## Physics 406: Homework 4

1. **Two spin systems in contact:** A result which we will use a lot in this course is Sterling's approximation for the logarithm of a factorial, which says that

$$\ln n! \simeq n \ln n - n$$
,

where the approximation becomes better and better as  $n \to \infty$ . We will see how to prove this result in a later lecture. For the moment we just assume it. (If you want to see the proof, it's given in Appendix A of Kittel and Kroemer.)

Now consider the spin system we looked at, composed of N "spins" that each point either up or down.

- (a) Write down an expression for the multiplicity g(N,s) of states that have a given value of the spin excess parameter s, defined by  $2s = N_{\uparrow} N_{\downarrow}$ . Hence write down an expression for the logarithm of the multiplicity as a function of s.
- (b) Apply Sterling's approximation and show that for large N

$$\ln g = N \ln N - (\frac{1}{2}N + s) \ln(\frac{1}{2}N + s) - (\frac{1}{2}N - s) \ln(\frac{1}{2}N - s). \tag{1}$$

(c) Noting that  $\ln(\frac{1}{2}N+s) = \ln(\frac{1}{2}N) + \ln(1+2s/N)$ , expand to second order in 2s/N and hence show that

$$g(N,s) \simeq 2^N e^{-2s^2/N}$$
. (2)

This is a Gaussian or normal distribution: the binomial distribution becomes a Gaussian distribution for large N.

- (d) Now suppose we have two identical systems of N spins. Using Eq. (2), write down an expression for the total multiplicity of both systems together as a function of the spin excesses  $s_1$  and  $s_2$  of the two systems, assuming for the moment that the two systems are not in contact with one another. (To make the calculations easier, you can assume N is even.) Eliminate  $s_2$  in favor of the total spin excess  $s = s_1 + s_2$  and hence show that the distribution of possible values of  $s_1$  is Gaussian.
- (e) Now suppose that both systems are in a magnetic field of intensity B, so that they have energies  $U_1 = 2mBs_1$  and  $U_2 = 2mBs_2$ , where m is the dipole moment of each spin. If the two systems are now in contact with one another, so that energy can flow between them, then  $s_1$  and  $s_2$  are no longer fixed, but the total energy must be conserved, so  $s = s_1 + s_2$  is constant. What is the most likely value of  $s_1$ ?
- 2. **Entropy of a set of harmonic oscillators:** Consider the ensemble that we discussed earlier of N quantum harmonic oscillators. Each can have energy  $\varepsilon_s = s\hbar\omega$ , where  $\hbar$  and  $\omega$  are constants and s is a non-negative integer.
  - (a) Write down the expression for the multiplicity g(N,n) of states of the entire ensemble that have total internal energy  $U = n\hbar\omega$ , in a form involving factorials and hence write down the dimensionless entropy  $\sigma$  of the system.
  - (b) When N is large we can, to a good approximation, replace N-1 by N. Apply Sterling's approximation and derive an approximate expression for  $\sigma$  for large N.

- (c) Recalling the definition of the temperature  $\tau$  (in energy units)  $\tau^{-1} = \partial \sigma / \partial U|_N$ , differentiate to get an expression for  $\tau$  in terms of the internal energy. (You have to consider n to be a continuous variable to do this calculation, which is strictly speaking not correct—it is an integer. Later in the course we'll see a better derivation of this result that doesn't require us to do this kludge.)
- (d) Rearrange to show that

$$U = \frac{N\hbar\omega}{\exp(\hbar\omega/\tau) - 1}.$$

This is the internal energy of a set of *N* harmonic oscillators and, as we will show, is also the correct expression for the energy of a set of bosons (e.g., photons) in a quantum gas. From this expression we will later derive the famous black-body radiation spectrum of Rayleigh and Planck.

- (e) What is the heat capacity of the system?
- 3. **Partition function of a simple system:** Suppose a simple system has states with three energies,  $-\epsilon$ , 0, and  $+\epsilon$ . The multiplicities of the states are  $g(-\epsilon) = 1$ , g(0) = 2, and  $g(\epsilon) = 1$ . The system is put in contact with a thermal reservoir at dimensionless temperature  $\tau$  and allowed to come to equilibrium.
  - (a) Calculate the partition function Z of the system.
  - (b) Calculate the average internal energy of the system as a function of temperature.
  - (c) Show that the heat capacity is

$$C = \frac{\varepsilon^2}{\tau^2 [1 + \cosh(\varepsilon/\tau)]}.$$