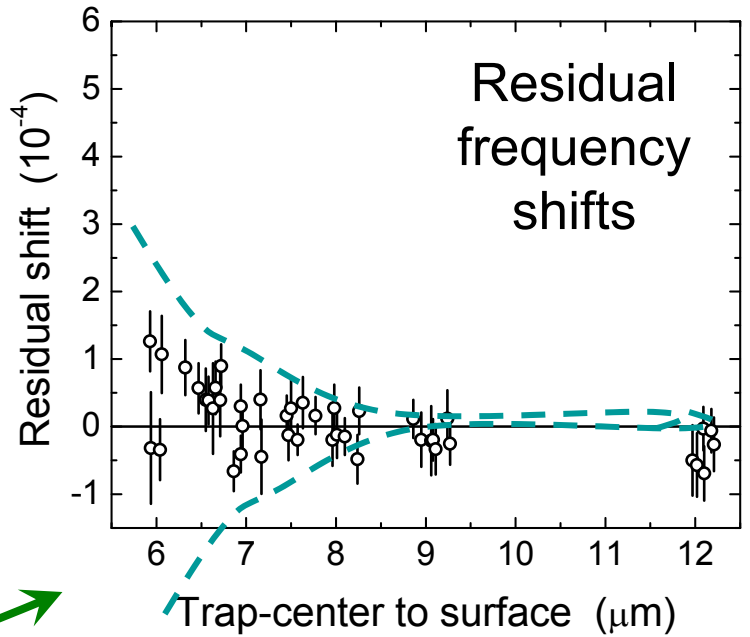
Residuals to the C-P force



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Subtract off
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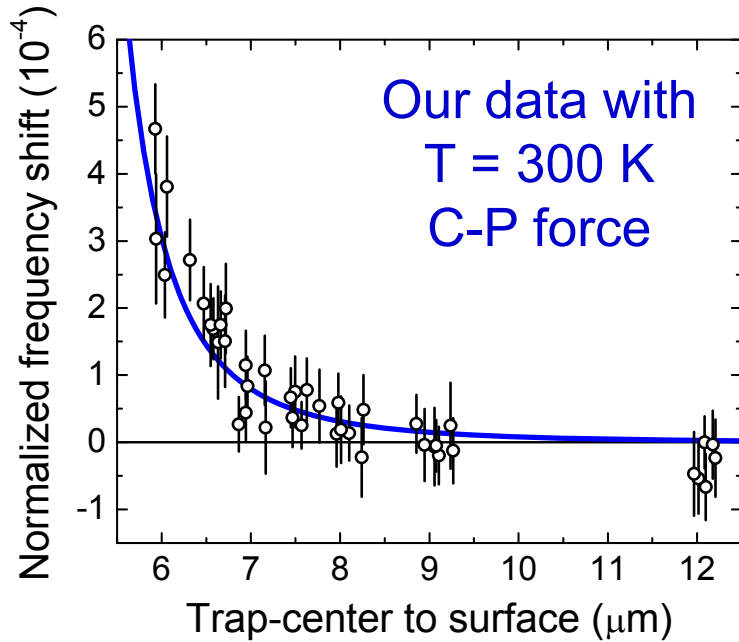
Use residuals to obtain a limit on the presence of additional forces

Trial value of λ
and α

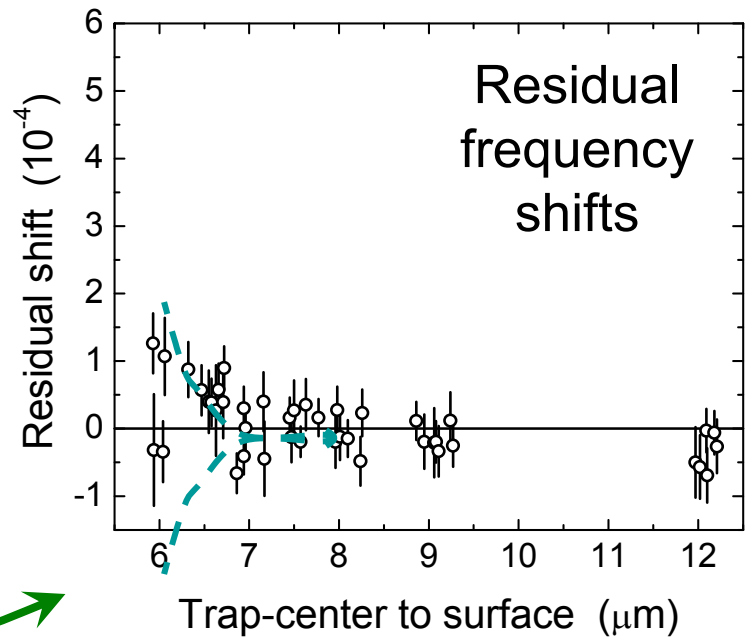
Residuals to the C-P force



The *absence* of forces in addition to C-P force allows us to obtain limits from our data:



Subtract off
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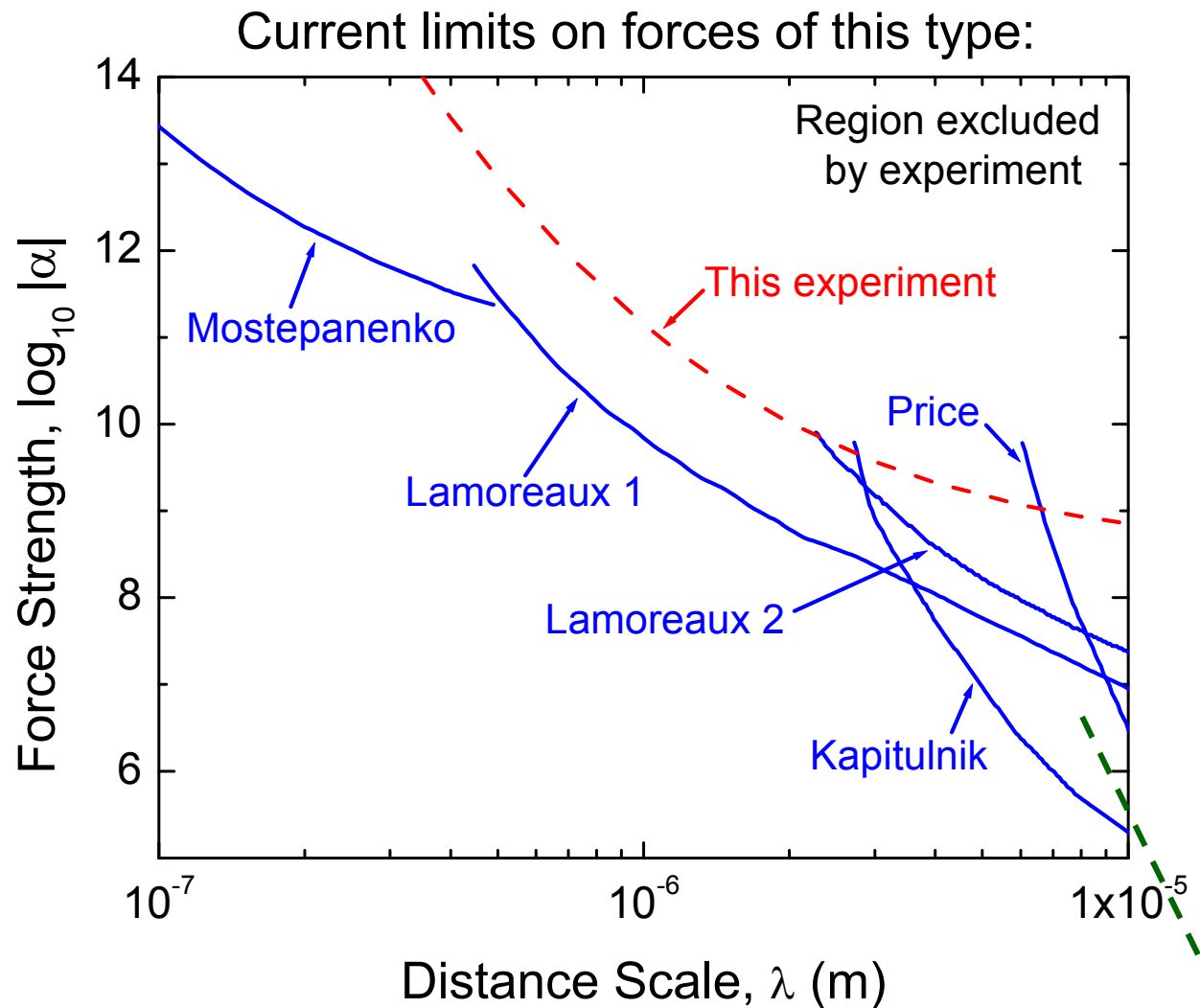
Use residuals to obtain a limit on the presence of additional forces

Smaller trial
value of λ

Residuals to the C-P force

$$U = \int_V \frac{Gm\rho dV}{r} (1 + \alpha e^{-r/\lambda})$$

- Very different type of measurement (atom-bulk vs. bulk-bulk)
- Our experiment does not* reach the current best limits in 0.3-10 μm range
- Experimental modifications could improve sensitivity by over an order of magnitude



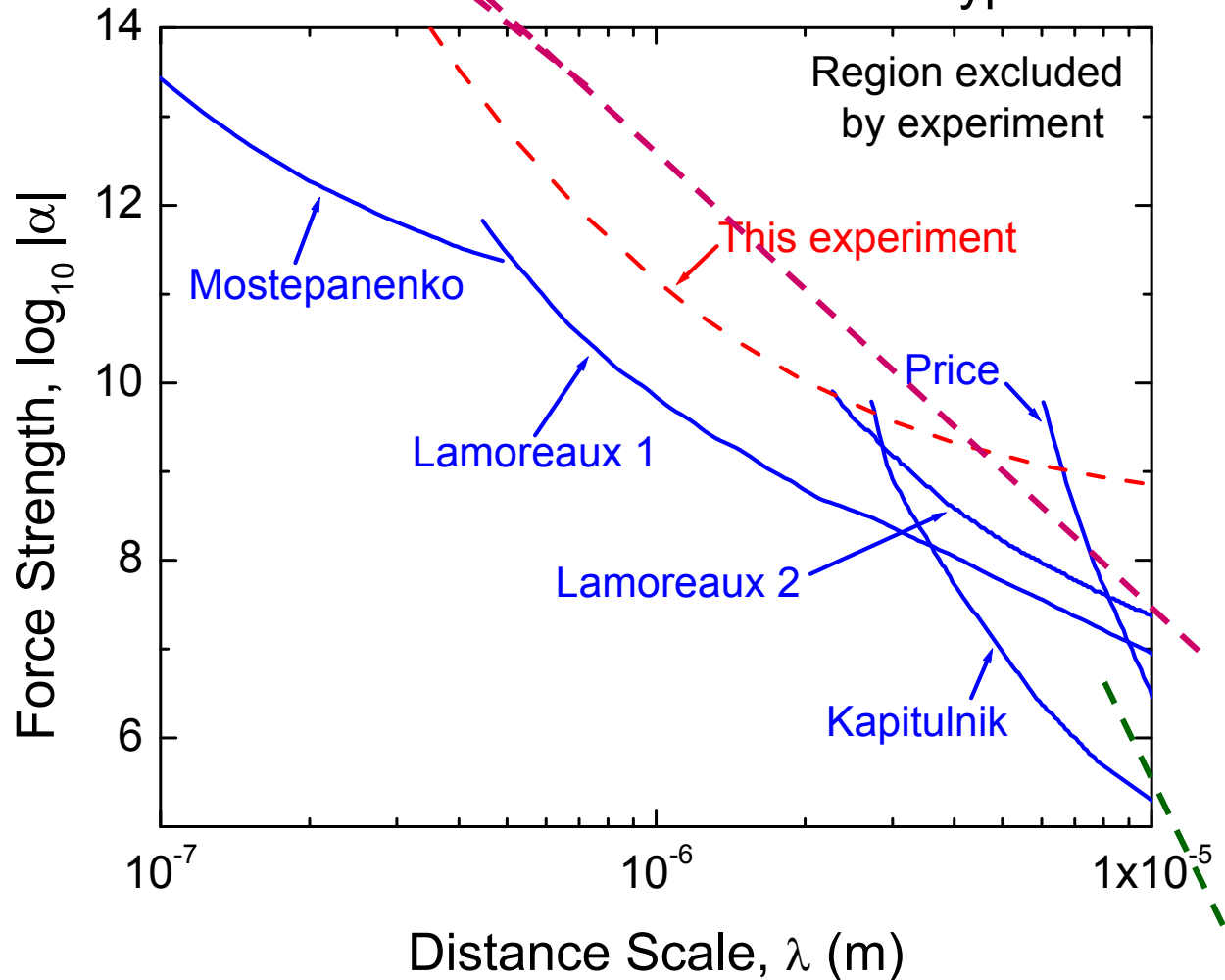
Limits on exotic forces This way to UW
 $\alpha=1$ →

$$U = \int_V \frac{Gm\rho dV}{r} (1 + \alpha e^{-r/\lambda})$$

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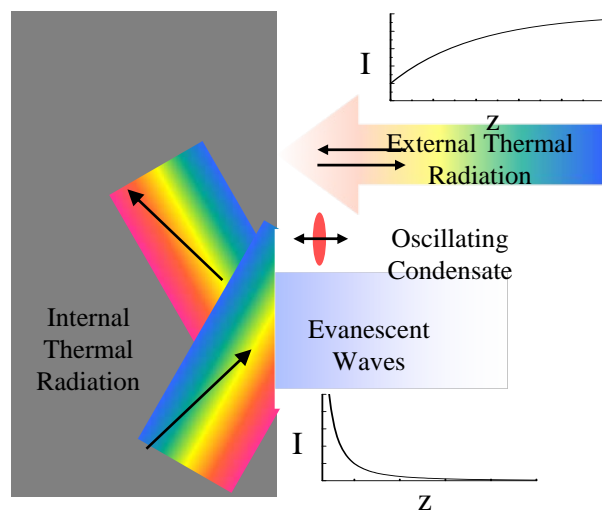
Casimir-Polder force on Atom

Current limits on forces of this type:

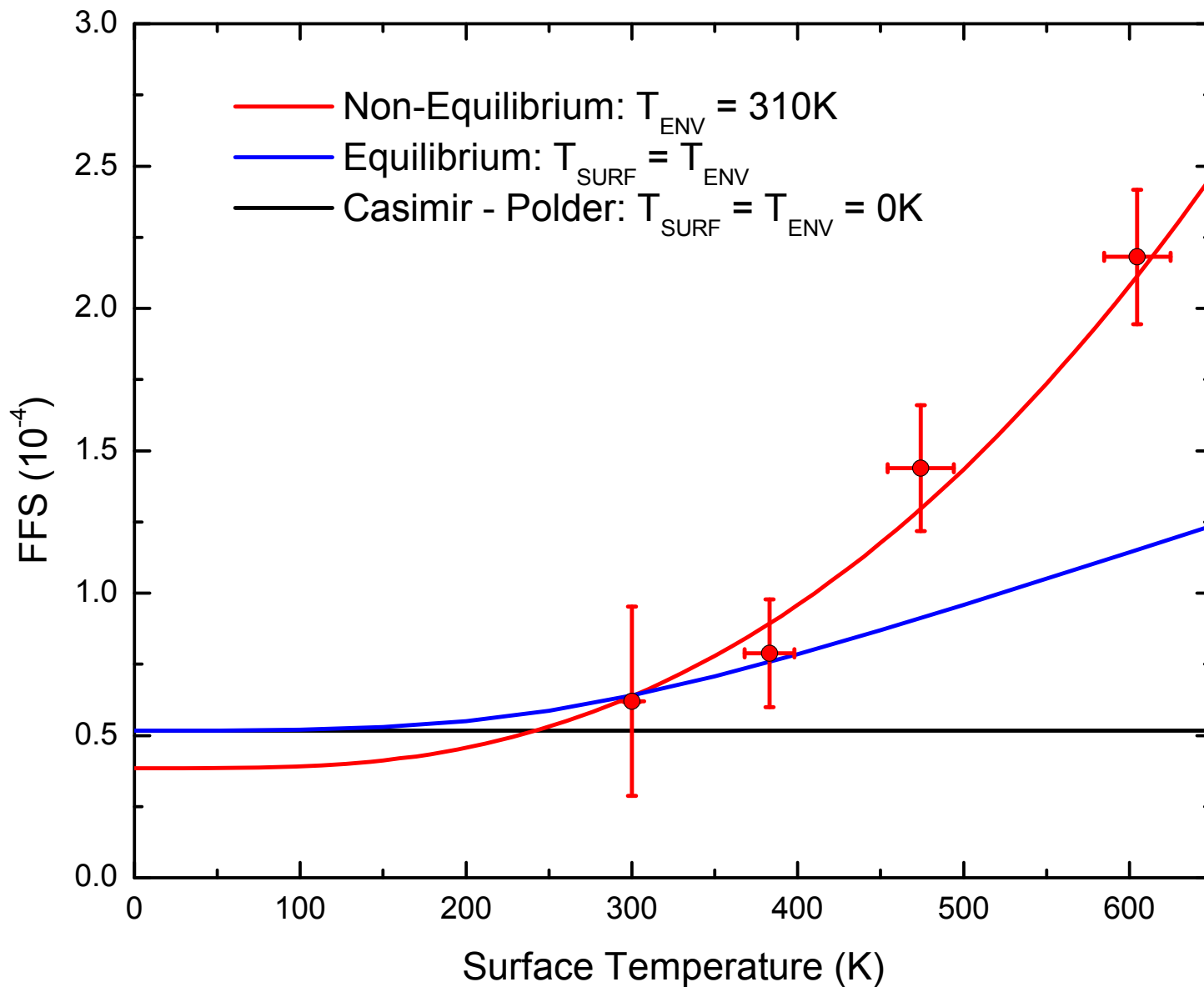


Limits on exotic forces This way to UW $\alpha=1$ \longrightarrow

Trento guys say by tinkering with temperature of “far away walls of experiment” relative to temperature of substrate, can change *sign* of total C-P force, make it repulsive.



Fixed distance to substrate, 7.5 microns



Q1:

Why hasn't the gravity from the energy of the zero-point fluctuations of fields imploded the universe?

Q2:

And caused us all to die?

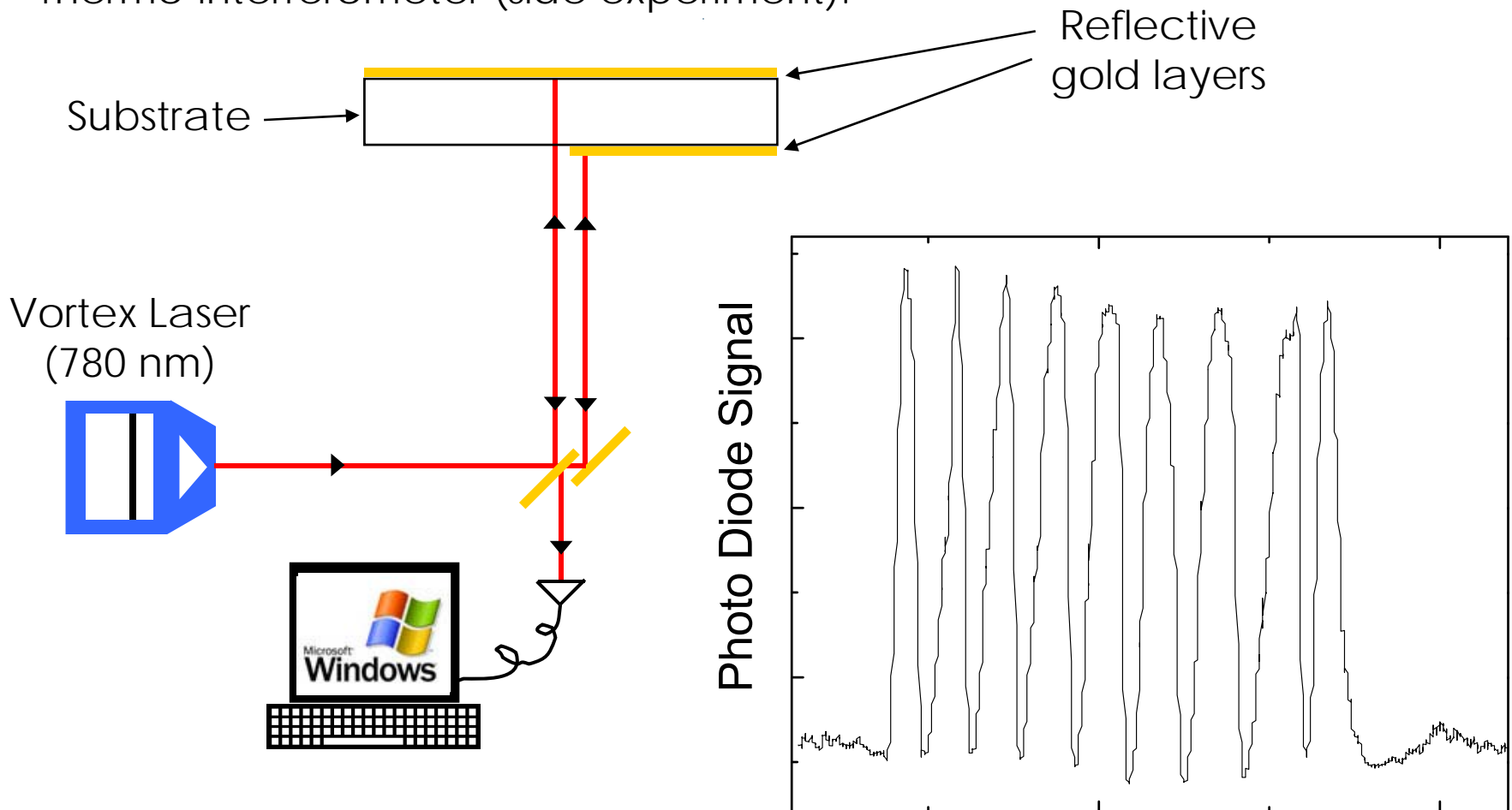
Q3:

Excrutiatingly painful deaths?

A: No clue. But maybe now we understand why the Casimir-Polder force between an atom and a surface is (usually) attractive, not repulsive.

Temperature Measurement

* Thermo-interferometer (side experiment):



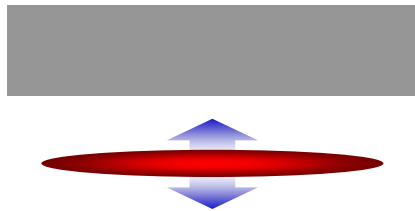
* For 1mm thick fused silica glass:

~30 °C/fringe

600 K



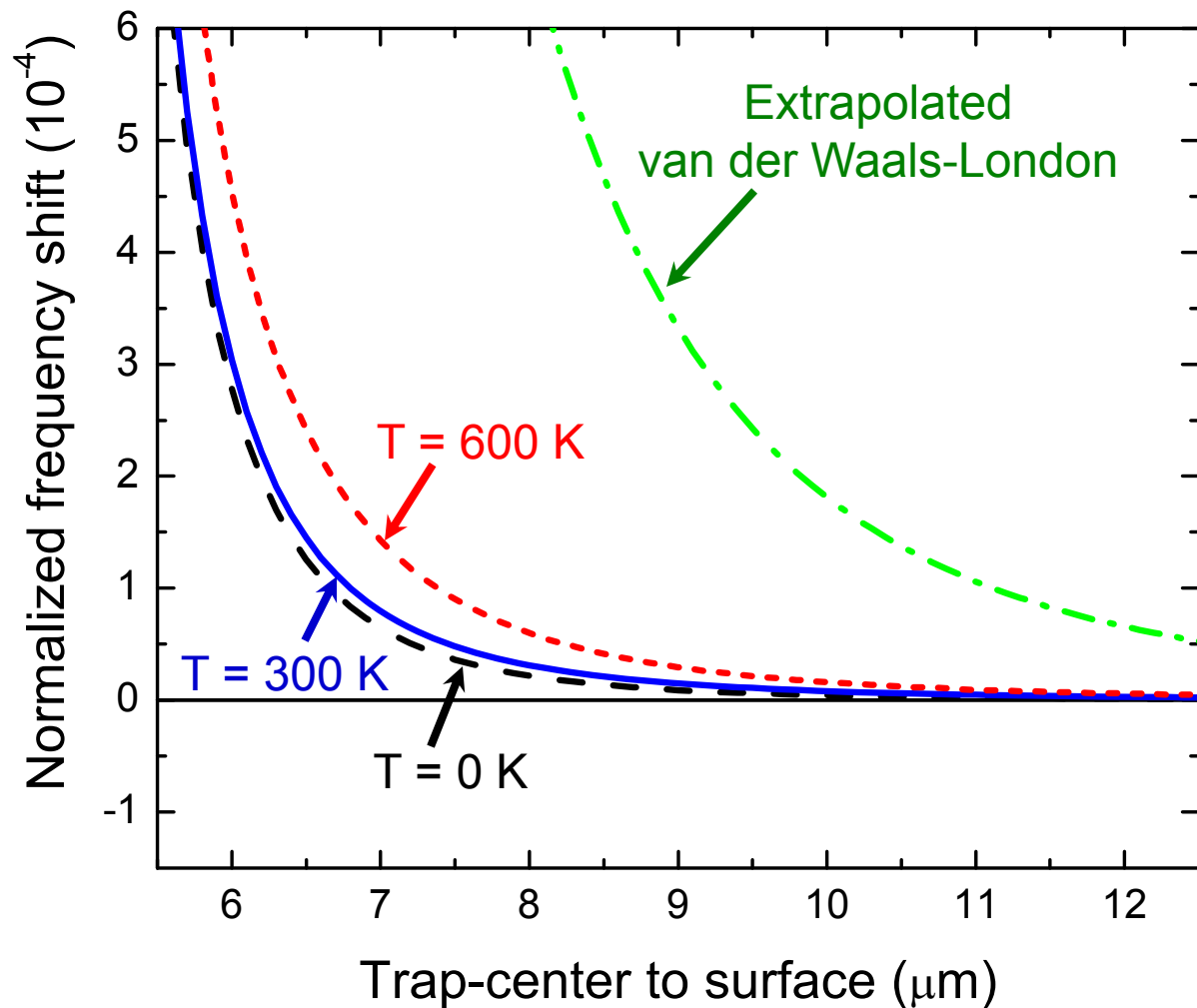
Predicted *FFS* from C-P force



Theory:

Antezza *et al.*

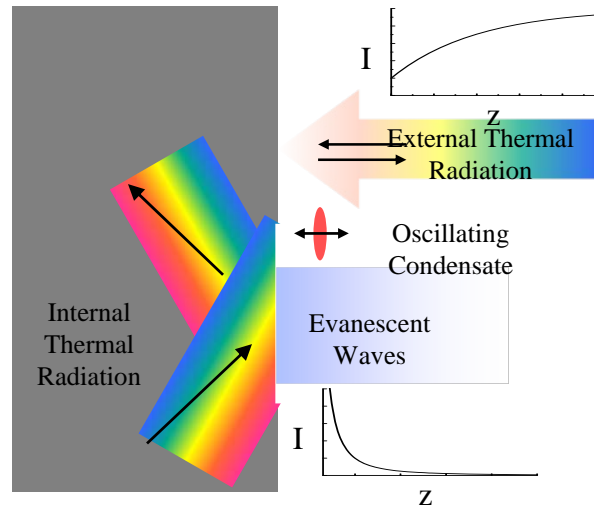
BEC width and oscillation amplitude accounted for in theory.



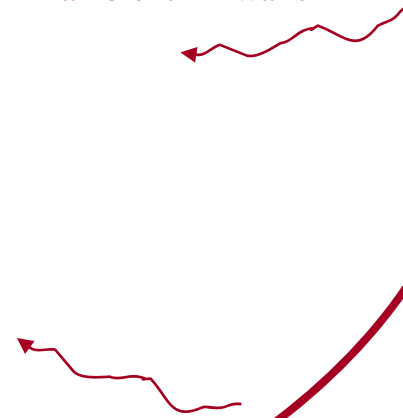
M. Antezza, L.P. Pitaevskii, and S. Stringari, Phys. Rev. A **70**, 053619 (2004)

Trento guys say by tinkering with temperature of the substrate relative to the “far away walls of experiment” can remove the cancellation of forces.

Evanescent thermal radiation at $T_{\text{substrate}}$



External thermal radiation at $T_{\text{ambient}} = T_{\text{walls}}$



Where are we going?

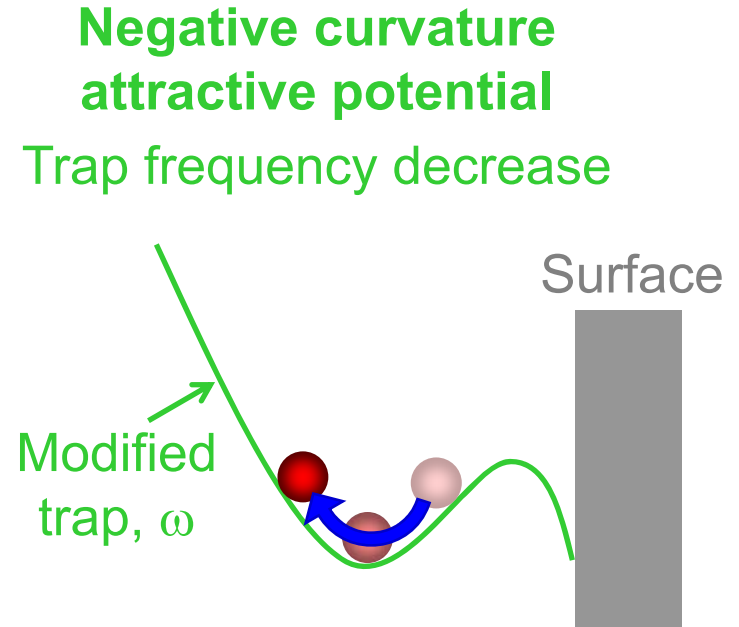
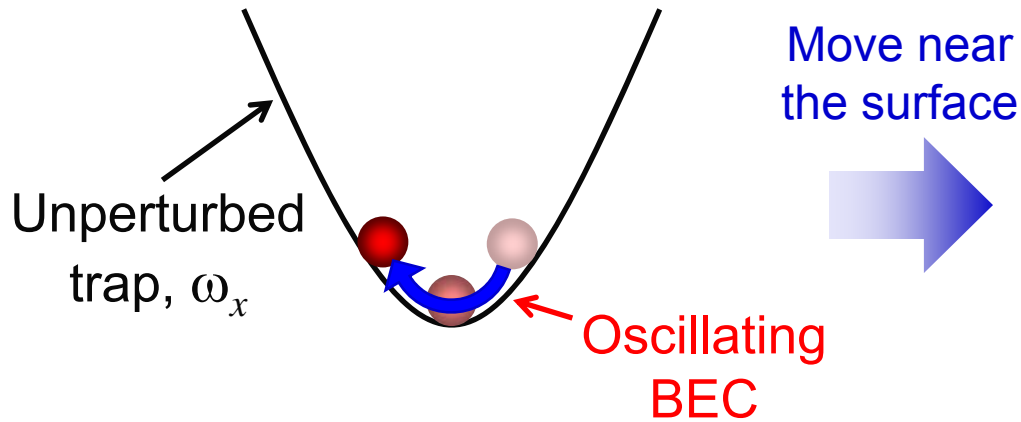
Heat environment, cool substrate, change sign of force?

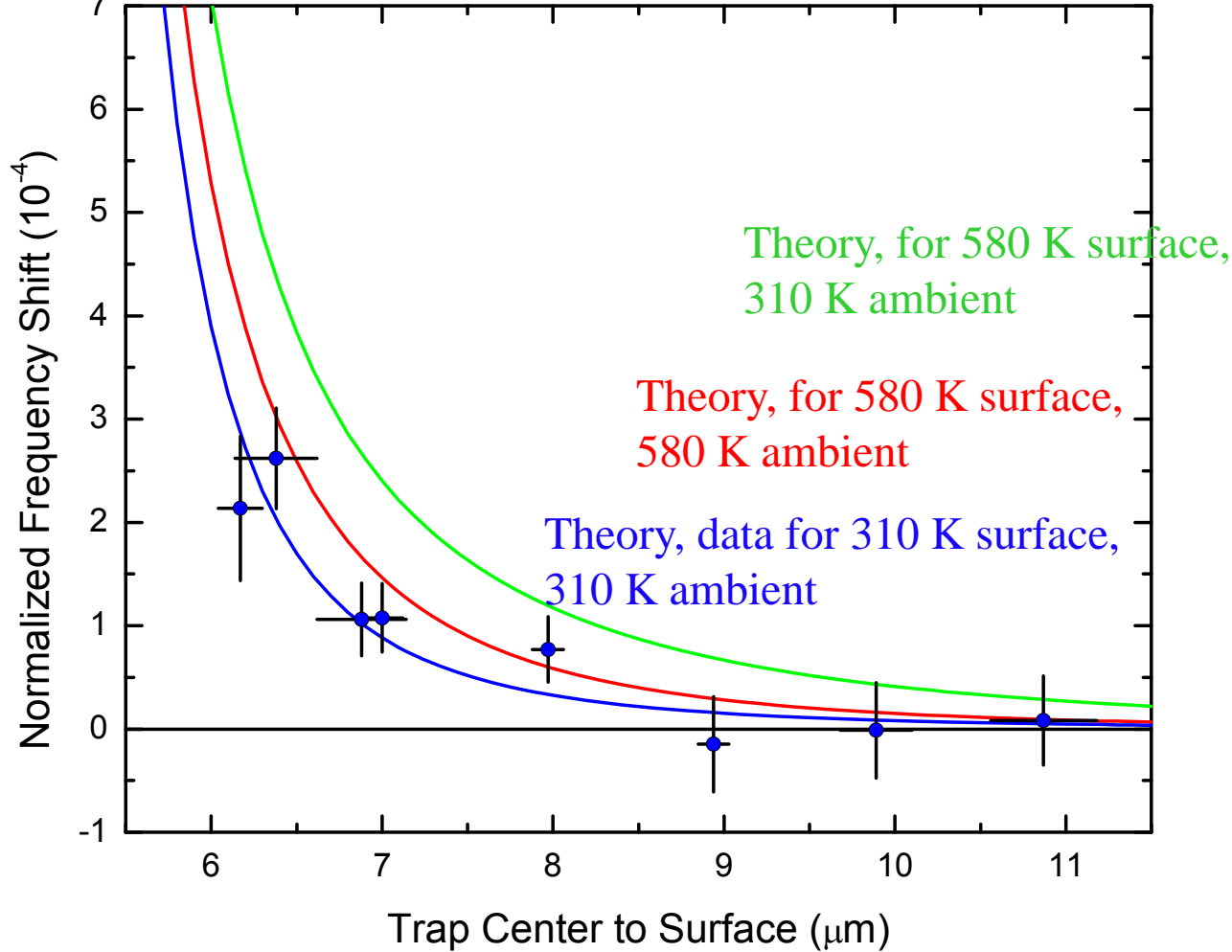
Try substrate with exotic dielectric properties, resonances.

Try spatially textured substrate.

Express trap frequency changes as normalized frequency shifts:

$$\frac{\omega_x - \omega}{\omega_x} \approx -\frac{1}{2\omega_x^2 m} \frac{d^2U}{dx^2}$$





Trento guys say by tinkering with temperature of the substrate relative to the “far away walls of experiment” can remove the cancellation of forces.